

Q1

$\int_0^{\pi/2} e^{ix \cos t} dt$: stationary phase for first term;
steepest descents for more.

$\int_0^1 \ln t e^{ixt} dt$: steepest descents (IBP & stationary phase don't work)

$\int_0^\infty t^{-1/2} e^{-t} dt$: $\int_0^\infty - \int_x^\infty$ and IBP on \int_x^∞ .

$\int_0^{\pi/2} e^{-x \sin^2 t} dt$: Laplace's method.

$\int_0^1 e^{ixt} e^{-1/t} dt$: steepest descents having set $s = e^{-1/t}$.

$\int_0^{10} \frac{e^{-xt}}{1+t} dt$: Taylor expand integrand & integrate term-by-term

$\int_0^{\pi/2} \frac{dt}{\cos^2 t + x \sin^2 t}$: $\int_0^{\pi/2} = \int_0^{\pi/2 - \delta} + \int_{\pi/2 - \delta}^{\pi/2}$, where $x \ll \delta \ll 1$.

$\int_0^1 \frac{\sin(xt)}{t} dt$: Taylor expand integrand & integrate term-by-term.

$\int_x^\infty t^{a-1} e^{-t} dt$: $\int_x^\infty = \int_0^\infty - \int_0^x$ and \int_0^x on \int_0^x for $\text{Re}(a) > 0$;

much more tricky for $\text{Re}(a) \leq 0$!

$\int_0^1 \frac{\ln t}{x+t} dt$: $\int_0^1 = \int_0^\delta + \int_\delta^1$, where $x \ll \delta \ll 1$.

Q2

$$I_1(x) = \int_{-1}^1 e^{-x \cos ht} dt \text{ as } x \rightarrow \infty.$$

Maximum of $\phi(t) = -\cos ht$ at $t = 0 \Rightarrow$ split integral

$$I_1(x) = \left(\underbrace{\int_{-1}^{-\epsilon}}_{I_{11}(x)} + \underbrace{\int_{-\epsilon}^{\epsilon}}_{I_{12}(x)} + \underbrace{\int_{\epsilon}^1}_{I_{11}(x)} \right) e^{-x \cos ht} dt, \text{ where } \epsilon \ll 1$$

by symmetry

$$I_{11}(x) = \int_{\epsilon}^1 e^{-x \cos ht} dt = O(e^{-x \cos h \epsilon}) = O(e^{-x} \cdot e^{-x \epsilon^2/2})$$

$$\begin{aligned} I_{12}(x) &= \int_{-\epsilon}^{\epsilon} e^{-x[1+t^2/2 + O(t^4)]} dt \quad (\text{OK } \because \epsilon \ll 1) \\ &= e^{-x} \int_{-\epsilon}^{\epsilon} e^{-xt^2/2} e^{O(xt^4)} dt \\ &= e^{-x} \int_{-\epsilon}^{\epsilon} e^{-xt^2/2} (1 + O(xt^4)) dt \quad (\text{provided } x\epsilon^4 \ll 1) \\ &= \frac{\sqrt{x}}{2} e^{-x} \int_{-\epsilon\sqrt{x}}^{\epsilon\sqrt{x}} e^{-s^2} (1 + O(s^4/x)) ds \quad (s = \sqrt{x}t) \\ &= \frac{\sqrt{x}}{2} e^{-x} \left[\int_{-\epsilon\sqrt{x}}^{\epsilon\sqrt{x}} e^{-s^2} ds + O(x^{-1}) \right] \\ &= \frac{\sqrt{x}}{2} e^{-x} \left[\int_{-\infty}^{\infty} e^{-s^2} ds + O(x^{-1}) \right] \quad (\text{provided } \epsilon\sqrt{x} \gg 1) \\ &= \frac{\sqrt{x}}{2} e^{-x} \left[\sqrt{\pi} + O(x^{-1}) \right] \quad \text{as } x \rightarrow \infty \end{aligned}$$

$$\epsilon\sqrt{x} \gg 1 \Rightarrow I_{11}(x) \ll I_{12}(x) \text{ as } x \rightarrow \infty$$

Expansions self consistent provided $x^{-1/2} \ll \epsilon \ll x^{-1/4}$ and therefore

$$\underline{\underline{I_1(x) \sim \frac{\sqrt{\pi}}{2} e^{-x} \text{ as } x \rightarrow \infty}}$$

$$I_2(x) = \int_{-\pi/2}^{\pi/2} e^{-x(t^2 - \sin^2 t)} dt \text{ as } x \rightarrow \infty.$$

$\phi(t) = \sin^2 t - t^2 < 0$ for $t \neq 0$, with $\phi(0) = 0$, so interior maximum at $t = 0$.

$\phi(t) = (t - \frac{t^3}{3} + O(t^5))^2 - t^2 = -\frac{1}{3}t^4 + O(t^6)$ as $t \rightarrow 0$, so this is a degenerate case.

Split: $I_2(x) = \left(\int_{-\pi/2}^{-\epsilon} + \int_{-\epsilon}^{\epsilon} + \int_{\epsilon}^{\pi/2} \right) e^{x\phi(t)} dt, \epsilon \ll 1$
 $I_{21}(x) \quad I_{22}(x) \quad I_{21}(x)$ by symmetry

$$I_{21}(x) = \int_{\epsilon}^{\pi/2} e^{x\phi(t)} dt = O(e^{x\phi(\epsilon)}) = O(e^{-x\epsilon^4})$$

$$\begin{aligned} I_{22}(x) &= \int_{-\epsilon}^{\epsilon} e^{-x(t^4/3 + O(t^6))} dt \quad (\text{OK } \because \epsilon \ll 1) \\ &= \int_{-\epsilon}^{\epsilon} e^{-xt^4/3} (1 + O(xt^6)) dt \quad (\text{provided } x\epsilon^6 \ll 1) \\ &= \left(\frac{3}{x}\right)^{1/4} \int_{-(x/3)^{1/4}\epsilon}^{(x/3)^{1/4}\epsilon} e^{-s^4} (1 + O(x^{-1/2}s^6)) ds \quad (s = (x/3)^{1/4}t) \\ &\quad \downarrow (\text{exp. small error for } x^{1/4}\epsilon \gg 1) \\ &= \left(\frac{3}{x}\right)^{1/4} \left[\int_{-\infty}^{\infty} e^{-s^4} ds + O(x^{-1/2}) \right] \quad (\text{provided } x^{1/4}\epsilon \gg 1) \\ &= \left(\frac{3}{x}\right)^{1/4} \left[2 \frac{\Gamma(1/4)}{4} + O(x^{-1/2}) \right] \quad (\text{by hint}) \end{aligned}$$

$x^{1/4} \epsilon \gg 1 \Rightarrow I_{21}(x) \ll I_{22}(x)$ as $x \rightarrow \infty$

Expansions self-consistent provided $x^{-1/4} \ll \epsilon \ll x^{-1/6}$ and therefore

$$\underline{\underline{I_2(x) \sim \frac{\Gamma(1/4)}{2} \left(\frac{3}{x}\right)^{1/4} \text{ as } x \rightarrow \infty}}$$

$$I_3(x) = \int_0^{\infty} e^{-2t - x/t^2} dt \text{ as } x \rightarrow \infty$$

$$\frac{d}{dt} \left(2t + \frac{x}{t^2} \right) = 2 \left(1 - \frac{x}{t^3} \right) = 0 \text{ when } t = x^{1/3}$$

⇒ this is a moveable maximum problem.

Let $y = x^{1/3}$, $t = x^{1/3} u$, then $I_3(x) = y \int_0^{\infty} e^{-y(2u + 1/u^2)} du$

$$\phi(u) = -2u - \frac{1}{u^2} \Rightarrow \phi'(u) = 2 \left(\frac{1}{u^3} - 1 \right) = 0 \text{ when } u = 1$$

$$\phi(u) = \phi(1) + \frac{1}{2!} \phi''(1)(u-1)^2 + O((u-1)^3) \text{ as } u \rightarrow 1, \text{ where } \phi(1) = -3, \phi''(1) = -6.$$

Split: $I_3(x) = \underbrace{\int_0^{1-\epsilon}}_{I_{31}(y)} + \underbrace{\int_{1-\epsilon}^{1+\epsilon}}_{I_{32}(y)} + \underbrace{\int_{1+\epsilon}^{\infty}}_{I_{33}(y)} y e^{y\phi(u)} du, \epsilon \ll 1$

$$I_{31}(y) = O(y e^{y\phi(1-\epsilon)}) = O(y e^{-3y} e^{-3y\epsilon^2})$$

$$I_{33}(y) = O(y e^{y\phi(1+\epsilon)}) = O(y e^{-3y} e^{-3y\epsilon^2})$$

$$\begin{aligned} I_{32}(y) &= \int_{1-\epsilon}^{1+\epsilon} y e^{y(-3 - 3(u-1)^2 + O((u-1)^3))} du \quad (\text{OK } \because \epsilon \ll 1) \\ &= y e^{-3y} \int_{1-\epsilon}^{1+\epsilon} e^{-3y(u-1)^2} (1 + O(y(u-1)^3)) du \quad (\text{provided } y\epsilon^3 \gg 1) \\ &= y e^{-3y} \frac{1}{\sqrt{3y}} \int_{-\epsilon\sqrt{3y}}^{\epsilon\sqrt{3y}} e^{-s^2} (1 + O(y^{-1/2} s^3)) ds \quad (\sqrt{3y}(u-1) = s) \\ &= \sqrt{\frac{y}{3}} e^{-3y} \left[\int_{-\infty}^{\infty} e^{-s^2} ds + O(y^{-1/2}) \right] \quad (\text{provided } \epsilon\sqrt{3y} \gg 1) \end{aligned}$$

$\epsilon\sqrt{3y} \gg 1 \Rightarrow I_{31}(y), I_{33}(y) \ll I_{32}(y) \text{ as } y \rightarrow \infty$

Expansions self-consistent provided $x^{-1/6} \ll \epsilon \ll x^{-1/9}$ and therefore

$$\underline{\underline{I_3(x) \sim \sqrt{\frac{\pi}{3}} x^{1/6} e^{-3x^{1/3}} \text{ as } x \rightarrow \infty}}$$

$$J_1(x) = \int_0^1 \cosh(t^2) e^{ixt^2} dt \quad \text{as } x \rightarrow \infty.$$

$\psi(t) = t^2$ has a stationary point at $t=0$, so split

$$J_1(x) = \left(\underbrace{\int_0^\varepsilon}_{J_{11}(x)} + \underbrace{\int_\varepsilon^1}_{J_{12}(x)} \right) \cosh(t^2) e^{ixt^2} dt, \quad \varepsilon \ll 1$$

$$J_{11}(x) = \int_0^\varepsilon [1 + o(t^4)] e^{ixt^2} dt \quad (\text{OK } \because \varepsilon \ll 1)$$

$$= \frac{1}{\sqrt{x}} \int_0^{\varepsilon\sqrt{x}} [1 + o(\frac{s^4}{x^2})] e^{is^2} ds \quad (s = \sqrt{x}t)$$

$$= \frac{1}{\sqrt{x}} \int_0^\infty e^{is^2} ds + O\left(\frac{1}{\varepsilon x}\right) \quad (\text{provided } \varepsilon\sqrt{x} \gg 1)$$

$O(1/(\text{sqrt}(x) \cdot \text{eps} \cdot \text{sqrt}(x)))$

the last equality because

$$\begin{aligned} \frac{1}{\sqrt{x}} \int_0^\infty e^{is^2} ds &= \frac{1}{\sqrt{x}} \int_{\varepsilon\sqrt{x}}^\infty \underbrace{\frac{1}{2is}}_f \underbrace{2ise^{is^2}}_{g'} ds \\ &= \frac{1}{\sqrt{x}} \left[\underbrace{\frac{1}{2is} e^{is^2}}_f \underbrace{\Big|_{\varepsilon\sqrt{x}}^\infty}_g - \int_{\varepsilon\sqrt{x}}^\infty \underbrace{\frac{-1}{2is^2}}_{f'} \underbrace{e^{is^2}}_g ds \right] \\ &= \underbrace{-\frac{1}{2i\varepsilon x} e^{i\varepsilon^2 x}}_{O(\frac{1}{\varepsilon x})} + \underbrace{\frac{1}{2i\varepsilon\sqrt{x}} \int_{\varepsilon\sqrt{x}}^\infty \frac{1}{s^2} e^{is^2} ds}_{O(\frac{1}{\varepsilon x})} \end{aligned}$$

$$\text{and } \frac{1}{x^{5/2}} \int_0^{\varepsilon\sqrt{x}} s^4 e^{is^2} ds = O\left(\frac{(\varepsilon\sqrt{x})^3}{x^{5/2}}\right) = O\left(\frac{\varepsilon^3}{x}\right) \ll 1$$

By parts ... first term

(provided $\varepsilon^3 \ll x$)

$$\begin{aligned} J_{12}(x) &= \int_\varepsilon^1 \underbrace{\frac{\cosh(t^2)}{2ixt}}_f \underbrace{2ixt e^{ixt^2}}_{g'} dt \\ &= \underbrace{\frac{\cosh(t^2)}{2ixt} e^{ixt^2}}_f \underbrace{\Big|_\varepsilon^1}_g - \int_\varepsilon^1 \underbrace{\frac{\partial}{\partial t} \left(\frac{\cosh(t^2)}{2ixt} \right)}_{f'} \underbrace{e^{ixt^2}}_g dt \\ &= O\left(\frac{1}{\varepsilon x}\right) \quad \text{with } O(1/x) \text{ as } x \rightarrow \infty \text{ by RLL} \end{aligned}$$

Expansions self-consistent provided $x^{-1/2} \ll \varepsilon \ll x^{-1/3}$ ($\Rightarrow x^{-1/2} \gg \frac{1}{\varepsilon x}$) and therefore (by hint)

$$\underline{\underline{J_1(x) \sim \frac{e^{i\pi/4}}{2} \sqrt{\frac{\pi}{x}} \text{ as } x \rightarrow \infty}}$$

$$J_2(x) = \text{Re} \left[J_4(x) = \int_0^1 \tan(t) e^{ixt^4} dt \right] \text{ as } x \rightarrow \infty$$

$\psi(t) = t^4$ has a stationary point at $t = 0$, so split

$$J_4(x) = \left(\int_0^\epsilon + \int_\epsilon^1 \right) \tan(t) e^{ixt^4} dt, \quad \epsilon \ll 1.$$

$$J_{41}(x) = \int_0^\epsilon (t + O(t^3)) e^{ixt^4} dt \quad (\text{OK } \because \epsilon \ll 1)$$

$$= x^{-1/4} \int_0^{x^{1/4}\epsilon} \left(\frac{s}{x^{1/4}} + O\left(\frac{s^3}{x^{3/4}}\right) \right) e^{is^4} ds \quad (s = x^{1/4}t)$$

$$= x^{-1/2} \int_0^\infty s e^{is^4} ds + O\left(\frac{1}{\epsilon^2 x}\right) \quad (\text{provided } x^{1/4}\epsilon \gg 1)$$

the last equality because

need $1/(\epsilon^2 x) \ll 1$

$$x^{-1/2} \int_{x^{1/4}\epsilon}^\infty s e^{is^4} ds = x^{-1/2} \int_{x^{1/4}\epsilon}^\infty \underbrace{\frac{1}{4is^3}}_f \underbrace{+ 4is^3 e^{is^4}}_{g'} ds$$

$$= x^{-1/2} \left[\underbrace{\frac{1}{4is^3} e^{is^4}}_{O\left(\frac{1}{x^{1/2}\epsilon^3}\right)} \Big|_{x^{1/4}\epsilon}^\infty - \int_{x^{1/4}\epsilon}^\infty \underbrace{\frac{-2}{4is^3}}_{f'} \underbrace{e^{is^4}}_g ds \right]$$

$$= O\left(\frac{1}{x\epsilon^2}\right)$$

$$\text{and } x^{-1} \int_0^{x^{1/4}\epsilon} s^3 e^{is^4} ds = \frac{1}{4i^2} e^{is^4} \Big|_0^{x^{1/4}\epsilon} = O\left(\frac{1}{x}\right)$$

$$J_{42}(x) = \int_\epsilon^1 \underbrace{\frac{\tan(t)}{4ixt^3}}_f \underbrace{4ixt^3 e^{ixt^4}}_{g'} dt$$

$$= \underbrace{\frac{\tan(t)}{4ixt^3} e^{ixt^4}}_g \Big|_\epsilon^1 - \int_\epsilon^1 \underbrace{\frac{\partial}{\partial t} \left(\frac{\tan(t)}{4ixt^3} \right)}_{f'} \underbrace{e^{ixt^4}}_g dt$$

$$= O\left(\frac{1}{x\epsilon^2}\right) \quad O(1/x) \text{ as } x \rightarrow \infty \text{ by RLL}$$

Expansions self consistent provided $\epsilon \gg x^{-1/4} (\Rightarrow x^{-1/2} \gg \frac{1}{\epsilon^2 x})$ and therefore (by hint)

$$J_2(x) \sim \frac{\cos \frac{\pi}{4}}{4} \Gamma\left(\frac{1}{4}\right) x^{-1/2} \text{ as } x \rightarrow \infty$$

$$J_3(x) = \int_0^1 e^{ix\psi(t)} dt \text{ as } x \rightarrow \infty$$

$$\psi(t) = t - \sin t = t - \left[t - \frac{t^3}{3!} + O(t^5) \right] = \frac{t^3}{6} + O(t^5) \text{ as } t \rightarrow 0$$

and $\psi'(t) = 1 - \cos t > 0$ for $0 < t \leq 1$, so $t=0$ is only stationary point and we split

$$J_3(x) = \underbrace{\left(\int_0^\varepsilon \right)}_{J_{31}(x)} + \underbrace{\left(\int_\varepsilon^1 \right)}_{J_{32}(x)} e^{ix\psi(t)} dt, \quad \varepsilon \ll 1.$$

$$\begin{aligned} J_{31}(x) &= \int_0^\varepsilon e^{ix \left[\frac{t^3}{6} + O(t^5) \right]} dt \quad (\text{OK } \because \varepsilon \ll 1) \\ &= \left(\frac{6}{x}\right)^{1/3} \int_0^{\varepsilon \left(\frac{x}{6}\right)^{1/3}} e^{is^3} \left(1 + O\left(x \left(\frac{s}{x^{1/3}}\right)^5\right)\right) ds \quad \left(s = \left(\frac{x}{6}\right)^{1/3} t\right) \\ &\quad \text{with } x \varepsilon^5 \ll 1 \\ &= \left(\frac{6}{x}\right)^{1/3} \int_0^\infty e^{is^3} ds + O\left(\frac{1}{\varepsilon^{1/2}}\right) \quad (\text{provided } x^{1/3} \varepsilon \gg 1) \end{aligned}$$

the last equality because

$$\begin{aligned} x^{-1/3} \int_{\varepsilon \left(\frac{x}{6}\right)^{1/3}}^\infty e^{is^3} ds &= x^{-1/3} \int_{\varepsilon \left(\frac{x}{6}\right)^{1/3}}^\infty \underbrace{\frac{1}{3is^2}}_f \underbrace{3is^2 e^{is^3}}_{g'} ds \\ &= x^{-1/3} \left[\underbrace{\frac{1}{3is^2}}_f \underbrace{e^{is^3}}_g \Big|_{\varepsilon \left(\frac{x}{6}\right)^{1/3}}^\infty - \int_{\varepsilon \left(\frac{x}{6}\right)^{1/3}}^\infty \underbrace{\frac{\partial}{\partial s} \left(\frac{1}{3is^2}\right)}_{f'} \underbrace{e^{is^3}}_g ds \right] \\ &= O\left(\frac{1}{\varepsilon^{1/2}}\right) \end{aligned}$$

$$\begin{aligned} \text{and } x^{-1} \int_0^{\varepsilon \left(\frac{x}{6}\right)^{1/3}} s^5 e^{is^3} ds &= x^{-1} \int_0^{\varepsilon \left(\frac{x}{6}\right)^{1/3}} \underbrace{\frac{s^3}{3i}}_f \underbrace{3is^2 e^{is^3}}_{g'} ds \\ &= x^{-1} \left[\underbrace{\frac{s^3}{3i}}_f \underbrace{e^{is^3}}_g \Big|_0^{\varepsilon \left(\frac{x}{6}\right)^{1/3}} - \int_0^{\varepsilon \left(\frac{x}{6}\right)^{1/3}} \underbrace{\frac{s^2}{i}}_{f'} \underbrace{e^{is^3}}_g ds \right] \\ &= O(\varepsilon^3) + O(1/x) \end{aligned}$$

Expansions self-consistent provided $x^{-1/3} \ll \varepsilon \ll x^{-1/5}$ and therefore (by hint)

$$\underline{\underline{J_3(x) \sim \left(\frac{2}{9}\right)^{1/3} \Gamma\left(\frac{1}{3}\right) e^{i\pi/6} x^{-1/3} \text{ as } x \rightarrow \infty}}$$

Q4

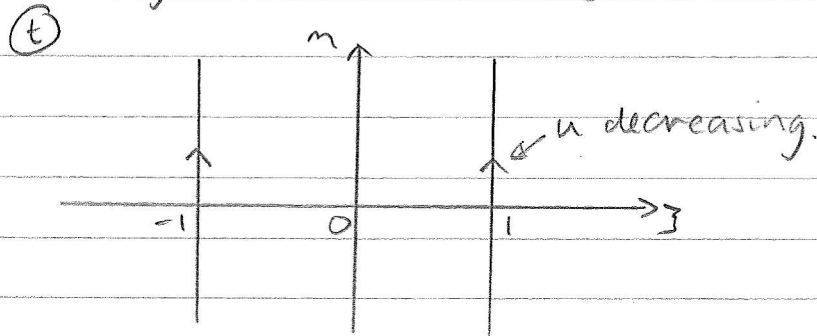
$$I(x) = \int_{-1}^1 f(t) e^{x\phi(t)} dt, \text{ where } f(t) = (1-t^2)^N, \phi(t) = it \quad (N \in \mathbb{N})$$

(a) Let $\phi(t) = u(\xi, \eta) + iv(\xi, \eta), t = \xi + i\eta$

$$\Rightarrow u = -\eta, v = \xi$$

On a steepest descent contour (SDC), the phase v is constant and $-\nabla u = -(u_\xi, u_\eta) = (0, 1)$ is a tangent pointing in the direction in which u decreases most rapidly.

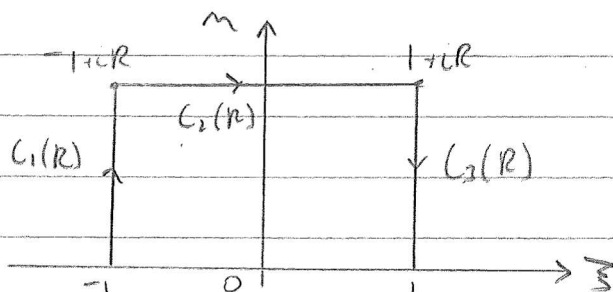
Thus, SDC through $t = \pm 1$ are $\xi = \pm 1$:



(b) Integrand a holomorphic function of t on \mathbb{C} , so by deformation theorem

$$I(x) = \left(\int_{C_1(R)} + \int_{C_2(R)} + \int_{C_3(R)} \right) f(t) e^{x\phi(t)} dt$$

where $C_i(R)$ are as shown:



Since $|f(t) e^{x\phi(t)}| = |f(t)| e^{-xR} \rightarrow 0$ as $R \rightarrow \infty$ on $C_2(R)$, let $R \rightarrow \infty$ to obtain

$$I(x) = \left(\int_{C_1(\infty)} + \int_{C_3(\infty)} \right) f(t) e^{x\phi(t)} dt = \left(\int_{-1}^{-1+i\infty} - \int_{1-i\infty}^{1+i\infty} \right) f(t) e^{x\phi(t)} dt$$

i.e. $I(x) = I_-(x) - I_+(x)$, where $I_+(x) = \int_{\pm 1}^{\pm 1+i\infty} f(t) e^{\alpha\phi(t)} dt$

(c) $I_+(x) = ie^{\pm i\alpha} \int_0^\infty f(\pm 1+is) e^{-\alpha s} ds \quad (t = \pm 1+is, s > 0)$

where $f(\pm 1+is) = [1 - (\pm 1+is)^2]^N = \begin{cases} (-is)^N (2+is)^N, & \oplus \\ (2+is)^N (is)^N, & \ominus \end{cases}$

Laplace's method:

$$I_+(x) = ie^{\pm i\alpha} \left(\int_0^\epsilon + \int_\epsilon^\infty \right) f(\pm 1+is) e^{-\alpha s} ds, \quad \epsilon \ll 1$$

\oplus first:

$$\int_0^\epsilon f(1+is) e^{-\alpha s} ds = \int_0^\epsilon (-is)^N (2+is)^N e^{-\alpha s} ds$$

$$= (-2i)^N \int_0^\epsilon (s^N + O(s^{N+1})) e^{-\alpha s} ds \quad (\because \epsilon \ll 1)$$

$$= \frac{(-2i)^N}{\alpha^{N+1}} \left[\int_0^{\epsilon\alpha} u^N e^{-u} du + O(\alpha^{-1}) \right] \quad (s = \frac{u}{\alpha})$$

$$= \frac{(-2i)^N}{\alpha^{N+1}} \left[\frac{\int_0^\infty u^N e^{-u} du}{\Gamma(N+1) = N!} + O(\alpha^{-1}) \right] \quad (\epsilon\alpha \gg 1)$$

$$= \frac{(-2i)^N N!}{\alpha^{N+1}} + O\left(\frac{1}{\alpha^{N+1}}\right)$$

$$\int_\epsilon^\infty f(1+is) e^{-\alpha s} ds = O(e^{-\epsilon\alpha})$$

Expansions self-consistent provided $\alpha^{-1} \ll \epsilon \ll 1$, so

$$I_+(x) \sim \frac{ie^{i\alpha} (-2i)^N N!}{\alpha^{N+1}} \text{ as } \alpha \rightarrow \infty$$

Similarly, $I_-(x) \sim \frac{ie^{-i\alpha} (2i)^N N!}{\alpha^{N+1}} \text{ as } \alpha \rightarrow \infty$.

$$\Rightarrow I(x) \sim \frac{2^N i^{N+1} N!}{\alpha^{N+1}} (e^{-i\alpha} - (-1)^N e^{i\alpha}) \text{ as } \alpha \rightarrow \infty$$

NB: real as required by symmetry of integrand!

Q5

(9)

$$\operatorname{erf}(z) = \frac{2}{\sqrt{\pi}} \int_0^z e^{-s^2} ds = \frac{2r}{\sqrt{\pi}} \int_0^{e^{i\theta}} e^{r^2 \phi(t)} dt$$

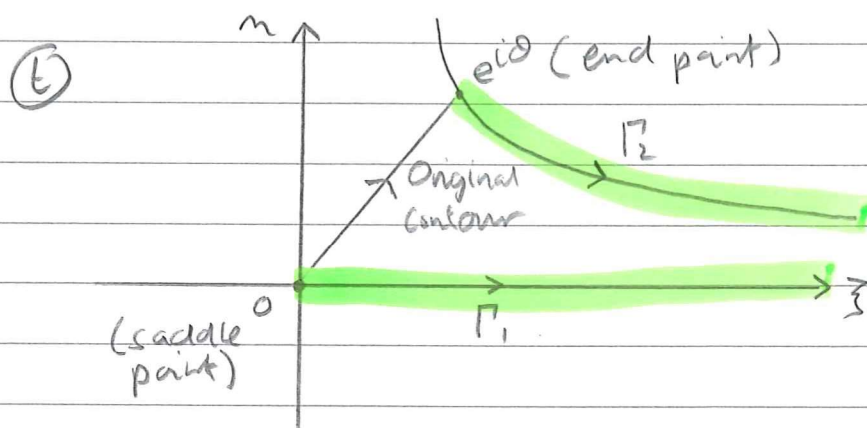
where $z = re^{i\theta}$, $s = rt$ and $\phi(t) = -t^2$.

Let $\phi(t) = u(\xi, \eta) + iv(\xi, \eta)$, $t = \xi + i\eta$

$$\Rightarrow u = \eta^2 - \xi^2, \quad v = -2\xi\eta, \quad \nabla u = 2(\xi, -\eta)$$

\Rightarrow • SDC through $t = 0$ is $\eta = 0$ (with $v = 0$)

• SDC through $t = e^{i\theta}$ ($0 < \theta < \pi/2$) is $2\xi\eta = 2\cos\theta\sin\theta = \sin 2\theta$
 $\xi > 0, \eta > 0$ (with $v = -\sin 2\theta$).



Deformation theorem $\Rightarrow \operatorname{erf}(z) = \left(\int_{\Gamma_1} - \int_{\Gamma_2} \right) \frac{2r}{\sqrt{\pi}} e^{r^2 \phi(t)} dt$

$$I_1(r) = \frac{2r}{\sqrt{\pi}} \int_0^\infty e^{-r^2 \xi^2} d\xi = \frac{2}{\sqrt{\pi}} \int_0^\infty e^{-\xi^2} d\xi = 1$$

$\xi = \frac{\hat{\xi}}{r}$

$$I_2(r, \theta) = \frac{2r}{\sqrt{\pi}} \int_{\cos\theta}^\infty e^{r^2(u(\xi, \eta(\xi)) - i\sin 2\theta)} (1 + i\eta'(\xi)) d\xi$$

because $\Gamma_2 = \{ \xi + i\eta(\xi) : \xi > \cos\theta \}$, where $\eta(\xi) = \frac{\sin 2\theta}{2\xi}$

$$\Rightarrow I_2(r, \theta) = \frac{2r}{\sqrt{\pi}} e^{-r^2 i \sin 2\theta} \int_{\cos\theta}^\infty F(\xi) e^{r^2 \Phi(\xi)} d\xi$$

where $F(\xi) = 1 + i\eta'(\xi) = 1 - \frac{i \sin 2\theta}{2\xi^2}$, $\Phi(\xi) = u(\xi, \eta(\xi)) = \frac{\sin^2 2\theta}{4\xi^2} - \xi^2$

Γ_2 a SDC $\Rightarrow \Phi(\zeta) = u(\zeta, v(\zeta))$ decreasing with ζ on Γ_2

Hence, Laplace's method gives

$$I_2(r, \theta) \sim -\frac{2r}{\sqrt{\pi}} e^{-r^2 i \sin 2\theta} \frac{F(\cos \theta) e^{r^2 \Phi(\cos \theta)}}{r^2 \Phi'(\cos \theta)} \quad \text{as } r \rightarrow \infty$$

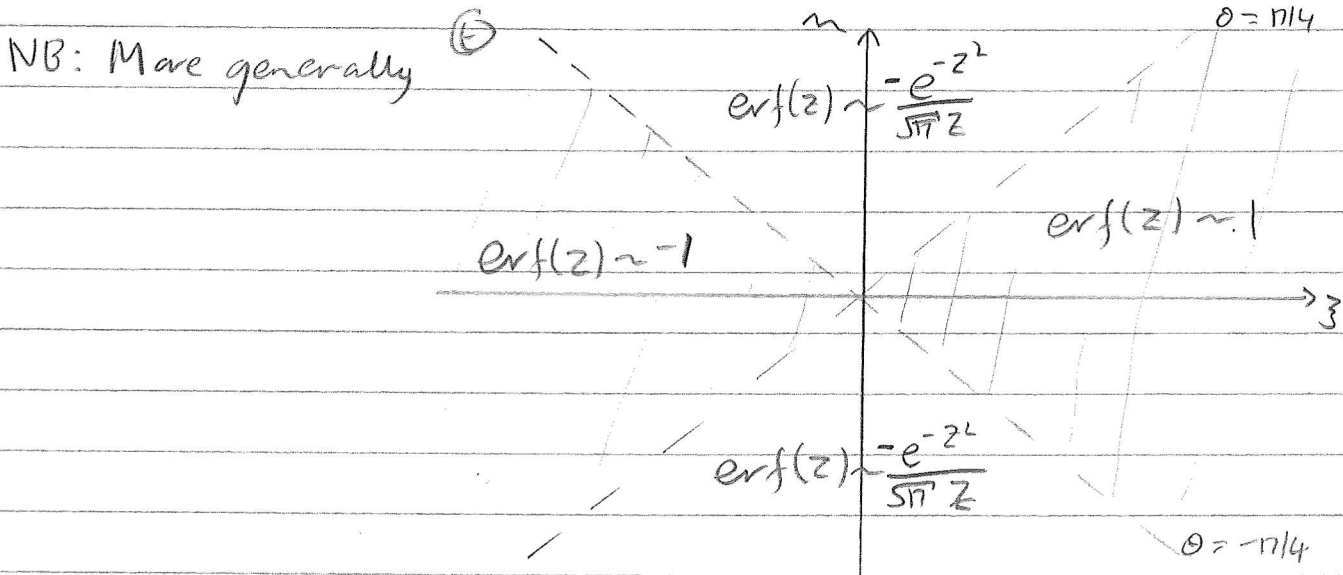
where $F(\cos \theta) = \frac{e^{-i\theta}}{\cos \theta}$, $\Phi(\cos \theta) = -\cos 2\theta$, $\Phi'(\cos \theta) = \frac{-2}{\cos \theta}$

$$\Rightarrow I_2(r, \theta) \sim +\frac{1}{\sqrt{\pi} r e^{i\theta}} e^{-r^2 e^{2i\theta}} \quad \text{as } r \rightarrow \infty.$$

Thus, $I_1 \sim 1$ and $I_2 \sim +\frac{1}{\sqrt{\pi} z} e^{-z^2}$ as $r = |z| \rightarrow \infty$ for $0 < \theta = \arg(z) < \pi/2$

$$|I_2| \sim \frac{1}{r} e^{-r^2 \cos 2\theta} \begin{cases} \ll 1 & \text{for } 0 < \theta \leq \pi/4 \\ \gg 1 & \text{for } \pi/4 < \theta < \pi/2 \end{cases}$$

$$\Rightarrow \text{erf}(z) \sim \begin{cases} 1 & \text{for } 0 < \theta \leq \pi/4 \\ -\frac{1}{\sqrt{\pi} z} e^{-z^2} & \text{for } \pi/4 < \theta < \pi/2 \end{cases} \quad \text{as } |z| \rightarrow \infty$$



- Different asymptotic expansion in different sectors called Stokes phenomenon
- While e^{-z^2} is entire, it has an essential singularity at ∞
- $\theta = \pm \pi/2$ are Stokes lines (across which topology of SDC changes).
- $|\theta| = \pi/4, 5\pi/4$ are anti-Stokes lines (across which dominance of end point and saddle point changes).

Q6

$$I(\varepsilon) = \int_0^1 \frac{f(x)}{x+\varepsilon} dx$$

$$= \left(\int_0^{\delta} + \int_{\delta}^1 \right) \frac{f(x)}{x+\varepsilon} dx, \text{ where } \varepsilon \ll \delta \ll 1.$$

$$I_1 = \int_0^{\delta/\varepsilon} \frac{f(\varepsilon y) dy}{y+1} \quad (x = \varepsilon y)$$

$$= \int_0^{\delta/\varepsilon} \frac{1}{y+1} (f(0) + \varepsilon y f'(0) + O(\varepsilon^2)) dy \quad (\because \varepsilon y \ll \delta \ll 1)$$

$$= f(0) \ln(y+1) \Big|_0^{\delta/\varepsilon} + O(\delta)$$

$$= f(0) \ln\left(1 + \frac{\delta}{\varepsilon}\right) + O(\delta)$$

$$= f(0) \ln\left(\frac{\delta}{\varepsilon}\right) + f(0) \ln\left(1 + \frac{\varepsilon}{\delta}\right) + O(\delta)$$

$$= -f(0) \ln \varepsilon + f(0) \ln \delta + O\left(\delta, \frac{\varepsilon}{\delta}\right)$$

$$I_2 = \int_{\delta}^1 \frac{f(x) dx}{x+\varepsilon}$$

$$= \int_{\delta}^1 \frac{f(x)}{x(1+\frac{\varepsilon}{x})} dx$$

$$= \int_{\delta}^1 \frac{f(x)}{x} \left(1 - \frac{\varepsilon}{x} + O(\varepsilon^2)\right) dx \quad (\because \frac{\varepsilon}{x} < \frac{\varepsilon}{\delta} \ll 1)$$

$$\sim \int_{\delta}^1 \frac{f(x) - f(0)}{x} dx + \int_{\delta}^1 \frac{f(0)}{x} dx + \dots$$

$$\sim -f(0) \ln \delta + \int_{\delta}^1 \frac{f(x) - f(0)}{x} dx + \dots$$

$$\Rightarrow I \sim -f(0) \ln \varepsilon + f(0) \ln \delta - f(0) \ln \delta + \int_{\delta}^1 \frac{f(x) - f(0)}{x} dx + \dots$$

$$\sim -f(0) \ln \varepsilon + \int_0^1 \frac{f(x) - f(0)}{x} dx + \dots \quad \text{as } \varepsilon \rightarrow 0$$

Q1

$\int_0^{\pi/2} e^{ix \cos t} dt$: stationary phase for first term;
steepest descents for more.

$\int_0^1 \ln t e^{ixt} dt$: steepest descents (IBP & stationary phase don't work)

$\int_0^{\infty} t^{-1/2} e^{-t} dt$: $\int_0^{\infty} - \int_x^{\infty}$ and IBP on \int_x^{∞} .

$\int_0^{\pi/2} e^{-x \sin^2 t} dt$: Laplace's method.

$\int_0^1 e^{ixt} e^{-1/t} dt$: steepest descents having set $s = e^{-1/t}$.

$\int_0^{10} \frac{e^{-xt}}{1+t} dt$: Taylor expand integrand & integrate term-by-term

$\int_0^{\pi/2} \frac{dt}{\cos^2 t + x \sin^2 t}$: $\int_0^{\pi/2} = \int_0^{\pi/2 - \delta} + \int_{\pi/2 - \delta}^{\pi/2}$, where $x \ll \delta \ll 1$.

$\int_0^1 \frac{\sin(xt)}{t} dt$: Taylor expand integrand & integrate term-by-term.

$\int_x^{\infty} t^{a-1} e^{-t} dt$: $\int_x^{\infty} = \int_0^{\infty} - \int_0^x$ and \int_0^x on \int_0^x for $\operatorname{Re}(a) > 0$;

much more tricky for $\operatorname{Re}(a) \leq 0$!

$\int_0^1 \frac{\ln t}{x+t} dt$: $\int_0^1 = \int_0^{\delta} + \int_{\delta}^1$, where $x \ll \delta \ll 1$.