

## Klick theorem

- One-point correlator in free theory:

$$\langle \phi(x_1) \rangle_0 = \frac{1}{2} \int d^d x'' \Delta(x_1 - x'') J(x'') + \frac{1}{2} \int d^d x' \Delta(x' - x_1) J(x') \Big|_{J=0} = 0$$

- Two-point correlator

$$\langle \phi(x_1) \phi(x_2) \rangle_0 = \frac{1}{2} \Delta(x_1 - x_2) + \frac{1}{2} \Delta(x_2 - x_1) = \Delta(x_1 - x_2)$$

- Higher-point functions:

- if  $N$  is odd :  $\langle \phi(x_1) \dots \phi(x_N) \rangle = 0$

based on the symmetry of the Lagrangian  $\phi \leftrightarrow -\phi$

- if  $N$  is even:

$$N=4: \langle \phi(x_1) \phi(x_2) \phi(x_3) \phi(x_4) \rangle_0 = \Delta(x_1 - x_2) \Delta(x_3 - x_4) \\ + \Delta(x_1 - x_3) \Delta(x_2 - x_4) \\ + \Delta(x_1 - x_4) \Delta(x_2 - x_3)$$

General even  $N$ :

$$\langle \phi(x_1) \dots \phi(x_N) \rangle_0 = \sum \Delta(x_{j_1} - x_{j'_1}) \dots \Delta(x_{j_{N/2}} - x_{j'_{N/2}})$$

sum runs over all distinct ways of grouping the set  $\{1, \dots, N\}$  in pairs

## Klick theorem

- Graphical notation:

$$x_1 \text{ --- } x_2 = \Delta(x_1 - x_2)$$

Example:  $N=4$

$$\langle \phi(x_1) \phi(x_2) \phi(x_3) \phi(x_4) \rangle_0 = \begin{array}{c} x_1 \text{ --- } x_2 \\ x_3 \text{ --- } x_4 \end{array} + \begin{array}{c} x_1 \text{ --- } x_3 \\ x_2 \text{ --- } x_4 \end{array} + \begin{array}{c} x_1 \text{ --- } x_4 \\ x_2 \text{ --- } x_3 \end{array}$$

# Feynman diagrams

- For interacting theories the Lagrangian splits

$$L = L_0 + L_I$$

↑ only quadratic in fields

- Correlation functions

$$G^{(n)}(y_1, \dots, y_n) = \frac{\int \mathcal{D}\phi \phi(y_1) \dots \phi(y_n) e^{-S_0 - \int L_I(x) d^D x}}{\int \mathcal{D}\phi e^{-S_0 - \int L_I(x) d^D x}}$$

We expand the numerator and denominator in powers of  $L_I$ :

$$e^{-\int L_I(x) d^D x} = \sum_{n=0}^{\infty} \frac{(-1)^n}{n!} \int d^D x_1 \dots d^D x_n \bullet L_I(x_1) \dots L_I(x_n)$$

All terms in the numerator are then of the form

$$\langle \phi(y_1) \dots \phi(y_n) L_I(x_1) \dots L_I(x_n) \rangle_0$$

→ these are correlators in free theory

Since  $L_I$  is a polynomial in  $\phi$  then it can be done by Wick's theorem.

- Example:  $\phi^4 \rightarrow L_I(x) = \frac{\lambda}{4!} \phi(x)^4$

Let us move fields at  $x$  a bit apart  $\phi(x)^4 \rightarrow \prod_{j=1}^4 \phi(x_j)$

$$G^{(2)}(y_1, y_2)_{\text{NUM}} = \begin{array}{c} y_1 \text{ --- } y_2 \\ + \left(-\frac{\lambda}{4!}\right) \begin{array}{c} y_1 \text{ --- } \text{---} \text{---} y_2 \\ \quad \quad \quad \cup \\ \quad \quad \quad x \end{array} \\ + \left(-\frac{\lambda}{4!}\right) \begin{array}{c} y_1 \text{ --- } \text{---} \text{---} y_2 \\ \quad \quad \quad \cup \\ \quad \quad \quad \cup \\ \quad \quad \quad x \end{array} \end{array}$$

4 · 3 = 12 3

$$= \Delta(y_1 - y_2) - 12 \frac{\lambda}{4!} \int d^D x \Delta(y_1 - x) \Delta(x - x) \Delta(x - y_2) - 3 \frac{\lambda}{4!} \Delta(y_1 - y_2) \int d^D x \Delta(x - x)^2$$

$$G^{(2)}(y_1, y_2)_{\text{NUM}} = \text{---} + \text{---} + \text{---}$$

Now, the denominator:

$$G^{(2)}(y_1, y_2)_{\text{DEN}} = 1 + 8$$

$$= 1 - \frac{3\lambda}{4!} \int \Delta(x-x)^2 d^D x$$

• The integral  $\int \Delta(x-x)^2 d^D x = \Delta^2(0) \int d^D x$  diverges

$$\int d^D x = V \cdot T \rightarrow \infty$$

This contribution is cancelled by the denominator.

We keep  $V \cdot T$  finite, divide and only then take  $V \cdot T$  to infinity.

→ diagrams containing pieces not connected to the external points (bubbles) can be ignored

• Second-order two-point function

$$\lambda^2: \text{---} + \text{---} + \text{---}$$

The last one:

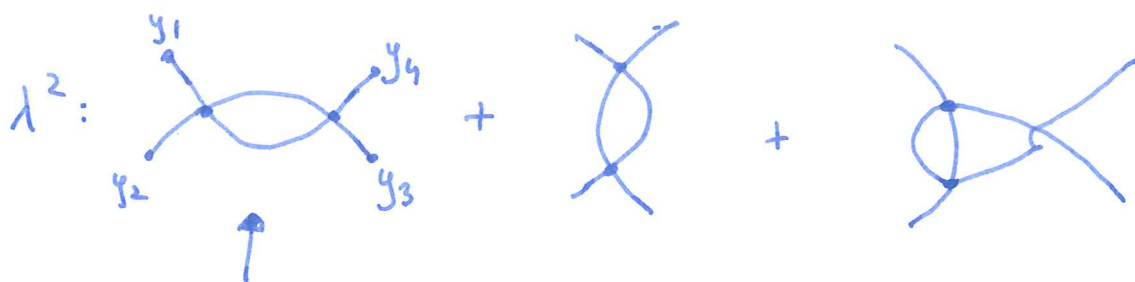


$$2 \cdot 4 \cdot 4 \cdot 3! \cdot \frac{1}{2!} \left(-\frac{\lambda}{4!}\right)^2 \int \Delta(y_1 - x_1) \Delta(x_1 - x_2)^3 \Delta(x_2 - y_2) d^D x_1 d^D x_2$$

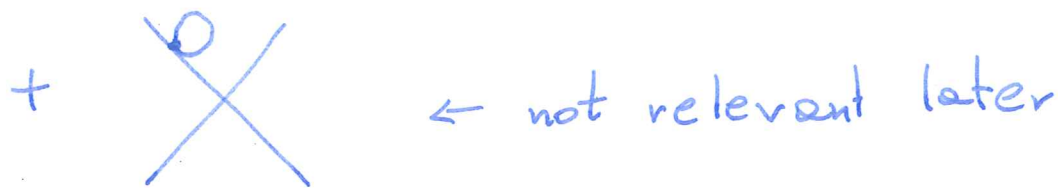
$\uparrow$                        $\uparrow$                        $\uparrow$   
 $x_1 \leftrightarrow x_2$      $x$  with  $y$      $x_1$  with  $x_2$

• Four-point correlation function (only connected contributions)

$$\lambda: \quad \begin{array}{c} \diagup \\ \diagdown \end{array} = \begin{array}{c} \diagdown \\ \diagup \end{array} \\ = 4! \frac{-\lambda}{4!} \int \prod_{j=1}^4 \Delta(y_j - x) d^D x$$



$$\frac{(-1)^2}{2} \int \Delta(y_1 - x_1) \Delta(y_2 - x_1) \Delta(x_1 - x_2)^2 \Delta(x_2 - y_3) \Delta(x_2 - y_4) d^D x_1 d^D x_2$$



## Feynman rules in position space

1. Draw all topologically distinct connected diagrams with  $N$  external lines and each internal vertex attached to 4 lines.
  2. To each line associate a factor  $\Delta(x'-x'')$
  3. To each internal vertex associate a factor  $-1$
  4. Integrate over internal vertices  $\prod_j \int d^D x_j$
  5. Multiply by the symmetry factor  $\frac{1}{\text{integer}}$
- Only the last point might be problematic.
    - generic diagram symmetry factor is 1.
    - non-generic diagrams:

$$\text{---} \bigcirc \text{---} \rightarrow \frac{1}{2} \quad (\text{twist the bubble})$$

$$\text{---} \bigcirc \bigcirc \text{---} \rightarrow \left(\frac{1}{2!}\right)^2$$

$$\text{---} \bigcirc \text{---} \rightarrow \frac{1}{3!} \quad (\text{permutation of 3 internal lines})$$

In general: symmetry factor is the inverse of number of elements in the symmetry group of a given diagram.

## General formula for symmetry factors in $\phi^4$

Define:

$S$  - number of self-connections

$D$  - number of double connections

$T$  - number of triple connections


$Q$  - number of quadruple connections


$N_{IVP}$  - number of identical vertex permutations

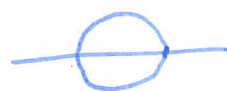
Then:

$$S_G = \frac{1}{2^{S+D} (3!)^T (4!)^Q N_{IVP}}$$

Example:

 :  $S=1, D=0, T=0, N_{IVP}=1$   
 $\Rightarrow S_G = \frac{1}{2}$

 :  $S=1, D=1, T=0, N_{IVP}=1$   
 $\Rightarrow S_G = \frac{1}{4}$

 :  $S=0, D=0, T=1, N_{IVP}=1$   
 $\Rightarrow S_G = \frac{1}{3!}$