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GR II Part C Math
(also MMathPhys, MScMTP)

Outline (roughly!)

- * Mathematical techniques from differential geometry
- * Field equations with matter
- * Linearized GR
 - gravitational field of an isolated body at large distances
 - gravitational waves
- * Black holes
 - (review of) Schwarzschild solution
 - Penrose diagrams
 - Kerr solution
 - time permitting: charged black holes

[For MMathPhys / MScMTP: watch for topics for oral presentations]

Chapter 1: Mathematical background

1.1 Mathematical model of GR

In Einstein's GR, the model for spacetime (space of events) is a 4 dim differentiable manifold M with

- a metric g with signature $(-+++)$

↳ symmetric (0,2) tensor which is differentiable and non-degenerate (ie $\det g \neq 0$ or equivalently $g_{ab} X^a Y^b = 0 \quad \forall Y \Rightarrow X=0$)

(g defines an inner product on $TM \cong \mathbb{R}^4$;

line element = infinitesimal distance between neighboring events: $ds^2 = g_{ab}(x) dx^a dx^b$)

AND • a connection ∇ which is metric ($\nabla g = 0$) and torsion free

$M \rightsquigarrow$ Riemannian geometry

Why do we want M to be a differentiable manifold? (3)

(A) Informal discussion of manifolds

↖ set of points + neighbourhood basis
manifold: (topological) space which locally "looks like" (homeomorphic to) \mathbb{R}^n

That is:

M can be covered by a collection of coordinate charts (U_i, X_i^a) $i=1 \rightarrow N$ (atlas) where

$U_i \subset M$ open sets in M
 ↖ "coordinate patch"
 "local coordinate neighborhood"

"coordinates" → $X_i^a : \begin{array}{ccc} U_i & \longrightarrow & \mathbb{R}^n \\ \downarrow & & \\ p & \longmapsto & (x_i^1, x_i^2, \dots, x_i^n) \end{array}$

with the requirement that on

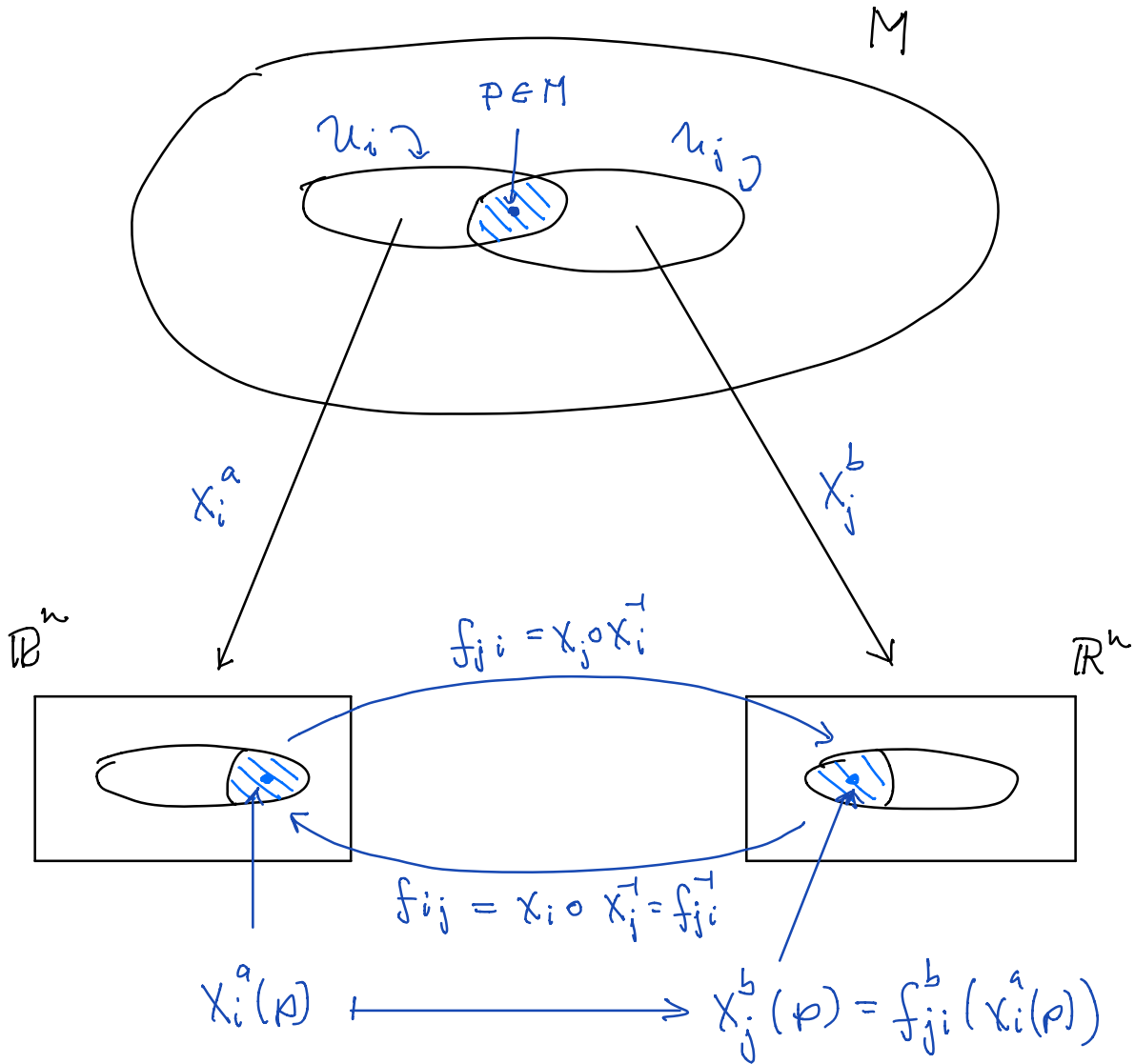
$$U_{ij} = U_i \cap U_j$$

we have transition functions

$$X_i^a = f_{ij}^a(X_j^b)$$

which are continuous

4



f, f^{-1} continuous, defined only on the overlap $U_i \cap U_j$

composition of f_{ji} and x_i same as x_j

$$f_{ji} \circ x_i = x_j$$

5

differentiable manifold :

transition functions are smooth
(ie differentiable to all orders)

For GR : we want a differentiable
manifold because of

Einstein's equivalence principle

ie special relativity holds at small
distances

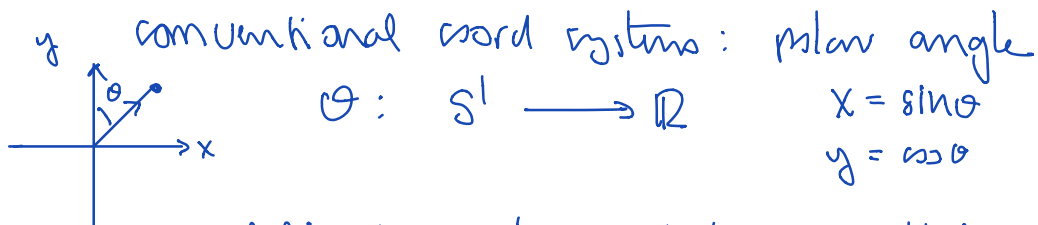
- coordinate patches $\sim \mathbb{R}^4$
- transition functions (change of coordinates) are smooth functions

Examples (see GR1)

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1) \mathbb{R}^n one coordinate patch

2) S^1 $(x, y) \in \mathbb{Q}^2$ st $x^2 + y^2 = 1$

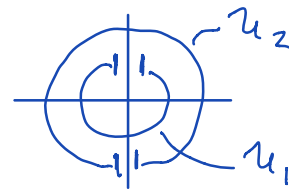


θ defined up to an integral multiple of 2π

$$\theta \simeq \theta + 2\pi n \quad n \in \mathbb{Z}$$

but $0 \leq \theta \leq 2\pi$ is not an open interval

Need two patches
to cover the circle



3) \vee , ∞ are not manifolds

(7)

4) Sphere $S^2: (x, y, z) \in \mathbb{R}^3$ with $x^2 + y^2 + z^2 = 1$

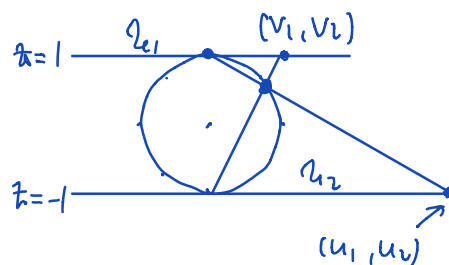
$\mathcal{U}_1: S^2 \setminus \{\text{north pole}\}$

$\mathcal{U}_2: S^2 \setminus \{\text{south pole}\}$

use stereographic projection

$$\mathcal{U}_1: p \mapsto (u_1, u_2) = \left(\frac{2x}{1-z}, \frac{2y}{1-z} \right)$$

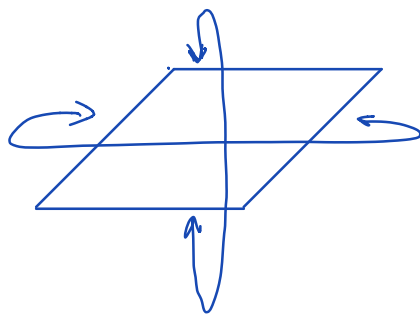
$$\mathcal{U}_2: p \mapsto (v_1, v_2) = \left(\frac{2x}{1+z}, \frac{2y}{1+z} \right)$$



$$\mathcal{U}_1 \cap \mathcal{U}_2 \quad -1 < z < 1$$

$$v_i = \frac{4u_i}{u_1^2 + u_2^2} \text{ is } C^\infty$$

5) T^2



4 patches is enough.

(B)

(8)

Tensors geometric objects on M
defined naturally by the manifold structure

In GR: we write equations in
terms of tensors (eg g , the metric,
 R , the curvature of the metric, etc)
(laws of nature in terms of tensors!)

Principle of relativity \leftrightarrow general covariance

laws of nature are covariant i.e. independent of
the choice of a coordinate system.

(p,q) tensor $T^{a_1 \dots a_p}_{b_1 \dots b_q}$

defined by its transformation properties under
coordinate transformations

indices: book-keeping of transformation laws

under a coordinate transformation $x^a \rightarrow \tilde{x}^a(x)$

$$T^{a_1 \dots a_p}_{b_1 \dots b_q} = \left(\partial_{c_1} \tilde{x}^{a_1} \dots \partial_{c_p} \tilde{x}^{a_p} \right) \left(\tilde{\partial}_{b_1}^{d_1} \dots \tilde{\partial}_{b_q}^{d_q} \right) T^{c_1 \dots c_p}_{d_1 \dots d_q}$$

Notation: $\partial_a = \partial / \partial x^a$

- Examples:
- scalar $(0,0)$ $\tilde{\phi}(\tilde{x}) = \phi(x)$
 - vector $(1,0)$ (contravariant vector)
 - $(0,1)$ 1-form, covector, covariant vector
 - the metric is a $(0,2)$ symmetric tensor

Tensor operations (part 1)

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(operations on tensors or between tensors to get other tensors)

- 1) addition: $T + S$; T, S same type
- 2) contraction: $(p, q) \rightarrow (p-1, q-1)$
sum over one upper index and one lower index
- 3) raise and lower indices with the metric
 $(p, q) \rightarrow (p-1, q+1)$ "lowering"
 $(p, q) \rightarrow (p+1, q-1)$ "raising"

(Needs a metric g on M . for example
 $TM \cong (TM)^*$ is an isomorphism given by
 $V_a = g_{ab} u^b$ $u^a = g^{ab} V_b$)

4) exterior derivative of p -forms

• if a scalar: $\partial_a f$ is a 1-form (covector)

• let $V = V_a dx^a$ be a 1-form

want
derivation
which
is a
tensor

$\partial_a V_b$ is not a tensor

The same is true more generally:

if T is (p, q) -tensor, $\partial_a T$ is not a tensor

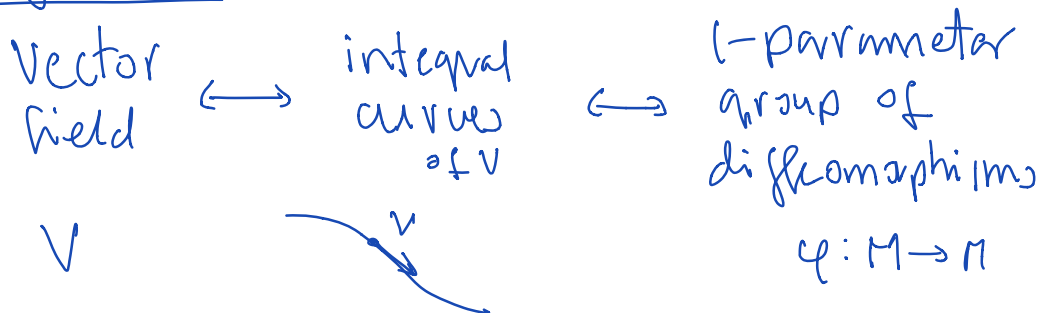
p -forms • However: if A is a p -form i.e. a totally antisymmetric $(0, p)$ tensor

exterior
derivative Then $dA = \frac{1}{p!} \partial_{[a} A_{a_1 \dots a_p]} dx^{b a_1 \dots a_p}$ is a $p+1$ form.

(Will discuss derivations $dx, [,], \nabla$ later)

(c) Vector fields, diffeomorphisms and integral curves

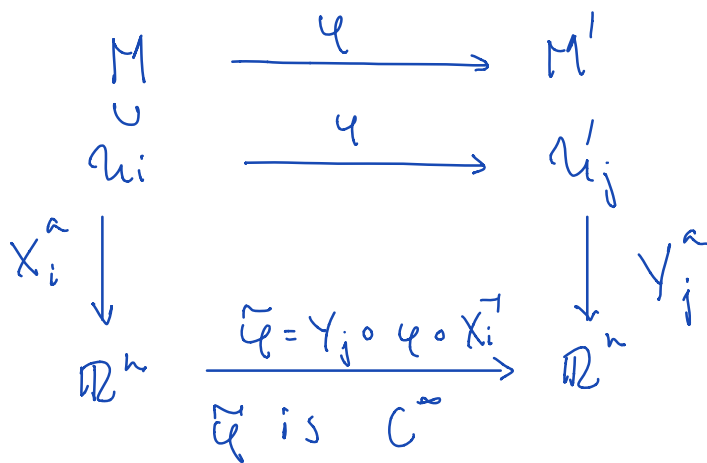
Key ideas:



A diffeomorphism between two differentiable manifolds M & M' is a C^∞ bijective map

$$\varphi: M \longrightarrow M'$$

which preserves the manifold structure



We think of two diffeomorphic manifolds as "equivalent"

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Examples

1) S^2 with coordinates (θ, ϕ)
 \rightarrow a rotation around the z -axis

2) T^2



3) $S^2 \rightarrow$ ellipse

$$(x, y, z) \leftrightarrow \left(\frac{x}{a}, \frac{y}{b}, \frac{z}{c} \right)$$
$$x^2 + y^2 + z^2 = 1$$

4) S^2 & T^2 are not diffeomorphic

(End of lecture 1)

lecture #2

(2)

A one-parameter group of diffeomorphisms of M is a differentiable map

$$\begin{aligned} \varphi: \mathbb{R} \times M &\longrightarrow M \\ (t, x^a) &\longmapsto \varphi^a(t, x) = \varphi_t^a(x) = x^a(t) \end{aligned}$$

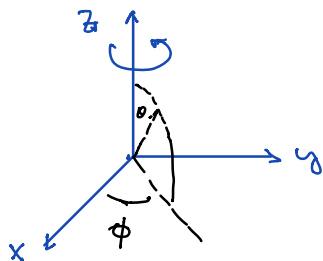
$t \in (a, b) \subset \mathbb{R}$

where $\varphi_t: M \longrightarrow M$ is a diffeomorphism with $\varphi_0 = \text{identity}$, and

$$\varphi_{s+t} = \varphi_s \circ \varphi_t \quad \text{group law}$$

Example: S^2 with coordinates (θ, ϕ)

$$\varphi_t(\theta, \phi) = (\theta, \phi + t) \quad \begin{array}{l} \text{rotation by } t \\ \text{around } z\text{-axis} \end{array}$$



orbits \longrightarrow circles
 $\theta = \text{constant}$

Remark: $\varphi^a(t, x) = x^a(t)$ defines a curve on M parametrized by t

Theorems: on the relation between vector fields and 1-parameter group of diffeomorphisms

$\varphi \rightarrow V$ (i) To each 1 parameter group of diffeomorphisms φ_t we can associate a vector field V by

$$V^a(x) = \frac{d}{dt} \varphi_t^a(x)$$

$V \rightarrow \varphi$ (ii) Conversely: let V be a smooth vector field on M . Then, there is a family of curves associated to $V|_p$ for each $p \in M$, such that one and only one curve passes through $p \in M$, and the tangent to this curve at p is V_p .

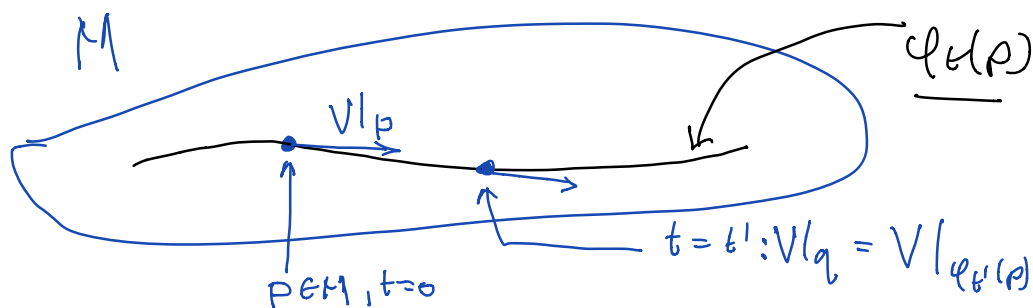
These curves are called integral curves of V

Proof of (i) what happens to $p \in M$
under the action of φ ?

(14)

$\varphi \rightarrow V$

Fix $p \in M$: $\varphi_t(p)$ is a unique
curve through p at $t=0$ ($\varphi_0(p) = p$)



φ_t traces a path on M (orbit of φ_t through p)

The components of the tangent vector V
to the curve at $\varphi_t(p)$ are given by

$$V^a(x(t)) = \frac{d}{dt} \varphi_t^a(x) \quad (\text{differentiate } \varphi)$$

in particular, at $t=0$ we have $V|_p$,
the tangent vector to the path at p .

$V \rightarrow \varphi$ Proof of (ii)

(15)

let x^a be coordinates in a neighborhood $p \in M$.

To find a curve through p with tangent vector $V|_p$ at p given φ , we solve the differential equation

$$V^a(x(t)) = \frac{d}{dt} \varphi^a(x) = \frac{d}{dt} \varphi^a(x(t)) \quad (\text{integrate})$$

for the components of V .

The solution exists and it is unique given the initial condition at $t=0$

$$\varphi_0(p) = p$$

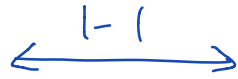
(by Picard's theorem on existence and uniqueness of 1st order differential equations with initial conditions)

Thus for every smooth vector field V , there is a family of integral curves on M .

In summary: For each $p \in M$

(15)

V
smooth vector
field on M



integral curves $x(t)$

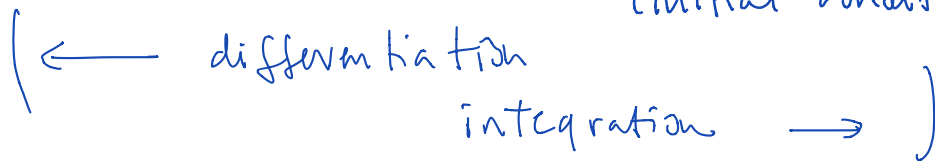
orbits of a 1-parameter

group of diffeomorphisms
 $\varphi(t, x) = \varphi_t^{\hat{}}(x) = \tilde{x}(t)$

given by

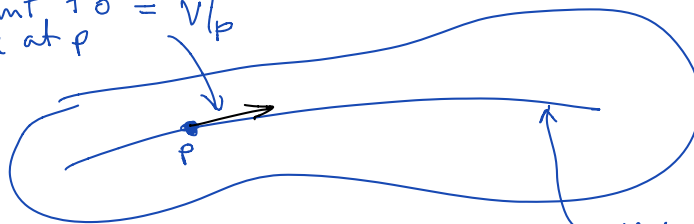
$$V^a(x(t)) = \frac{d}{dt} \varphi_t^{\hat{}}(x), \quad \varphi_0(p) = p$$

(initial conditions)



\times there is a unique curve $\varphi_t(p)$ through each $p \in M$
at $t=0$ st $V|_{\varphi_t(p)}$ is tangent to the point $\varphi_t(p)$

vector
tangent to
curve at $p = V|_p$



$\varphi_t(p) = x(t)$
integral curve
of V
(path of p on M)

Examples

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① $M = \mathbb{R}^2$ coordinates (x, y)

Find the integral curves of $X = \frac{\partial}{\partial x}$

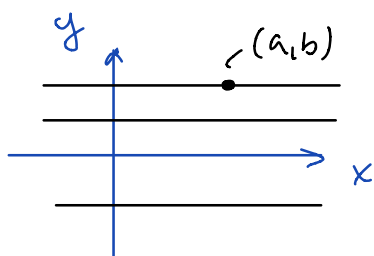
solve $V^a = \frac{d}{dt} \varphi_t^a$, $V^1 = 1$, $V^2 = 0$

$\varphi_t^a = x^a(t)$: $\varphi_t^1 = x(t)$, $\varphi_t^2 = y(t)$

and solve $1 = \frac{dx}{dt}$, $0 = \frac{dy}{dt}$

Integrating: $x(t) = t + a$ $\left\{ \begin{array}{l} a, b \\ \text{constants} \end{array} \right.$
 $y(t) = b$

so $\varphi_t = (t + a, b)$ is a unique curve through each point $(a, b) \in \mathbb{R}^2$



Integral curves are lines parallel to the x -axis

$X = \frac{\partial}{\partial x} \rightsquigarrow$ translations in the x -direction

$$(2) M = \mathbb{R}^2$$

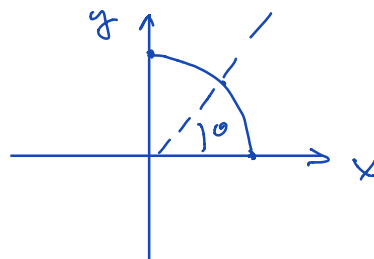
(17)

(i) Find a vector V st its integral curves are circles centered at the origin

Consider a circle of radius r

$$x = r \cos \theta$$

$$y = r \sin \theta$$



$$V^a = \frac{dx^a}{d\theta} \Rightarrow \begin{aligned} V^1 &= -r \sin \theta = -y \\ V^2 &= r \cos \theta = x \end{aligned}$$

$$\text{so } V = x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} = \frac{\partial}{\partial \theta}$$

\hookrightarrow generates rotations around the origin

(ii) Conversely: given $V = x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x}$

What are the integral curves of V ?

$$\text{solve: } -y(t) = \frac{dx}{dt}, \quad x(t) = \frac{dy}{dt} \quad (\text{coords } (x, y))$$

$$\text{so } x \frac{dx}{dt} + y \frac{dy}{dt} = 0$$

$$\text{so } x^2 + y^2 = \text{constant} \quad \text{circles centered at the origin}$$

① Tensor operations (part 2)

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(derivations)

5) Lie bracket of two vector fields
fields u, v is another vector field

$$Z = [u, v]$$

with components

$$Z^a = u^b \partial_b v^a - v^b \partial_b u^a$$

or equivalently

$$Z(f) = Z^a \partial_a f = u(V(f)) - v(U(f))$$

(one can prove that in fact Z is a vector)

Properties: let X, Y, Z be vector fields

- $[X, Y] = -[Y, X]$, $[X, X] = 0 \quad \forall X$
- bilinearity

$$[X, \alpha Y + \beta Z] = \alpha [X, Y] + \beta [X, Z]$$

for any constants α, β

- Jacobi identity

$$0 = [X, [Y, Z]] + [Y, [Z, X]] + [Z, [X, Y]]$$

(vectors & the binary operation $[\cdot, \cdot]$ form a Lie-algebra)

(geometric interpretation later)

want to generalize the notion of directional derivative along a vector X (19)

6) Definition: Lie derivative

let X be a vector field

(i) the Lie derivative of a function f along X is

$$L_X f = X(f) = X^a \partial_a f$$

ie the directional derivative of f along X

(ii) The Lie derivative of a vector field Y along X is

$$L_X Y = [X, Y]$$

(iii) The Lie derivative $L_X T$ of a (p,q) tensor T along X is defined by requiring the Lie derivative to satisfy the Leibnitz rule over tensor products and contractions ie

$$L_X (T \dots S \dots) = T \dots L_X S \dots + S \dots L_X T \dots$$

Note: $L_X T$ is of same type as T

One finds (by induction)

$$\begin{aligned} \mathcal{L}_X T^{a_1 \dots a_p}_{b_1 \dots b_q} &= X^c \partial_c T^{a_1 \dots a_p}_{b_1 \dots b_q} \\ &\quad - \sum_{i=1}^p T^{a_1 \dots a_{i-1} c a_{i+1} \dots a_p}_{b_1 \dots b_q} \partial_c X^{a_i} \\ &\quad + \sum_{i=1}^q T^{a_1 \dots a_p}_{b_1 \dots b_{i-1} c b_{i+1} \dots b_q} \partial_{b_i} X^c \end{aligned}$$

For example: let S be a 1-form. Then

$$\mathcal{L}_X S_a = X^b \partial_b S_a + S_b \partial_a X^b$$

Proof: Let Q be a vector field

Then $\mathcal{L}_X(\underbrace{Q^a S_a}_{\text{scalar}}) = X^b \partial_b (Q^a S_a)$

so $\mathcal{L}_X(Q^a S_a) = X^b (\partial_b Q^a) S_a + Q^a X^b \partial_b S_a$ ①

On the other hand: requiring that \mathcal{L}_X satisfies the Leibnitz rule we have

$$\begin{aligned} \mathcal{L}_X(Q^a S_a) &= (\mathcal{L}_X Q^a) S_a + Q^a \mathcal{L}_X S_a \\ &\stackrel{(ii)}{=} (X^b (\partial_b Q^a) - Q^b \partial_b X^a) S_a + Q^a \mathcal{L}_X S_a \quad \text{②} \end{aligned}$$

Equating ① & ②

$$Q^a \mathcal{L}_X S_a = Q^a (X^b \partial_b S_a + S_b \partial_a X^b) \quad \text{must be true } \forall Q$$

$$\Leftrightarrow \mathcal{L}_X S_a = X^b \partial_b S_a + S_b \partial_a X^b$$

Property:

$$d[x, y] = dx dy - dy dx$$

(sheet 1, and what is the meaning of this)

Later:

$L_X T \sim$ change of T along X
(change along curves generated by X)

\sim directional derivative of T along X

Remark: all operations defined so far (except (b)) are independent of a metric on M or a connection on M (depend only on the manifold structure)

7) Covariant differentiation ∇ (GR1)

$\partial_a T$ is not in general a tensor

Want: differentiation st

(p, q) tensor $\rightarrow (p, q+1)$ tensor

Need extra structure on M

Definition: a covariant derivative ∇ on M

is a linear ($\nabla(T+S) = \nabla T + \nabla S$) map st for any (p, q) tensor T , ∇T is a $(p, q+1)$ tensor which satisfies

(i) $\nabla_a f = \partial_a f$ (0,1) tensor $\nabla_X f = X(f)$

(ii) $\nabla_a V^b = \partial_a V^b + \Gamma^b_{ac} V^c$ where

Γ (connection coefficients) transform under coordinate transformations st

Γ is not a (1,1) tensor!

$\nabla_a V^b$ is a (1,1) tensor

[under $x^a \rightarrow \tilde{x}^a(x)$

$\tilde{\Gamma}^c_{ab} = \partial_f \tilde{x}^c \tilde{\partial}_a \tilde{x}^d \tilde{\partial}_b \tilde{x}^e \Gamma^f_{de} + (\partial_d \tilde{x}^c) \tilde{\partial}_a \tilde{\partial}_b \tilde{x}^d$]

(iii) Extend definition to other tensors by requiring that ∇ satisfies the Leibnitz rule

$$\nabla_a (T^{...} S^{...}) = (\nabla_a T^{...}) S^{...} + T^{...} \nabla_a S^{...}$$

Examples

$$\nabla_a W_b = \partial_a W_b - \Gamma_{ab}^c W_c \quad (1,1) \text{ tensor}$$

$$\nabla_c g_{ab} = \partial_c g_{ab} - \Gamma_{ca}^d g_{db} - \Gamma_{cb}^d g_{ad} \quad (2,2) \text{ tensor}$$

Definition: the torsion T of ∇ is defined as

$$T_{ab}^c = 2 \Gamma_{ca}^b - 2 \Gamma_{cb}^a$$

T is $(1,2)$ -tensor

GR: M is a differentiable manifold endowed with a torsion free ($T=0$)

covariant derivative ∇

$$\text{Then } T=0 \iff \Gamma_{ab}^c = \Gamma_{ba}^c$$

(Einstein-Cartan studied effects of $T \neq 0$)

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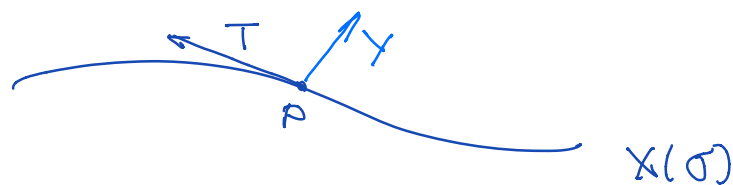
Geometrically: ∇ induces the notion of parallel transport along a curve $X^a(\sigma)$.

Suppose the curve has tangent vector T^a

Then

$$\frac{DY}{D\sigma} = \nabla_T Y = T^a \nabla_a Y$$

parallel transport
of Y along the
curve determined
by T



* allows for the identification of tangent spaces along T

* geodesics $\nabla_T T = 0$ (affine parametrization)

curve that parallel transports its own tangent vector

GN: TL geodesics \rightarrow trajectory of a free falling body in the gravitational field g

(E) Going back to the mathematical model of GR

metric
 \downarrow
 grav. field

M is a differentiable manifold with a torsion free connection ∇ and a metric g st $\nabla g = 0$

Theorem: Fundamental theorem of Riemannian geometry

Let g be a metric on a differentiable manifold M . Then there is a unique torsion free covariant derivative ∇ on M st

$$\nabla g = 0$$

This is the Levi-Civita connection

$$\Gamma^a_{bc} = \frac{1}{2} g^{ad} (\partial_b g_{cd} + \partial_c g_{bd} - \partial_d g_{bc})$$

Proof: GR 1

(see also Hughton & Tod for proof using \mathcal{L}_X)

Remark: $\nabla g = 0$ implies that the operation of "raising" and "lowering" indices with the metric commutes with ∇ .

Definition Curvature of ∇

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Assuming that Γ is once differentiable on each coordinate patch, we define the (1,3) tensor

$$(\nabla_b \nabla_c - \nabla_c \nabla_b) V^a = - \underline{R_{bcd}{}^a} V^d$$

for any vector field V on M

$R_{bcd}{}^a$: Riemann curvature tensor

($R_{bcd}{}^a \leftrightarrow$ gravitational field
analogue to F_{ab} in Maxwell's theory of EM)

One finds:

$$-R_{abc}{}^d = \partial_a \Gamma_{bc}^d + \Gamma_{ae}^d \Gamma_{bc}^e - (a \leftrightarrow b)$$

$$(\text{clearly } R_{abc}{}^d = -R_{bca}{}^d)$$

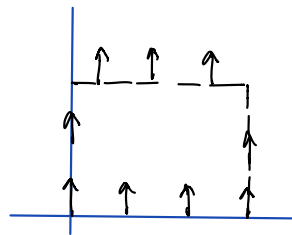
For a torsion free connection (GR 1)

$$(i) \quad R_{cabcd} = 0$$

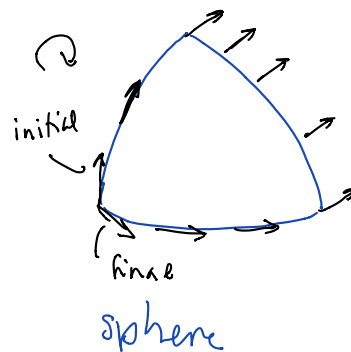
$$(ii) \quad \nabla_{[a} R_{bc]d}{}^e = 0 \quad (\text{Bianchi identity})$$

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The curvature is a "measure" of the change of a vector field when parallel transported around a loop in M



flat space



sphere

GR1

The geodesic deviation is an important interpretation of curvature as the gravitational field.

(we will use this when studying gravitational waves)

Describes the relative motion of a pair of nearby geodesics

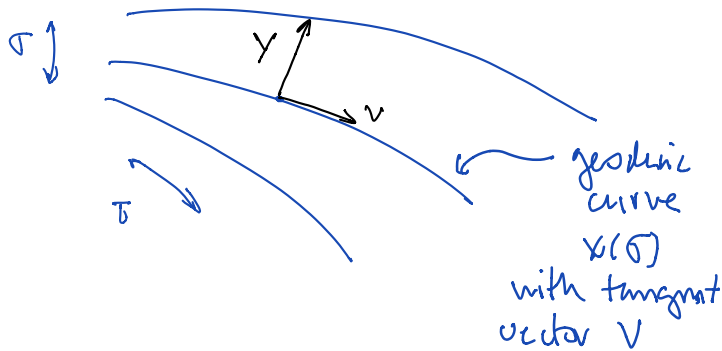
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Consider a one-parameter family of (affinely parametrised) geodesics

$$X^a(\tau, \sigma)$$

proper time τ

σ labels member of family



$$V^a = \frac{dX^a}{d\tau}$$

$$\Theta = V^a \nabla_a V^b = \nabla_V V^b$$

let Y be a vector field which measures the infinitesimal displacement of nearby geodesics as σ varies so Y is tangent to curves with $\tau = \text{constant}$.

That is: Y represents relative motion

Then:
$$\frac{D^2 Y^a}{D\tau^2} = \nabla_V^2 Y^a = -R_{bcd}{}^a V^b Y^c V^d$$

LHS = relative acceleration of nearby particles in free fall

Geodesic deviation: relative motion reveals presence of the gravitational field (curvature) as relative acceleration of nearby particles in free fall.

1.2 Killing vectors and isometries

From now on: M is a differentiable manifold with metric g and a torsion free connection ∇ st $\nabla g = 0$

Recall that $L_X T$ is defined using only the manifold structure

However for a torsion free connection

$$L_X f = X^a \partial_a f = X^a \nabla_a f$$

$$L_X Y^a = [X, Y]^a = X^b \nabla_b Y^a - Y^b \nabla_b X^a$$

(extra terms cancel if $T=0$ ie $\Gamma^a_{bc} = \Gamma^a_{cb}$)

etc...

Consider $L_X g$: we have (exercise)

$$L_X g_{ab} = X^c \nabla_c g_{ab} + g_{cb} \nabla_a X^c + g_{ac} \nabla_b X^c$$

If $\nabla g = 0$ then: $L_X g_{ab} = 2 \nabla_{(a} X_{b)}$

Definition: A Killing vector field is a vector X

st $L_X g = 0$ ie $\nabla_{(a} X_{b)} = 0$

(X is a direction along which g is unchanged)

Properties:

(a) closure under the Lie bracket
If L & K are Killing vectors so is $[L, K]$

(b) $\nabla_a \nabla_b K_c = -R_{bca}{}^d K_d$
for any Killing vector K .

(c) Killing vectors give rise to
geodesic integrals of motion

ie let V be tangent to a geodesic.

Then, $K_a V^a$ is a constant along the geodesic.

[easier to show: gives 1st order diff eq for $x^a(\tau)$
vs 2nd order Euler Lagrange eqs]

Proof

$$(a) \mathcal{L}_{[K, L]} g = \mathcal{L}_K \mathcal{L}_L g - \mathcal{L}_L \mathcal{L}_K g = 0$$

(b) For any vector field

$$\nabla_{[a} \nabla_{b]} K_c = \frac{1}{2} R_{abc}{}^d K_d$$

Recall Bianchi identity: $R_{cab}{}^d = 0$. Then

$$0 = \nabla_{[a} \nabla_{b]} K_c + \nabla_{[c} \nabla_{a]} K_b + \nabla_{[c} \nabla_{b]} K_a$$

$$0 = \nabla_a \nabla_b K_c - \nabla_b \nabla_a K_c + \nabla_b \nabla_c K_a - \nabla_c \nabla_b K_a + \nabla_c \nabla_a K_b - \nabla_a \nabla_c K_b$$

$$\begin{array}{l} \text{Using Killing} \\ \text{equations} \end{array} \Rightarrow \begin{array}{l} + \nabla_b \nabla_c K_a \\ - \nabla_c \nabla_a K_b \\ + \nabla_a \nabla_b K_c \end{array}$$

$$= 2 (\nabla_a \nabla_b K_c + \nabla_b \nabla_c K_a - \nabla_c \nabla_b K_a)$$

$$\Rightarrow \nabla_a \nabla_b K_c = -2 \nabla_{[c} \nabla_{b]} K_a = -R_{bca}{}^d K_d$$

(c) GRL 1

let V be tangent to an affinely parametrised geodesic $x^a(\tau)$. Then

$$\begin{aligned} \frac{d}{d\tau} (K_a V^a) &= \frac{dx^b}{d\tau} \partial_b (K_a V^a) = V^b \nabla_b (K_a V^a) \\ &= V^b (\nabla_b K_a) V^a + K_a \tilde{V}^b \nabla_b V^a \rightarrow 0 \text{ by geodesic equation} \\ &= V^a V^b \nabla_b K_a = 2 V^a V^b \nabla_{(b} V_{a)} \\ &= 0 \text{ for } V \text{ a Killing vector} \end{aligned}$$

so $K_a V^a$ constant along the geodesic. //

Definition: an isometry is a diffeomorphism of M which leaves the metric invariant

$$\mathcal{L}_V g = 0$$

where V is the (Killing) vector which generates the diffeomorphisms.

Note: extensive discussion of isometries in problem sheets.

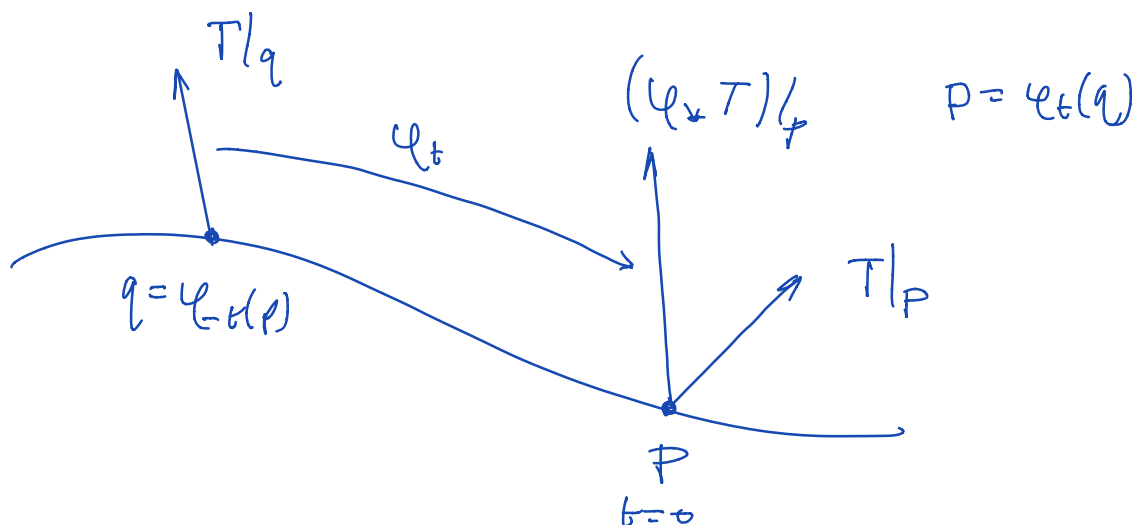
Let T be a tensor type (p, q) and $\textcircled{32}$
 V a smooth vector field on M .

We know that at each point $p \in M$ V
 defines a 1-parameter group of diffeomorphisms
 φ_t (φ_t takes a point $p \in M$ a parameter
 distance t along the integral curves of V)

The diffeomorphism φ_t induces a linear
 map φ_{t*} on tensors (push forward)

$$\varphi_{t*} : T|_q \longrightarrow \varphi_{t*} T|_{\varphi_t(q)}$$

where $\varphi_{t*} T|_{\varphi_t(q)}$ is the tensor T pushed
 along the integral curve (flow of V) a
 distance t (φ_{t*} preserves type)



Definition: The Lie derivative of T along a vector V , denoted $\mathcal{L}_V T$, is the rate of change of T along the integral curves of V , that is, it is the rate of change of T under ϕ_t

$$\mathcal{L}_V T|_p = \lim_{t \rightarrow 0} \frac{1}{t} (T|_p - \phi_{t*} T|_p)$$

Of course this needs to be equivalent to our previous definition. (see page 19)

Properties

✓ ① $d_v T$ maps (p, q) tensors into (p, q) tensors

② $d_v T$ is linear: \forall constants α, β

✓
$$d_v (\alpha T + \beta S) = \alpha d_v(T) + \beta d_v(S)$$

③ d_v obeys the Leibnitz rule

✓
$$d_v(TS) = T d_v S + (d_v T) S$$

④ $d_v f = V(f)$

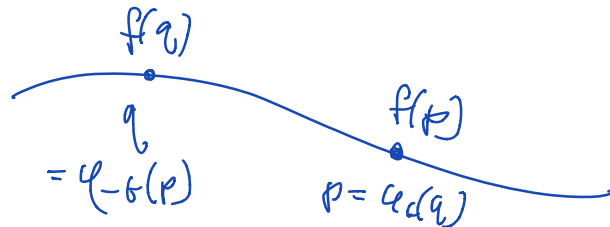
⑤ $d_v u = [u, V]$

①, ②, ③ : proofs as in calculus

Proof of (4)

(35)

$$d_v f|_p = \lim_{t \rightarrow 0} \frac{1}{t} (f|_p - \psi_{t*} f|_p)$$



$$\begin{aligned} \text{In this case } (\psi_{t*} f)(p) &\equiv f(q) \\ &= f(\psi_{-t}(p)) \end{aligned}$$

$$\text{ie } \psi_{t*} \circ f = f \circ \psi_{-t}$$

ie $\psi_{t*} f$ = function on M whose value at $p = \psi_t(q)$ is the value of the function at q

Hence

$$d_v f|_p = \lim_{t \rightarrow 0} \frac{1}{t} (f|_p - f|_{\psi_{-t}(p)}) = \left. \frac{df}{dt} \right|_p$$

$p \sim q + t$

[Compare : let $\{x^a\}$ be local coords in a neighborhood of p st and let V be the integral curve through p with $\frac{dx^a}{dt} = V^a(x(t)) \Rightarrow d_v f = V^a \partial_a f = \frac{df}{dt} = V(t)] //$

Proof of ⑤

Let u be a vector on M and f a smooth function on M

Recall notation $u(f) = u^a \partial_a f$

$$L_v(u(f)) = V(u(f)) \quad \text{by } \textcircled{4}$$

Using the Leibnitz rule

$$L_v(u(f)) = (L_v u)(f) + u(L_v f)$$

Thus

$$V(u(f)) = (L_v u)(f) + u(V(f))$$

$$\text{Hence} \quad L_v u = [V, u]$$

Remark: One can also prove this directly from the definition. In this case

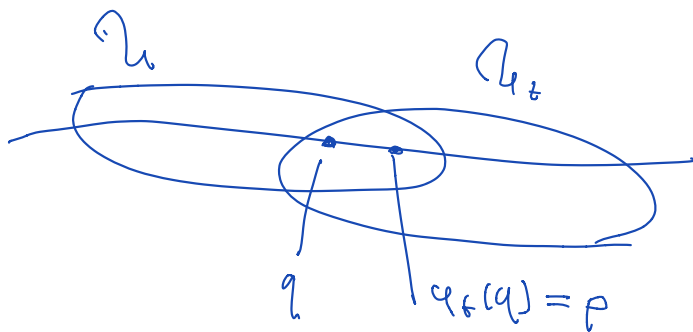
$$\begin{array}{ccc} \varphi_{t*} : TM|_q & \longrightarrow & TM|_{\varphi_t(q)=p} \\ \downarrow \psi & & \downarrow \psi \\ u_q & \longmapsto & (\varphi_{t*} u)|_{\varphi_t(q)} \end{array}$$

where $(\varphi_{t*} u)|_{\varphi_t(q)}$ is the vector tangent to the integral curves of u at $\varphi_t(q) \in M$.



Equivalently: φ_t (or any diffeomorphism) (3+)
induces a coordinate transformation

$$\begin{array}{ccc}
 \varphi_t: & M & \longrightarrow & M \\
 & U & & U \\
 & u & \longrightarrow & u_t = \varphi_t(u) \\
 & \psi & & \psi \\
 & q & \longmapsto & \varphi_t(q) = p \\
 & \downarrow & & \downarrow \\
 & X^a(q) & \longmapsto & \tilde{X}^a(\varphi_t(q)) = \tilde{X}^a(p) \\
 & & \tilde{X} \circ \varphi_t \circ X^{-1} &
 \end{array}$$



Consider an infinitesimal coordinate change induced by a 1-parameter group of diffeomorphisms φ_t :

$$\begin{array}{ccc}
 \tilde{X}^a & = & X^a + t \left. \frac{dX^a}{dt} \right|_p + \dots \\
 \uparrow \text{of } p & & \uparrow \text{of } q \\
 & & \tilde{V}^a|_p = \left. \frac{dX^a}{dt} \right|_p
 \end{array}$$

(38)

Then: for a function f

$$\begin{aligned}(\psi_{t*} f)(p) &= f(q) \\ &= f(\tilde{x} - t v^a(x)) \\ &= f(\tilde{x}) - t v^a(x) \partial_a f + \dots \\ &= f(p) - t v^a(x) \partial_a f + \dots\end{aligned}$$

$$\text{so } v^a(x) \partial_a f = \text{div} f = \frac{df}{dt}$$

Remark: one can do this for other tensors
In particular:

$$(\psi_{t*} T)_p = \tilde{T}(q) = \tilde{T}(\tilde{x} - t v^a(x))$$

In summary:

Killing vectors \leftrightarrow isometries

$$\nabla_a V_b = 0$$

$$d_V g = 0$$

"metric invariant"

- * Extensive discussion of isometries in problem sheets, with examples!
- * Important consequence for g :
allows for an expression for g which is independent of one of the coordinates

Corollary of theorem page 13

(40)

Let V be a non-zero vector. Then there exist coordinates on M $\{x^a\} = \{t, y^i\}$
 $a=0, 1, \dots, n-1$ \leftarrow $i=1, \dots, n-1$

st locally $V = \partial/\partial t$.

Proof: Consider a hypersurface $\Sigma_0 \subset M$ defined by a function $f(\bar{x}) = 0$ where \bar{x}^a are coordinates on M .

WLOG: take $\Sigma_0 \subset V$ st V is not tangent to Σ_0

Choose $\bar{x}^0 = t = f(\bar{x})$ so Σ_0 is the hypersurface $t=0$ with coordinates $(0, \underline{y})$ where $\underline{y} = (y_1, y_2, \dots, y_{n-1})$

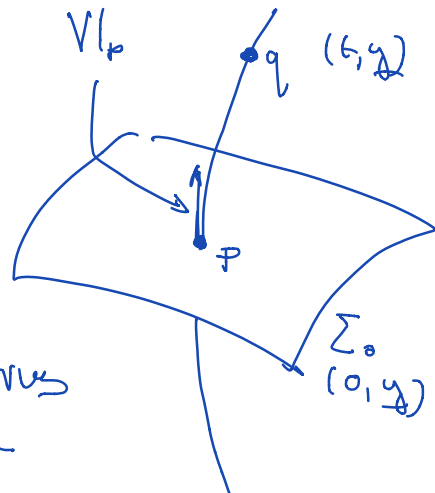
Given $V = \partial_t$: there is unique integral curve of V through each point $p \in \Sigma_0$ (with tangent $V|_p$ at that point) which is the unique soln of

$$\frac{dx^a(t)}{dt} = V^a = 1; \quad \frac{dy^i}{dt} = V^i = 0$$

with initial condition

$$x^a(p) = (0, \underline{y})$$

Thus $\underline{y} = \text{constant}$ on these curves and the coordinates of q are (t, \underline{y}) . //



[Remark: $\text{Span}\{V\}$ is tangent to $\textcircled{4}$
 a foliation by curves $\gamma = \text{constant}$
 (the leaves which are the integral
 curves of V with parameter t)
 $V \rightarrow$ tangent to the leaves]

Important property of the metric
 when there are Killing vectors:

Let $K \neq 0$ be a vector field on M . Then
 There are coordinates (x^0, x^1, x^2, x^3) st
 $K = K^0 \partial_a = \partial_0$ ie $K^0 = 1, K^i = 0 \quad i=1,2,3$

K a Killing vector: $\partial_0 g_{ab} = 0$

that is, the metric does not depend on the
 coordinate x^0 (ie g does not change along
 the integral curves of V)

Proof: K is a Killing vector iff $\mathcal{L}_K g_{ab} = 0$

$$\text{iff } 0 = K^c \partial_c g_{ab} + \cancel{g_{ca} \partial_b K^c} + \cancel{g_{cb} \partial_a K^c}$$

$$\text{iff } 0 = K^0 \partial_0 g_{ab} = \partial_0 g_{ab} //$$

Next: one can generalize the notion of integral curves.

let $W = \text{span}\{u_1, \dots, u_n\}$

set of linearly independent vector fields

Then $\forall p \in M, W|_p \subset TM|_p$

Question: let $p \in M$. Is there

$S \subset M$ (embedded submanifold) through p st if $TS|_p = W|_p$, then for each $q \in S$ we have

$$TS|_q = W|_q \quad ?$$

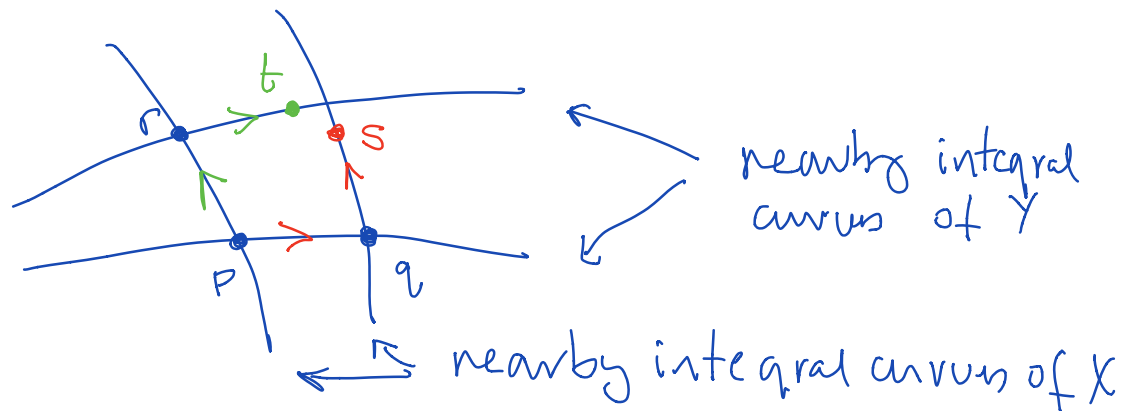
$n = 1 \quad S \rightarrow$ integral curves

$n > 1 \quad \underline{\text{NEED}} \quad [u_i, u_j] \in W \quad \forall i, j = 1, \dots, n$

1.3 Geometry of $[X, Y]$

(43)

Let X, Y be smooth vector fields on M .
Consider the integral curves of X & Y



Let $p \in M$: compare $YX(p)$ & $XY(p)$

XY: let $q = p + \epsilon Y(p)$
 $s = q + \epsilon X(q)$

$$s = p + \epsilon Y(p) + \epsilon X(p + \epsilon Y(p))$$

$$= p + \epsilon Y(p) + \epsilon X(p) + \epsilon^2 XY(p)$$

YX: let $r = p + \epsilon X(p)$
 $t = r + \epsilon Y(r)$

$$t = p + \epsilon X(p) + \epsilon Y(p) + \epsilon^2 YX(p)$$

$$\therefore s - t = \epsilon^2 [X, Y] \Big|_p$$

[Notation: $X(p) \rightarrow$ action of 1-parameter group of diffeomorphisms on p i.e. $\varphi_t(p)$]

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$[X, Y] = 0$ (or more generally

$[X, Y] \in \text{span}\{X, Y\}$ as we will see later)

\Rightarrow the set of all points that can be reached along integral curves of X & Y from a given point $p \in M$ forms a 2-dim submanifold $S \subset M$ through p

The tangent space to S at any $q \in S$ is the same as the tangent space to S at p :

let $W = \text{span}\{X, Y\}$ & $TS|_p = W|_p \subset TM|_p$

then $TS|_q = W|_q \quad \forall q \in S$.

lecture #5

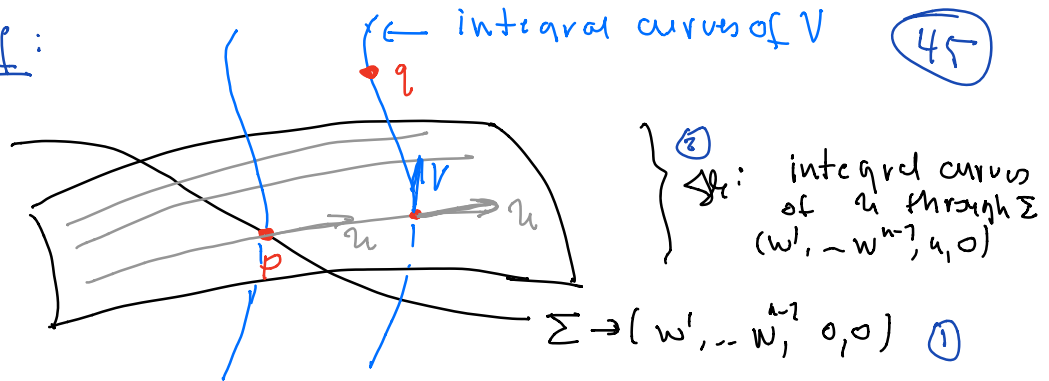
Theorem: (Frobenius for $k=2$)

Let u, v be linearly independent smooth vector fields on M . Then there are coords

$\{w^1, \dots, w^{n-2}, u, v\}$

st $u = \frac{\partial}{\partial u}$, $v = \frac{\partial}{\partial v}$ iff $[u, v] = 0$

Proof:



- ① Let Σ be a $(n-2)$ -surface. Then
 $\forall p \in \Sigma$ choose coords $(w^1, \dots, w^{n-2}, 0, 0)$
- ② Let \mathcal{H} be a hypersurface consisting of integral curves of U through Σ .
 let u : parameter length wrt U for the integral curves of U through Σ
 Then: coords of points on \mathcal{H} $(w^1, \dots, w^{n-2}, u, 0)$
- ③ Let q be a point near \mathcal{H} on an integral curve of V through \mathcal{H} .
 let v : parameter length wrt V for the integral curves of V through \mathcal{H}
 coords of q : $(w^1, w^2, \dots, w^{n-2}, u, v)$

Tangent vectors at q_i :

(246)

$$V = \partial_v, \quad U = a \partial_u + b \partial_v + s^i \frac{\partial}{\partial w^i}$$

$$\text{st } \underset{v=0}{\text{on } \mathcal{H}_i} \left. \begin{array}{l} s^i = 0 \quad \forall i=1, \dots, n-2 \\ b=0, a=1 \end{array} \right\} \begin{array}{l} U \text{ tangent} \\ \text{to its} \\ \text{integral} \\ \text{curves} \end{array}$$

Let f be a function on M

$$\begin{aligned} [U, V](f) &= U(V(f)) - V(U(f)) \\ &= U(\partial_v(f)) - \partial_v(a \partial_u f + b \partial_v f + s^i \partial_i f) \\ &= U(\partial_v(f)) - (\partial_v a) \partial_u f - (\partial_v b) \partial_v f - (\partial_v s^i) \partial_i f - \cancel{\partial_u(a \partial_v f)} \\ &= -(\partial_v a) \partial_u f - (\partial_v b) \partial_v f - (\partial_v s^i) \partial_i f \end{aligned}$$

Then $[U, V](f) = 0$ for any arbitrary differentiable function f .

$$\text{iff } \partial_v a = 0, \quad \partial_v b = 0, \quad \partial_v s^i = 0 \quad \forall i$$

(because $\partial_u, \partial_v, \partial_i$ are linearly independent)

$$\text{iff } a = a(u, w^i), \quad b = b(u, w^i), \quad s^i = s^i(u, w^i)$$

ie a, b, s^i independent of v

On \mathcal{H}_i , where $v=0$: $U = \partial_u$

$$\text{ie } a=1, \quad b=0, \quad s^i=0$$

$$\therefore b=0, \quad s^i=0 \quad \forall i, \quad \text{and } a=1 \quad \& \quad U = \partial_u$$

Remark:

(47)

Lemma: Let \hat{u} and \hat{v} be linearly independent vector fields such that

$$[\hat{u}, \hat{v}] \in \text{Span} \{ \hat{u}, \hat{v} \}$$

Thm: there exist u, v with

$$\text{Span} \{ u, v \} = \text{Span} \{ \hat{u}, \hat{v} \}$$

and $[u, v] = 0$

Proof: let $u = \lambda \hat{u}$, $v = \mu \hat{v}$

$$\begin{aligned} [u, v](f) &= [\lambda \hat{u}, \mu \hat{v}](f) = \lambda \hat{u}(\mu \hat{v}(f)) - \mu \hat{v}(\lambda \hat{u}(f)) \\ &= \lambda \hat{u}(\mu) \hat{v}(f) + \lambda \mu \hat{u} \hat{v}(f) - \mu \hat{v}(\lambda) \hat{u}(f) - \mu \lambda \hat{v} \hat{u}(f) \\ &= \lambda \hat{u}(\mu) \hat{v}(f) - \mu \hat{v}(\lambda) \hat{u}(f) + \lambda \mu [\hat{u}, \hat{v}](f) \end{aligned}$$

But $[\hat{u}, \hat{v}] \in \text{Span} \{ \hat{u}, \hat{v} \}$, thus for some a, b

$$[\hat{u}, \hat{v}](f) = a \hat{u} + b \hat{v}$$

Thus:

$$\begin{aligned} [u, v](f) &= \lambda \mu \left\{ \frac{1}{\mu} \hat{u}(\mu) \hat{v}(f) - \frac{1}{\lambda} \hat{v}(\lambda) \hat{u}(f) + a \hat{u}(f) + b \hat{v}(f) \right\} \\ &= \lambda \mu \left\{ \left(\frac{1}{\mu} \hat{u}(\mu) + b \right) \hat{v} + \left(-\frac{1}{\lambda} \hat{v}(\lambda) + a \right) \hat{u} \right\} (f) \end{aligned}$$

For $[u, v] = 0 \quad \forall f$: solve for λ, μ

$$\hat{u}(\mu) = -\mu b \qquad \hat{v}(\lambda) = \lambda a$$

ie $[u, v] = 0$ for λ, μ ($u = \lambda \hat{u}$, $v = \mu \hat{v}$)

$$\text{st } \hat{u}(b \mu) = -b \qquad \hat{v}(a \lambda) = a \quad //$$

Corollary:

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$[X, Y] \in \text{span}\{X, Y\}$ iff
the $\text{span}\{X, Y\}$ is tangent to a family
of 2-surfaces (integral surfaces)

On the surfaces $\{w_1, \dots, w^{n-2}\}$ is constant
and $\{u, v\}$ parametrize the surfaces
(We say that the $\text{span}\{X, Y\}$ is integrable)

Proof: by lemma & theorem.

Theorem: Frobenius

Let $W = \text{span}\{X_1, \dots, X_k\}$

↪ set of linearly independent
smooth vector fields.

$[X_i, X_j] \in W$ iff W is tangent to a
family of k -dim surfaces (integral surfaces).

The integral surfaces are parametrized
by coordinates (u_1, \dots, u_k) and given by
 $(w_1, \dots, w^{n-k}) = \text{constant}$.

Proof: induction on k

// not
examenable

(49)

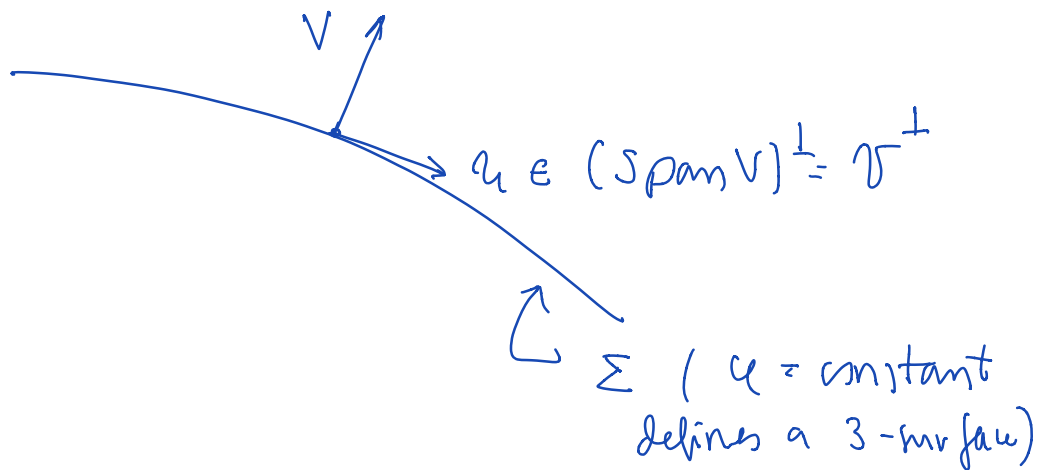
Dual versions

Theorem: let $V_a \neq 0$ be a smooth 1-form

$V_a \nabla_b V_c = 0$ iff there are functions φ, ψ st $V_a = \psi \partial_a \varphi$ //

Definition: such V_a is perpendicular to the hypersurface Σ defined by $\varphi = \text{constant}$ and it is said to be hypersurface orthogonal (HSO).

Recall: V_a is perpendicular to the hypersurface Σ ($\varphi = \text{constant}$) iff $V_a u^a = 0$ \forall vectors u which are tangent to Σ



Proof

(50)

(\Leftarrow) straightforward & computational

Suppose there are functions ψ, φ st

$$V_a = \psi \partial_a \varphi$$

$$\begin{aligned} V_a \nabla_{cb} V_{c\gamma} &= \psi \partial_a \varphi \nabla_{cb} (\psi \partial_{c\gamma} \varphi) \\ &= \psi \partial_a \varphi (\partial_{cb} (\psi \partial_{c\gamma} \varphi) - \cancel{\Gamma_{cbb}^d} \psi \partial_d \varphi) \\ &= \psi \partial_a \varphi (\partial_{cb} \psi \partial_{c\gamma} \varphi + \cancel{\psi \partial_{cb} \partial_{c\gamma} \varphi}) \\ &= \psi \partial_a \varphi \partial_{cb} \psi \partial_{c\gamma} \varphi \end{aligned}$$

$$\exists V_a \nabla_{cb} V_{c\gamma} = \exists \psi \partial_a \varphi \partial_b \psi \partial_{c\gamma} \varphi = 0$$

(\Rightarrow) Let $\mathcal{V}^\perp = (\text{Span}\{V\})^\perp$ orthogonal complement of $\text{Span}\{V\}$ in TM

Recall: a vector $W \in \mathcal{V}^\perp$
if $g_{ab} W^a X^b = W_a X^b = 0, \forall X \in \text{Span}\{V\}$

Claim: $V_{ca} \nabla_b V_{c\gamma} = 0 \Rightarrow [X, Y] \in \mathcal{V}^\perp$
 $\forall X, Y$ linearly indep vectors in \mathcal{V}^\perp

If true: by Frobenius, there is a family of $(n-1)$ -surfaces tangent to V^\perp and V is perpendicular to these surfaces.

Then we can pick ψ st $\psi = \text{constant}$ on each member of this family, and

$$\partial_a \psi \quad (\text{"gradient"})$$

is perpendicular to these surfaces i.e.

$$V_a = \psi \partial_a \psi \quad \text{for some function } \psi$$

Proof of the claim: (TP: if V st $V_a \nabla_b V_c = 0$
 $\Rightarrow [X, Y] \in V^\perp, \forall X, Y$ lin. indep vectors in V^\perp)

want to prove that this vanishes.

$$\begin{aligned} [X, Y]^a V_a &= (X^b \nabla_b Y^a - Y^b \nabla_b X^a) V_a \\ &= X^b (\nabla_b (Y^a V_a) - Y^a \nabla_b V_a) - Y^b (\nabla_b (X^a V_a) - X^a \nabla_b V_a) \\ &= -X^a Y^b (\nabla_a V_b - \nabla_b V_a) \end{aligned}$$

$$\text{so } [X, Y]^a V_a = -2 X^a Y^b \nabla_{[a} V_{b]} \quad (*)$$

Now: $V_{ca} \nabla_b V_{ca} = 0$ if $\nabla_b V_{ca} = 0$ (52)

$$X^b Y^c \nabla_{ca} V_{ca} = 0 \quad \forall X, Y \in \mathcal{V}^\perp$$

So

$$0 = X^b Y^c \nabla_{ca} V_{ca}$$

$$= \frac{1}{3} X^b Y^c (V_a \nabla_{cb} V_{ca} + \cancel{V_b \nabla_{cc} V_{ca}} + \cancel{V_c \nabla_{ca} V_{cb}})$$

\uparrow as $X^b V_b = 0$ \uparrow as $Y^c V_c = 0$

$$(*) \rightarrow = \frac{1}{3} V_a \left(-\frac{1}{2} [X, Y]^b V_b \right)$$

$\therefore [X, Y]^b V_b = 0$ i.e. $[X, Y] \in \mathcal{V}^\perp$
 and \mathcal{V}^\perp (is integrable) is tangent
 to a family of $(n-1)$ -surfaces //

not in the lectures

General: Let $\mathcal{V} = \text{Span} \{ \underbrace{V^{(1)}, \dots, V^{(h)}}_{\text{set of linearly independent 1-forms}} \}$

Let $X, Y \in \mathcal{V}^\perp$. Then

$$[X, Y] \in \mathcal{V}^\perp \text{ if } [X, Y]^a V_a = -2 X^a Y^b \nabla_{ca} V_b = 0 \quad \forall V \in \mathcal{V}$$

$$\text{if } \nabla_{ca} V_b = \sum_{i=1}^h \beta^{(i)}_{ca} V_b^{(i)}, \quad \forall V \in \mathcal{V}$$

where $\beta^{(i)}$ are arbitrary 1-forms

if there is a family of $(n-h)$ hypersurfaces
 tangent to \mathcal{V}^\perp

(53)

Important example

Let K be HSO (but not necessarily Killing)
then $g_{0i} = 0 \quad i=1, 2, 3$

Proof: $K_a = g_{ab} K^b = \psi \partial_a \varphi$

(K perpendicular to a 3 dim hypersurface defined by $\varphi = \text{constant}$)

We can choose coordinates st $\varphi = x^0$
and $K = \partial_0$

Then $K_0 = g_{00} K^0 = g_{00} = \psi \partial_0 \varphi = \psi$

$$K_i = g_{i0} K^0 = g_{i0} = \psi \partial_i \varphi = 0$$

$\uparrow x^0$

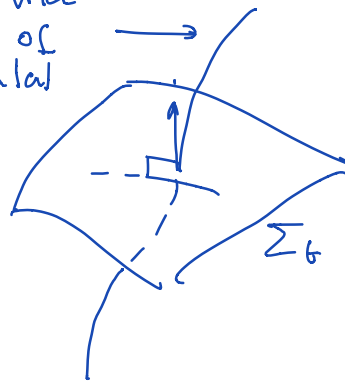
$$\therefore ds^2 = \psi dx^{*0} + g_{ij} dx^i dx^j \quad //$$

If K is also Killing then g_{ab} is independent of x^0

Applications:

- cosmology

time-like
geodesics of
fundamental
observers
 u is HSO



$t = \text{constant}$

- Schwarzschild black hole

\rightsquigarrow static soln of Einstein's eqs in vacuum

\hookrightarrow there is a HSO TL KV

- stationary solns: TL K-V.

- Think about the case where

$$W = \text{Span}\{K_1, \dots, K_n\}$$

K_i are Killing vectors and the fact that $[K_i, K_j]$ is also a Killing vector

End of Chapter 1