

chapter 4 - Introduction to spatial variation

Principle of Mass Balance

$$\text{rate of change} = \text{net movement or flux} + \text{net rate of production}$$

4.1 Derivation of the reaction-diffusion equations

chemical species C_1, \dots, C_m

concentrations c_1, \dots, c_m

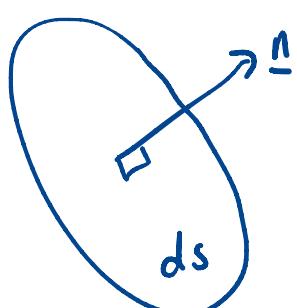
$$\text{w/o diffusion: } \frac{dc_i}{dt} = R_i(c_1, \dots, c_m)$$

$\underbrace{\quad}_{\text{net production rate}}$

- per unit volume.

concentration $c_i(x, t)$

FLUX $q_i(x, t)$



flux: defined s.t.

amount of chemical flowing through dS in the infinitesimal time interval dt

$$= n \cdot q dS dt$$

Fick's law of diffusion:

$$\frac{q}{t} = - D \nabla C$$

diffusion coeff.

For any closed volume $V \in D$ (fixed in time and space),
with boundary ∂V

$$\frac{d}{dt} \int_V c_i dV = - \int_{\partial V} q \cdot \underline{n} dS + \int_V R_i(C_1, \dots, C_m) dV$$

total amount
 c_i in V
flux out
through ∂V
net production

Divergence thm gives

$$\begin{aligned} \frac{d}{dt} \int_V c_i dV &= - \int_V \nabla \cdot q dV + \int_V R_i(C_1, \dots, C_m) dV \\ &= \int_V \left\{ \nabla \cdot (D_i \nabla c_i) + R_i(C_1, \dots, C_m) \right\} dV \end{aligned}$$

For any closed volume $V \in D$ with boundary ∂V

$$\int_V \left\{ \frac{\partial c_i}{\partial t} - \nabla \cdot (D_i \nabla c_i) - R_i \right\} dV = 0$$

therefore

$$\boxed{\frac{\partial c_i}{\partial t} = \nabla \cdot (D_i \nabla c_i) + R_i} \quad \begin{matrix} x \in D \\ i = 1, \dots, m \end{matrix}$$

System of reaction-diffusion eqns.

Require : ICs and BCs to close the system.

think further about the step \Rightarrow

suppose $\frac{\partial c_i}{\partial t} - \nabla \cdot (D_i \nabla c_i) - R_i \neq 0$ some

WLOG suppose the above is positive at $\underline{x}^* \in D$

Then $\exists \varepsilon > 0$ s.t. $\frac{\partial c_i}{\partial t} - \nabla \cdot (D_i \nabla c_i) - R_i > 0$

for $\underline{x} \in B_\varepsilon(\underline{x}^*)$.

In this case

$$\int_{B_\varepsilon(\underline{x}^*)} \left[\frac{\partial c_i}{\partial t} - \nabla \cdot (D_i \nabla c_i) - R_i \right] dV > 0$$

contradicts our original assumption, and so
the RD eqns hold.

① Single species, 1D, no reactions

$$\frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial x^2} \quad (\text{diffusion eqn})$$

$$\text{length scale } L \quad \text{timescale } T = \frac{L^2}{D}$$

$$\text{cell } L \sim 10^{-5} \text{ m} = 10^{-3} \text{ cm}$$

$$D \sim 10^{-7} \text{ cm}^2 \text{ s}^{-1}$$

$$\text{Then } T \sim \frac{10^{-6} \text{ cm}^2}{10^{-7} \text{ cm}^2 \text{ s}^{-1}} \sim 10 \text{ s}$$

$$\text{If } L \mapsto 10L \quad T \mapsto 10^2 T$$

$$L \mapsto 100L \quad T \mapsto 10^4 T$$

4.2 Chemotaxis

chemotaxis - movements up chemical gradients

$$\text{diffusive flux } J_D = -D_n \nabla n$$

$$\text{chemotactic flux } J_C = n \underline{\chi(c)} \nabla c = n \nabla \underline{\Phi(c)}$$

chemotactic coefficient

monotonic inc

$$\text{eg } \chi(c) = \chi,$$

$$\frac{\chi_0}{c}, \frac{\chi_0}{(k+c)^2}$$

$$\begin{aligned} \text{Total flux : } J &= J_D + J_C \\ &= -D_n \nabla n + n \chi(c) \nabla c \end{aligned}$$

Typical model

$$\frac{\partial n}{\partial t} = \nabla \cdot (D_n \nabla n) - \nabla \cdot (n \chi(c) \nabla c) + \underbrace{f(n, c)}$$

$$\frac{\partial c}{\partial t} = \nabla \cdot (D_c \nabla c) + \lambda n - \mu c$$

eg. logistic growth.

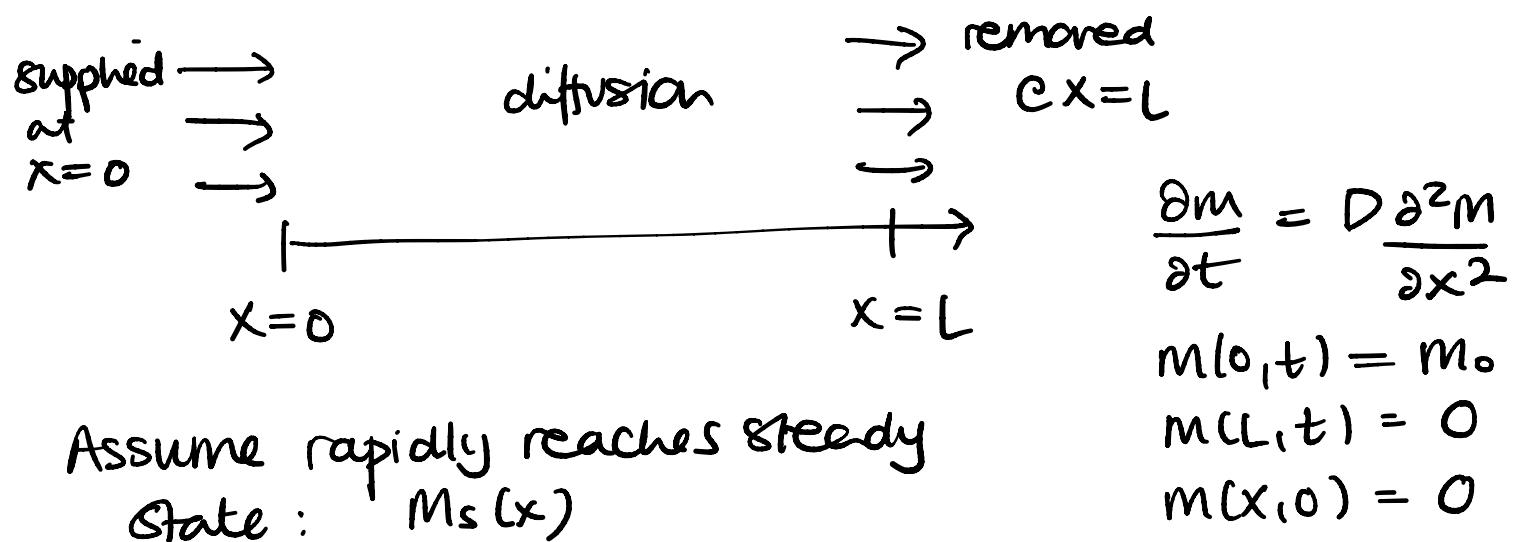
4.3 Positional information and pattern formation

Theories

① Positional information - Wolpert
"French Flag Model"

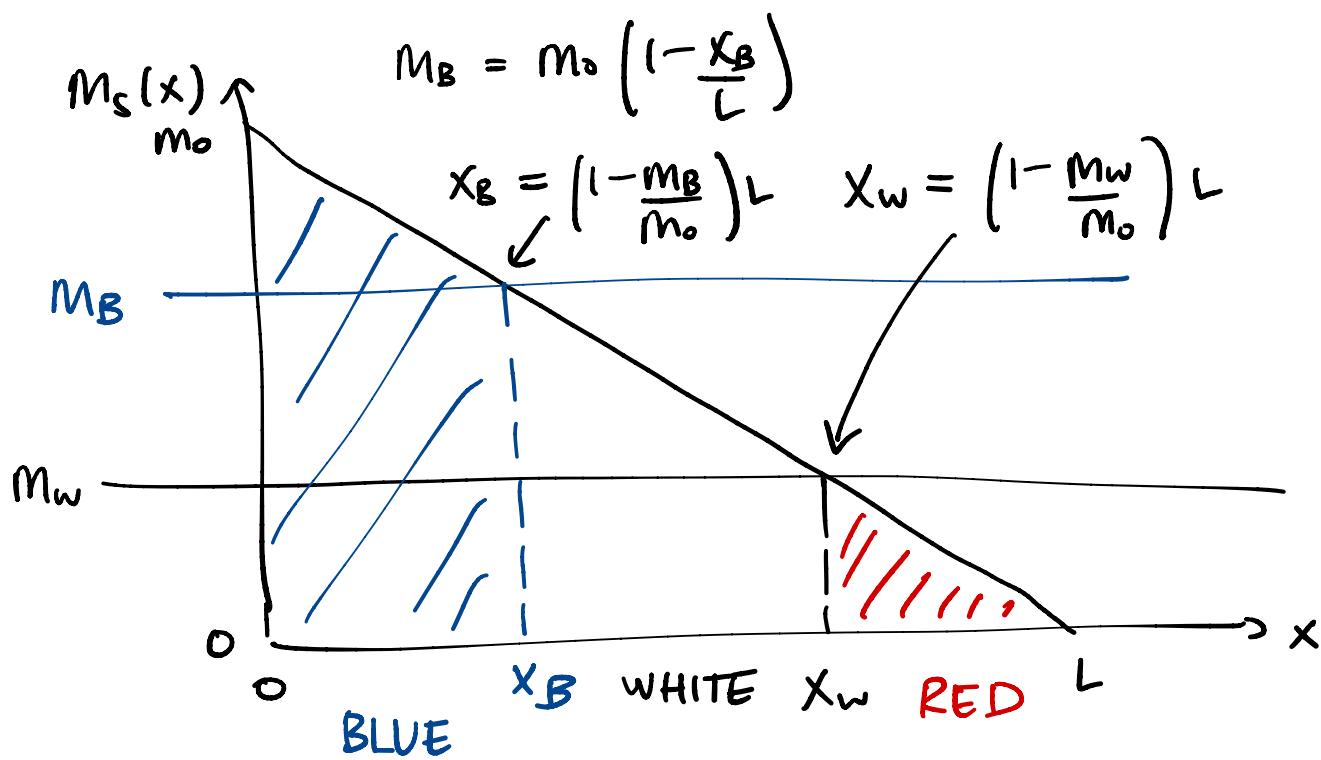
② Diffusion-driven instability - Thuring (§ 6)

the French Flag Model



Assume rapidly reaches steady state : $m_s(x)$

$$\frac{d^2 m_s}{dx^2} = 0 \Rightarrow m_s(x) = m_0 \left(1 - \frac{x}{L}\right)$$



4.4 Minimum domains for spatial structure

$$\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2} + f(u)$$



$$f(u) = ru \left(1 - \frac{u}{q}\right) - \frac{u^2}{1+u^2}$$

Boundary conditions $u(0, t) = 0$
 $u(L, t) = 0$

Q - If we start with a small initial distribution, will there be an outbreak, or will the popⁿ die out?

Suppose $0 \leq u(x, 0) \ll 1$ $f(u) \approx f'(0)u = ru$

Then $\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2} + \underbrace{f'(0)u}_{=r} u$

Seek a solution $u(x, t) = \sum_{n=1}^{\infty} a_n(t) \sin\left(\frac{n\pi x}{L}\right)$

then $\frac{da_n}{dt} = \left[-\frac{Dn^2\pi^2}{L^2} + f'(0) \right] a_n$
 $\underbrace{-\frac{Dn^2\pi^2}{L^2}}$
 $\underbrace{+ f'(0)}$
 $\underbrace{a_n}$

$$\Rightarrow a_n(t) = a_n(0) e^{-\sigma_n t}$$

$$\therefore u(x, t) = \sum_{n=1}^{\infty} a_n(0) e^{\underbrace{(f'(0) - \frac{Dn^2\pi^2}{L^2})}_{} t} \sin\left(\frac{n\pi x}{L}\right) + O_n$$

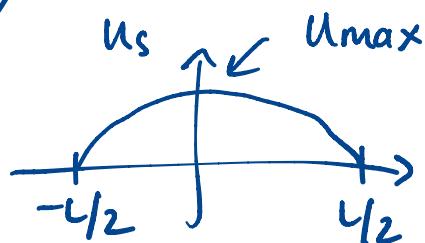
For the outbreak to die out : $0_n < 0 \forall n$

$$f'(0) - \frac{Dn^2\pi^2}{L^2} < 0 \quad \forall n$$

i.e. $L \leq \sqrt{\frac{D\pi^2}{f'(0)}} := L_{crit}$

Q - how does the steady state pattern depend on the domain size?

shift coordinates



At steady state $0 = D \frac{\partial^2 u}{\partial x^2} + f(u)$

Multiply by $\frac{\partial u}{\partial x}$, and integrate :

$$0 = \int D \left(\frac{\partial u}{\partial x} \right) \left(\frac{\partial^2 u}{\partial x^2} \right) dx + \int f(u) \frac{\partial u}{\partial x} dx$$

$$\frac{1}{2} D \left(\frac{\partial u}{\partial x} \right)^2 + F(u) = F(u_{max})$$

$\left(\text{for } x > 0, \frac{\partial u}{\partial x} < 0 \right)$

$$F(u) := \int_0^u f(y) dy$$

$$F'(u) = f(u)$$

$$\frac{\partial u}{\partial x} = \pm \sqrt{\frac{2}{D}} \sqrt{F(u_{max}) - F(u)}$$

$$2 \int_0^{L/2} dx = - \sqrt{2D} \int_{u_{max}}^0 \frac{1}{\sqrt{F(u_{max}) - F(\bar{u})}} d\bar{u}$$

$$\frac{L}{\sqrt{2D}} = \int_0^{u_{max}} \frac{1}{\sqrt{F(u_{max}) - F(\bar{u})}} d\bar{u}$$

