

(i) L finite dimensional \mathfrak{g} -module

$L^* = \{f: L \rightarrow \mathbb{C}\}$ as a vector space. The action on L^* is given by $(x \cdot f)(l) = -f(x \cdot l)$, $x \in \mathfrak{g}$, $l \in L$, $f \in L^*$.

Define $A^\perp \subset L^*$ by $A^\perp = \{f \in L^* \mid f(a) = 0 \forall a \in A\}$ for every submodule $A \subseteq L$. Then A^\perp is a submodule for L^* . Moreover if A is proper in L , then A^\perp is a proper submodule of L^* . Hence if L is not simple, L^* is not simple.

Conversely, if $B \subset L^*$ is a submodule, define $B^\perp \subset L$ by $B^\perp = \{l \in L \mid f(l) = 0 \forall f \in B\}$. ~~etc.~~ etc.

(ii) $L(\lambda)$, $\lambda \in \mathcal{P}^+$ simple \mathfrak{g} -module. By (i), $L(\lambda)^*$ is also a simple \mathfrak{g} -module. It is easy to see that the weights of $L(\lambda)^*$ are precisely the negatives of the weights of $L(\lambda)$ (~~use~~ pick a basis of weight vectors of $L(\lambda)$ and take the dual basis in $L(\lambda)^*$; these are weight vectors of $L(\lambda)^*$). Moreover, we know that

$$\dim L(\lambda)_\mu = \dim L(\lambda)_{w(\mu)} \quad \forall w \in W, \mu \text{ weight of } L(\lambda)$$

This means that $\dim L(\lambda)_{w(\lambda)} = 1 \quad \forall w \in W$. In particular, $\dim L(\lambda)_{w_0(\lambda)} = 1$, hence $\dim L(\lambda)^*_{-w_0(\lambda)} = 1$.

~~Notice~~ Notice $-w_0(\lambda) \in \mathcal{P}^+$ iff $\lambda \in \mathcal{P}^+$. This is because $-w_0(\alpha) > 0$ for all $\alpha > 0$ (positive root). Moreover $-w_0(\lambda)$ is a maximal weight, otherwise $\exists \mu \notin \text{weight of } L(\lambda)$ (so $-\mu$ is a weight of $L(\lambda)^*$) s.t. $-\mu - (-w_0(\lambda)) \in \mathcal{Q}^+$. Apply $-w_0$ remembering that $-w_0(\mathcal{Q}^+) = \mathcal{Q}^+ \Rightarrow w_0(\mu) - \lambda \in \mathcal{Q}^+$. But $w_0(\mu)$ is also a weight of $L(\lambda)$, contradiction!

(iii) Claim 0 occurs as a weight of $L(\lambda)$ if and only if λ is a sum of $\sqrt{\text{roots}}$. This follows immediately from the following remark: By using $\mathfrak{sl}(2)$ -theory for each \mathfrak{sl}_α , α positive root, one can show that μ is a weight of $L(\lambda)$ if and only if $w(\mu) \leq \lambda$ for all $w \in W$ (one direction is immediate since μ weight $\Rightarrow w(\mu)$ is a weight)

② Weyl's dimension formula: $\lambda \in \mathcal{P}^+$

$$\dim L(\lambda) = \prod_{\alpha \in \Phi^+} \frac{\langle \lambda + \rho, \alpha \rangle}{\langle \rho, \alpha \rangle} \quad k \in \mathbb{N}$$

Take $\lambda = (k-1)\rho \in \mathcal{P}^+$. Then $\langle \lambda + \rho, \alpha \rangle = k \langle \rho, \alpha \rangle$
 $\Rightarrow \dim L((k-1)\rho) = k^{|\Phi^+|}$

③ $\omega_1, \dots, \omega_n$ are an $\mathbb{Z}_{\geq 0}$ -basis of \mathcal{P}^+ , i.e., every $\lambda \in \mathcal{P}^+$ can be written uniquely as $\lambda = a_1 \omega_1 + \dots + a_n \omega_n$ for $a_i \in \mathbb{Z}_{\geq 0}$.
 It is easy to see that ~~the~~ the weight vectors of a tensor product $V \otimes U$ are precisely ~~the~~ $v \otimes u$ where v, u are weight vectors of V and U , respectively. Moreover, the weights of $V \otimes U$ are ~~the~~ $\lambda + \mu$, λ weight of V , μ weight of U . Hence ~~if~~ if λ, μ are highest weights of V, U , resp., then $\lambda + \mu$ is a highest weight of $V \otimes U$. From complete reducibility then $L(\lambda + \mu)$ occurs as a summand of $L(\lambda) \otimes L(\mu)$. By induction, it follows that $L(\lambda)$ is a direct summand of $L(\omega_1)^{\otimes a_1} \otimes \dots \otimes L(\omega_n)^{\otimes a_n}$ when $\lambda = a_1 \omega_1 + \dots + a_n \omega_n$.

④ (i) We think of $\mathfrak{g} \cong \{h \in \mathbb{C}^n \mid \sum h_i = 0\}$ and $\mathfrak{g}^* = \mathbb{C}^n / \langle \varepsilon_1 + \dots + \varepsilon_n \rangle$.
 Then $w_i = \varepsilon_1 + \dots + \varepsilon_i \quad 1 \leq i \leq n-1$. From a previous problem sheet (or easily by a direct computation) $e_1 \dots e_i$ is a highest weight vector with weight w_i of $\wedge^i V$. This means that $L(w_i)$ is a summand of $\wedge^i V$. We compute dimensions to show they are equal (there are many other ways to show equality).

$\dim \wedge^i V = \binom{n}{i}$. Use Weyl's dim formula to compute $\dim L(w_i)$:
 $\rho = (n, n-1, \dots, 1)$, $\alpha_{ij} = \varepsilon_i - \varepsilon_j \quad i < j$
 so $\langle \rho, \alpha_{ij} \rangle = [n - (i-1)] - [n - (j-1)] = j - i$ etc

The case $L(\omega_2)$ is similar (in fact it's just the table before to⁻⁴⁻ which we apply $-\omega_0 = -(13)$)

(iii) If we write the coordinates of h as (x_1, x_2, x_3) with $x_1 + x_2 + x_3 = 0$

then

$$\text{ch}_{L(\omega_1)}(x_1, x_2, x_3) = \frac{e^{4x_1+2x_2+x_3} - e^{2x_1+4x_2+x_3} - e^{4x_1+x_2+2x_3} - e^{x_1+2x_2+4x_3} + e^{x_1+4x_2+2x_3} + e^{2x_1+x_2+4x_3}}{e^{3x_1+2x_2+x_3} - e^{2x_1+3x_2+x_3} - e^{3x_1+x_2+2x_3} - e^{x_1+2x_2+3x_3} + e^{x_1+3x_2+2x_3} + e^{2x_1+2+3x_3}}$$

as functions $\mathbb{C}^3 \rightarrow \mathbb{C}$.

$$\xi = e^{x_1} + e^{x_2} + e^{x_3}$$

and similarly for $L(\omega_2)$.

(6) (i) Easy

(ii) h is spanned by $H_i = E_{ii} - E_{\alpha_i, \alpha_i}$. Define the dual basis $\{\epsilon_i\}$ of h^* s.t. $\epsilon_i(H_j) = \delta_{ij}$.

The roots are:

- $\epsilon_i - \epsilon_j, i \neq j$ with root vector $E_{ij} - E_{\alpha_j, \alpha_i}$
- $\epsilon_i + \epsilon_j, i < j$ $E_{i, \alpha_j} + E_{j, \alpha_i}$
- $-\epsilon_i - \epsilon_j, i < j$ $E_{\alpha_i, j} + E_{\alpha_j, i}$
- $2\epsilon_i$ E_{i, α_i}
- $-2\epsilon_i$ $-E_{\alpha_i, i}$

(iii) The choice of positive roots is $\{\epsilon_i - \epsilon_j, i < j\} \cup \{\epsilon_i + \epsilon_j, i < j\} \cup \{2\epsilon_i\}$.

(iv) The simple roots corresponding to (iii) are

$$\epsilon_1 - \epsilon_2, \epsilon_2 - \epsilon_3, \dots, \epsilon_{n-1} - \epsilon_n, 2\epsilon_n$$



Dynkin diagram.

The fundamental weights (dual to the coroots! which are $H_1 - H_2, H_2 - H_3, \dots, H_{n-1} - H_n, \frac{1}{2}H_n$) are:

~~$\omega_1 = \epsilon_1$~~ $\omega_1 = \epsilon_1, \omega_2 = \epsilon_1 + \epsilon_2, \dots, \omega_n = \epsilon_1 + \dots + \epsilon_n.$

(v) The ~~roots~~ ^{weights} of V are $\epsilon_1, \dots, \epsilon_n, -\epsilon_1, \dots, -\epsilon_n$ with weight vectors: $e_1, \dots, e_n, e_{n+1}, \dots, e_{2n}$, respectively.

The highest weight is $\epsilon_1 = \omega_1$. It's easy to see that V is irreducible, hence $V = L(\omega_1)$.

(One elegant way to see irreducibility is to realize that the Weyl group acts by permutations and sign changes in the usual coordinates; hence the weights form a single W -orbit, so the representation must be irreducible.)

(vi) $\Lambda^2 V$ has weight vectors $e_k \wedge e_l$ with $k < l \leq 2n$ and multiplicity 1.
 weights ~~$\epsilon_i + \epsilon_j$~~ $\epsilon_i + \epsilon_j, 1 \leq i < j \leq n$
 $\epsilon_i - \epsilon_j, 1 \leq i \neq j \leq n.$
 $-(\epsilon_i + \epsilon_j), 1 \leq i < j \leq n.$
 0
 n

check: $4 \binom{n}{2} + n = 4 \cdot \frac{n(n-1)}{2} + n = 2n(n-1) + n = n(2n-1) = \binom{2n}{2}$ OK ✓

The vector $e_1 \wedge e_2$ is a highest weight vector with weight $\epsilon_1 + \epsilon_2$ $\Rightarrow \Lambda^2 V$ contains $L(\omega_2)$

Using the W -action, we can see that all the roots in \otimes are in the W -orbit of ω_2 hence the only weights that may not be in $L(\omega_2)$ are 0-weights.

Using for example the Weyl dim formula, we may find of $L(\omega_2)$
 $\dim L(\omega_2) = (n-1)(2n+1)$ (with $P = (n, n-1, \dots, 2, 1)$)

$\Rightarrow \dim \Lambda^2 V - \dim L(\omega_2) = 1$ and only the trivial has $\dim = 1$

The trivial rep in $\Lambda^2 V$ is spanned by $e_1 \wedge e_{n+1} + e_2 \wedge e_{n+2} + \dots + e_n \wedge e_{2n}$

(vii) The weights of $W = L(\omega_2)$ are:

$$\epsilon_1 + \epsilon_2, \epsilon_1 - \epsilon_2, \cancel{\epsilon_1 + \epsilon_2}, 0, -\epsilon_1 + \epsilon_2, -\epsilon_1 - \epsilon_2$$

all with multiplicity 1 in $sp(4)$.

Easy to verify Weyl's dimension.

(viii) Since $\epsilon_1 + \dots + \epsilon_k = \omega_k$ is a highest weight of $\Lambda^k V$ it follows that $L(\omega_k)$ is a submodule of $\Lambda^k V$.

$\phi_k: \Lambda^k V \rightarrow \Lambda^{k-2} V$. Notice that $\Lambda^{k-2} V$ does not have the weight ω_k , hence $L(\omega_k) \subseteq \ker \phi_k$. So it is sufficient to show $\ker \phi_k$ is irreducible. (One has to check that ϕ_k is a \mathfrak{g} -hom. first to know that $\ker \phi_k$ is a \mathfrak{g} -module.)

One way to prove this is via restriction to $sl(n) \subset sp(2n)$

$$\left\{ \begin{matrix} A & 0 \\ 0 & -A^t \end{matrix} \right\}$$

as in Fulton-Harris, see attached.

As before, we denote this by Γ_{a_1, \dots, a_n} :

$$\Gamma_{a_1, \dots, a_n} = \Gamma_{a_1 L_1 + a_2(L_1 + L_2) + \dots + a_n(L_1 + \dots + L_n)}.$$

These exhaust all irreducible representations of $\mathfrak{sp}_{2n}\mathbb{C}$.

We can find the irreducible representation $V^{(k)} = \Gamma_{0, \dots, 1, \dots, 0}$ with highest weight $L_1 + \dots + L_k$ easily enough. Clearly, it will be contained in the k th exterior power $\wedge^k V$ of the standard representation. Moreover, we have a natural contraction map

$$\varphi_k: \wedge^k V \rightarrow \wedge^{k-2} V$$

defined by

$$\varphi_k(v_1 \wedge \dots \wedge v_k) = \sum_{\substack{I \subset \{1, \dots, k\} \\ |I|=2}} Q(v_i, v_j) (-1)^{j+i-1} v_1 \wedge \dots \wedge \hat{v}_i \wedge \dots \wedge \hat{v}_j \wedge \dots \wedge v_k$$

(see §B.3 of Appendix B for an intrinsic definition and explanation). Since the representation $\wedge^{k-2} V$ does not have the weight $L_1 + \dots + L_k$, the irreducible representation with this highest weight will have to be contained in the kernel of this map. We claim now that conversely

Theorem 17.5. For $1 \leq k \leq n$, the kernel of the map φ_k is exactly the irreducible representation $V^{(k)} = \Gamma_{0, \dots, 0, 1, 0, \dots, 0}$ with highest weight $L_1 + \dots + L_k$.

PROOF. Clearly, it is enough to show that the kernel of φ_k is an irreducible representation of $\mathfrak{sp}_{2n}\mathbb{C}$. We will do this by restricting to a subalgebra of $\mathfrak{sp}_{2n}\mathbb{C}$ isomorphic to $\mathfrak{sl}_n\mathbb{C}$, and using what we have learned about representations of $\mathfrak{sl}_n\mathbb{C}$.

To describe this copy of $\mathfrak{sl}_n\mathbb{C}$ inside $\mathfrak{sp}_{2n}\mathbb{C}$, consider the subgroup $G \subset \mathfrak{Sp}_{2n}\mathbb{C}$ of transformations of the space $V = \mathbb{C}^{2n}$ preserving the skew form Q introduced in Lecture 16 and preserving as well the decomposition $V = \mathbb{C}\{e_1, \dots, e_n\} \oplus \mathbb{C}\{e_{n+1}, \dots, e_{2n}\}$. These can act arbitrarily on the first factor, as long as they do the opposite on the second, in coordinates, they are the matrices

$$G = \left\{ \begin{pmatrix} X & 0 \\ 0 & -X^{-1} \end{pmatrix}, X \in \text{GL}_n\mathbb{C} \right\}.$$

We have, correspondingly, a subalgebra

$$\mathfrak{s} = \left\{ \begin{pmatrix} A & 0 \\ 0 & -A \end{pmatrix}, A \in \mathfrak{sl}_n\mathbb{C} \right\} \subset \mathfrak{sp}_{2n}\mathbb{C}$$

isomorphic to $\mathfrak{sl}_n\mathbb{C}$.

Now, denote by W the standard representation of $\mathfrak{sl}_n\mathbb{C}$. The restriction of the representation V of $\mathfrak{sp}_{2n}\mathbb{C}$ to the subalgebra \mathfrak{s} then splits

$$V = W \oplus W^*$$

into a direct sum of W and its dual, and we have, correspondingly,

$$\wedge^k V = \bigoplus_{a+b=k} (\wedge^a W \otimes \wedge^b W^*),$$

How does the tensor product $\wedge^a W \otimes \wedge^b W^*$ decompose as a representation of $\mathfrak{sl}_n\mathbb{C}$? We know the answer to this from the discussion in Lecture 15 (see Exercise 15.30): we have contraction maps

$$\Psi_{a,b}: \wedge^a W \otimes \wedge^b W^* \rightarrow \wedge^{a-1} W \otimes \wedge^{b-1} W^*,$$

and the kernel of $\Psi_{a,b}$ is the irreducible representation $W^{(a,b)} = \Gamma_{0, \dots, 0, 1, 0, \dots, 0, 1, 0, \dots}$ with (if, say, $a \leq n - b$) highest weight $2L_1 + \dots + 2L_a + L_{a+1} + \dots + L_{n-b}$. The restriction of $\wedge^k V$ to \mathfrak{s} is thus given by

$$\wedge^k V = \bigoplus_{\substack{a+b \leq k \\ a+b \equiv k(2)}} W^{(a,b)}$$

and by the same token,

$$\text{Ker}(\varphi_k) = \bigoplus_{a+b=k} W^{(a,b)}.$$

Note that the actual highest weight factor in the summand $W^{(a,b)} \subset \text{Ker}(\varphi_k) \subset \wedge^k V$ is the vector

$$\begin{aligned} w^{(a,b)} &= e_1 \wedge \dots \wedge e_a \wedge e_{2n-b+1} \wedge \dots \wedge e_{2n} \\ &= e_1 \wedge \dots \wedge e_a \wedge e_{2n-k+a+1} \wedge \dots \wedge e_{2n}. \end{aligned}$$

Exercise 17.6. Show that more generally the highest weight vector in any summand $W^{(a,b)} \subset \wedge^k V$ is the vector

$$\begin{aligned} w^{(a,b)} &= e_1 \wedge \dots \wedge e_a \wedge e_{2n-k+a+1} \wedge \dots \wedge e_{2n} \wedge Q^{(k-a-b)/2} \\ &= e_1 \wedge \dots \wedge e_a \wedge e_{2n-k+a+1} \wedge \dots \wedge e_{2n} \wedge \left(\sum (e_i \wedge e_{n+i}) \right)^{(k-a-b)/2}. \end{aligned}$$

By the above, any subspace of $\text{Ker}(\varphi_k)$ invariant under $\mathfrak{sp}_{2n}\mathbb{C}$ must be a direct sum, over a subset of pairs (a, b) with $a + b = k$, of subspaces $W^{(a,b)}$. But now (supposing for the moment that $k < n$) we observe that the element

$$Z_{a,n-b} = E_{2n-b,a} + E_{n+a,n-b} \in \mathfrak{sp}_{2n}\mathbb{C}$$

carries the vector $w^{(a,b)}$ into $w^{(a-1,b+1)}$ and, likewise,

$$Y_{a+1,n-b+1} = E_{a+1,2n-b+1} + E_{n-b+1,n+a+1} \in \mathfrak{sp}_{2n}\mathbb{C}$$

carries $w^{(a,b)}$ to $w^{(a+1,b-1)}$. In case $a + b = k = n$, we see similarly that

$$Y_a = E_{n+a,a} \in \mathfrak{sp}_{2n}\mathbb{C}$$

carries the vector $w^{(a,b)}$ into $w^{(a-1,b+1)}$, and

$$U_{a+1} = E_{a+1,n+a+1} \in \mathfrak{sp}_{2n}\mathbb{C}$$

carries $w^{(a,b)}$ to $w^{(a+1,b-1)}$. Thus, any representation of $\mathfrak{sp}_{2n}\mathbb{C}$ contained in $\text{Ker}(\varphi_k)$ and containing any one of the factors $W^{(a,b)}$ will contain them all, and we are done. \square