## B5.6: Nonlinear Systems-Sheet 1 (solutions)

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**Q1** To determine the subspaces of the matrix

$$A = \begin{bmatrix} 0 & 2 & 0 \\ -2 & 0 & 0 \\ 2 & 0 & 6 \end{bmatrix},\tag{1}$$

we find its eigenvalues and eigenvectors, which are:

$$\lambda_1 = 6 \quad \mathbf{v}^{(1)} = (0, 0, 1)^T,$$
  

$$\lambda_2 = 2i \quad \mathbf{v}^{(2)} = i(0, 10, -1)^T + (10, 0, -3)^T,$$
  

$$\lambda_3 = -2i \quad \mathbf{v}^{(3)} = -i(0, 10, -1)^T + (10, 0, -3)^T.$$

Therefore, the stable  $E^s$ , unstable  $E^u$ , and center  $E^c$  subspaces of (1) are:

$$E^{s} = \emptyset,$$
  

$$E^{u} = \operatorname{span}\{(0,0,1)^{T}\},$$
  

$$E^{c} = \operatorname{span}\{(0,10,-1)^{T}, (10,0,-3)^{T}\};$$

that is, the empty set, the z-axis, and the plane specified by z = -3x/10 - y/10.

## Q2 (invariant set) For the system:

$$\dot{x} = -x, \tag{2}$$

$$\dot{y} = -y + x^2, \tag{3}$$

$$\dot{z} = z + x^2 \tag{4}$$

we define a quantity  $Q(t) = z + x^2/3$  and find, upon differentiating:

$$\dot{Q} = z + x^2/3,$$

which vanishes only if  $z = -x^2/3$ . Therefore,  $z = -x^2/3$  is an invariant set of the system. The other invariant set is the fixed point at (0, 0, 0), which, by linearization, is a saddle node because it has eigenvalues and eigenvectors given by:

$$\lambda_1 = 1 \quad \mathbf{v}^{(1)} = (0, 0, 1)^T, \lambda_2 = -1 \quad \mathbf{v}^{(2)} = (1, 0, 0)^T, \lambda_3 = -1 \quad \mathbf{v}^{(3)} = (0, 1, 0)^T.$$

The phase portrait of the system is shown in Fig. 2.



Figure 1: Phase portrait of (2)-(4) on the planes x = y = 0 and  $z = -x^2/3$ .

Q3 (attracting set) We consider the system:

$$\dot{x} = -y + x(1 - z^2 - x^2 - y^2),$$
 (5)

$$\dot{y} = x + y(1 - z^2 - x^2 - y^2),$$
 (6)

$$\dot{z} = 0 \tag{7}$$

We introduce the radius in cylindrical coordinates  $r^2 = x^2 + y^2$  and differentiate with respect to time:

$$r\dot{r} = x\dot{x} + y\dot{y}, \quad \Longrightarrow \quad \dot{r} = r(1 - r^2 - z^2),$$

which suggests r = 0 (i.e z-axis) is the  $\alpha$ -limit set (asymptotically unstable in this case) whilst  $r = \sqrt{1 - z^2}$  is the  $\omega$ -limit set (asymptotically stable in this case). Each circle is a level set in  $z \in [-1, 1]$ , which together form a unit sphere that is an attracting set (i.e. the orbits within a shell of radius  $\rho = \sqrt{1 - r^2 - z^2} \in [1 - \varepsilon, 1 + \varepsilon] \quad \forall \varepsilon > 0$  will tend to the unit sphere as  $t \to \infty$ ). Defining the unit sphere as  $S^1$ , the domain of attraction would be  $D(S^1) = \mathbb{R}^2 \times [-1, 1] \setminus \{0\}$ , given that  $\dot{z} = 0$ . In turn, the domain of attraction of the part of z-axis with |z| > 1 is given by  $\mathbb{R}^2 \setminus D(S^1)$ .

We note that there does not exist a dense orbit  $\Gamma$  on the unit sphere such that,  $\forall \Gamma$  also on the unit sphere, there are points  $x \in \Gamma$  and  $\tilde{x} \in \Gamma$  satisfying  $|x - \tilde{x}| < \varepsilon \quad \forall \varepsilon > 0$ . Orbits on the unit sphere are confined to circles at a particular z and will, therefore, never come arbitrarily close together.

Q4 (attracting set) We study the system:

$$\dot{r} = r(1-r), \quad \dot{\theta} = \sin^2 \frac{\theta}{2},$$
(8)

We note that fixed points occur at  $r = \{0, 1\}$  and  $\theta = 2\pi n$  for  $n \in \mathbb{Z} \cup \{0\}$ , which correspond to two fixed points in Cartesian coordinates, (x, y) = (0, 0) and (1, 0). The radius r = 0, or the fixed point (0, 0), is the  $\alpha$ -limit set whilst r = 1 is asymptotically stable. Due to (8), the trajectories proceed around the unit circle in an anti-clockwise direction and asymptotically approach (1, 0). Therefore, (1, 0) is the  $\omega$ -limit point in the system for all initial conditions except the fixed point at the origin (i.e.  $x_0 \neq 0$ ).

Despite this, however,  $\theta = 0$  is unstable according to (8). Perturbing the state by  $\theta = \varepsilon > 0$  for  $\varepsilon \ll 1$  will result in the trajectory proceeding anti-clockwise around the unit circle until it reaches (1,0) again.

The unit circle is an attracting set because it is an invariant set and all trajectories that are within the radius  $r \in [1-\varepsilon, 1+\varepsilon]$ , for  $\varepsilon > 0$ , will end up within the set. The fixed point (1,0) is also an attracting set since the same neighborhood as above satisfies the definition. The domain of attraction for both the unit circle and (1, 0) is  $D(r = 1) = D((1, 0)) = \mathbb{R}^2 \setminus \{0\}$ .

**Q5\*** (the variational equation and its adjoint): (i) We are asked to show that  $u(t) = \dot{x}(t)$  is a solution of the linear equation

$$\dot{y} = Df(\bar{x})y,\tag{9}$$

where  $\bar{x}(t)$  is a particular solution of  $\dot{x} = f(x)$ . We note that:

$$\dot{\bar{x}} = f(\bar{x}),$$
  
 $\implies \ddot{\bar{x}} = Df(\bar{x})\dot{\bar{x}},$ 

which is (9) for  $y = u = \dot{x}$ , as required. The vector  $\dot{x}(t)$  is tangential to the flow by construction and is analogous to the velocity of a particle traveling on a curve whose position is given by  $\bar{x}(t)$ .

(*ii*) For planar flows, we require two linearly independent solutions to (9). From (*i*) we know that one of this solutions is  $u(t) = \dot{\bar{x}}(t)$ . To determine the other solution, which we will call v(t), we suppose it has the general form:

$$v(t) = g(t)\tau(t) + h(t)n(t), \qquad (10)$$

where  $\tau(t) = \dot{x}(t)$  is a vector tangential to the flow traced out by  $\bar{x}(t)$  and  $n(t) = \tau(t)^{\perp}$  is the corresponding normal. Note that, for a prescribed  $\bar{x}(t)$ , both  $\tau(t)$  and n(t) are known and they are orthogonal to each other. Given that v(t) is a solution to linear system (9), we differentiate (10) with respect to t to obtain:

$$\dot{g}\tau + g\dot{\tau} + \dot{h}n + h\dot{n} = Df(\bar{x})g\tau + Df(\bar{x})hn, \implies \dot{g}\tau + \dot{h}n + h\dot{n} = Df(\bar{x})hn,$$
(11)

where the latter step results from the definition of  $\tau(t) = \dot{x}(t)$  and (9). Multiplying (11) by *n*, we use orthogonality to obtain:

$$\dot{h} = \frac{h}{|n|^2} (nDf(\bar{x})n - \dot{n}n) = K_1(t)h, \implies h(t) = c_1 \exp\left(\int K_1(t) \, dt\right),$$
(12)

for  $c_1 \in \mathbb{R}$ . Note that  $K_1(t)$  is completely known for a given  $\bar{x}(t)$ . Taking (11) and multiplying through by  $\tau$ , we now obtain:

$$\dot{h} = \frac{h}{|\tau|^2} (\tau D f(\bar{x}) n - \dot{n} \cdot \tau) = K_2(t) h, \implies h(t) = \int K_2 h \, dt + c_2, \quad (13)$$

for  $c_2 \in \mathbb{R}$ . Similarly,  $K_2(t)$  is completely known. We obtain the explicit form of v(t) by combining (10), (12) and (13) above, so that the complete solution of (9) is

$$y(t) = a_1 u(t) + a_2 v(t)$$

for  $a_1, a_2 \in \mathbb{R}$ .