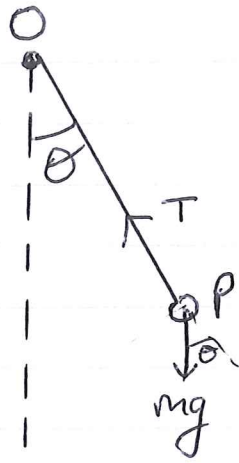


SWING

Aim & Questions

Simplifying Assumptions

- (A1) The "swing-rider" system is approximated by a pendulum with variable length $l(t)$.
- (A2) The rope is taken to be massless, & the rider is represented by a particle of mass m at the end of the rope (standing will move the mass up, hence $l(t)$ becomes smaller).
- (A3) Air resistance is the only other force (as well as tension in the rope) - i.e. we ignore friction at the pivot.
- (A4) System is modelled in 2D.
- (A5) Rope is modelled as a rigid rod (light rod)

Figure 5.1.

Equations of Motion (Ignore Air Resistance) Assume $l(t)$ fixed

Method 1 : Newton's II Law.

Use (r, θ) coords. So P has coords $l(t) \underline{e}_r = L \underline{e}_r$

[Recall $\underline{e}_r = \sin \theta \underline{i} + \cos \theta \underline{j}$
 $\underline{e}_\theta = +\cos \theta \underline{i} - \sin \theta \underline{j}$
 $\dot{\underline{e}}_r = \dot{\theta} \underline{e}_\theta, \dot{\underline{e}}_\theta = -\dot{\theta} \underline{e}_r$]

$$\underline{F} = m \underline{a} \Rightarrow (mg \cos \theta - T) \underline{e}_r - mg \sin \theta \underline{e}_\theta = m \underline{\ddot{r}} \quad (1.1)$$

$$\underline{\dot{r}} = L \dot{\theta} \underline{e}_\theta$$

$$\underline{\ddot{r}} = L \ddot{\theta} \underline{e}_\theta + L (-\dot{\theta})^2 \underline{e}_r$$

$$\therefore mg \cos \theta - T = -mL \dot{\theta}^2 \quad (1.2)$$

$$\underline{-mg \sin \theta = mL \ddot{\theta}} \quad \left[\text{or } \frac{d}{dt}(L^2 \dot{\theta}) = -gL \sin \theta \right] \quad (1.3)$$

Suppose $L = l(t)$ (not fixed):

$$\underline{\dot{r}} = \dot{l} \underline{e}_r + l \dot{\theta} \underline{e}_\theta$$

$$\underline{\ddot{r}} = \ddot{l} \underline{e}_r + \dot{l} \dot{\theta} \underline{e}_\theta + \dot{l} \dot{\theta} \underline{e}_\theta + l \ddot{\theta} \underline{e}_\theta + l \dot{\theta}^2 (-\underline{e}_r) + \dot{l} \dot{\theta} \underline{e}_\theta$$

$$\therefore \text{From (1.1): } mg \cos \theta - T = m \ddot{l} - m l \dot{\theta}^2 \quad (1.4)$$

$$\underline{-mg \sin \theta = \dot{l} \dot{\theta} + l \ddot{\theta}} \quad (1.5)$$

$$(1.5) \Rightarrow \frac{d}{dt}(l^2 \dot{\theta}) = -gl \sin \theta$$

Method 2: Ang Mom.

$$(1.1) \Rightarrow m \ddot{\underline{r}} = -mg \sin \theta \underline{e}_\theta + (mg \cos \theta - T) \underline{e}_r$$

$$\underline{r} \wedge \ddot{\underline{r}} = -r g \sin \theta \underline{e}_r \wedge \underline{e}_\theta + 0 \quad (1.6)$$

LHS of (1.6) is $\frac{d}{dt} (\underline{r} \wedge \dot{\underline{r}})$, so (1.6) becomes

$$\frac{d}{dt} (l^2 \dot{\theta}) \underline{e}_r \wedge \underline{e}_\theta = -g \sin \theta l \underline{e}_r \wedge \underline{e}_\theta,$$

$$\text{ie } \frac{d}{dt} (l^2 \dot{\theta}) = -g l \sin \theta \quad (1.7)$$

(rate of change of any mom) = torque (moment of external forces)

$$\left[\begin{array}{l} \underline{F} = m \ddot{\underline{r}} \\ \underline{r} \wedge \underline{F} = m \underline{r} \wedge \ddot{\underline{r}} = \frac{d}{dt} (\underline{r} \wedge m \dot{\underline{r}}) = \underline{r} \wedge \underline{F} \end{array} \right]$$

Method 3: Conservation of Energy

$$\underline{F} = m\underline{\ddot{r}}$$

$$\Rightarrow \underline{F} \cdot \underline{\dot{r}} = m\underline{\ddot{r}} \cdot \underline{\dot{r}} = \frac{d}{dt} \left(\frac{m\underline{\dot{r}}^2}{2} \right)$$

$$\therefore \text{Here, we have } \frac{md}{2dt} (\dot{l}^2 + l^2 \dot{\theta}^2) = (mg \cos \theta - T) \dot{l} - mg \sin \theta l \dot{\theta} \quad (1.8)$$

N.B. If $\dot{l} = 0$ then we have $\frac{d}{dt} \left(\frac{ml^2 \dot{\theta}^2}{2} \right) = \frac{d}{dt} (mgl \cos \theta)$

So, we have the unknown $mg \cos \theta - T$.

But (1.4) $\Rightarrow$ this can be replaced by $m\ddot{l} - ml\dot{\theta}^2$

$\therefore$ (1.8) becomes

$$\frac{1}{2} \frac{d}{dt} (\dot{l}^2 + l^2 \dot{\theta}^2) = l\ddot{l} - l\dot{\theta}^2 - mg \sin \theta l \dot{\theta}$$

ie $\cancel{\dot{l}^2} + \cancel{\frac{2}{2} l \dot{l} \dot{\theta}^2} + \frac{2}{2} l^2 \dot{\theta} \ddot{\theta} = \cancel{\dot{l}^2} - \cancel{l \dot{\theta}^2} - mg \sin \theta l \dot{\theta}$

$$2l\dot{l} \dot{\theta}^2 + l^2 \dot{\theta} \ddot{\theta} = -g \sin \theta l \dot{\theta}$$

$$2l\ddot{l} + l^2 \ddot{\theta} = -g \sin \theta$$

$$\text{ie } \frac{d}{dt} (l^2 \dot{\theta}) = -gl \sin \theta \quad (1.9)$$

So, we eventually have $\frac{d}{dt}(l^2 \dot{\theta}) = -gl \sin \theta$

Non-dimensionalise

$$\text{Let } \hat{l} = \frac{l}{L_0}, \quad \tau = \frac{t}{T_0} \quad L_0, T_0 \text{ fixed.}$$

$$\frac{d}{dt} = \frac{d}{d\tau} \frac{1}{T_0}$$

$$\therefore \text{ we have } \frac{L_0^2}{T_0^2} \frac{d}{d\tau} (\hat{l}^2 \dot{\theta}) = -g L_0 \hat{l} \sin \theta \quad \equiv \frac{d}{d\tau}$$

$$\Rightarrow \frac{d}{d\tau} (\hat{l}^2 \dot{\theta}) = -\frac{g T_0^2}{L_0} \hat{l} \sin \theta$$

I.C. $\frac{L(t)}{L_0} = \hat{l}$, say, so set $L_0 = L$ so that $\hat{l}(0) = 1$.

$$\therefore \text{ Set } T_0^2 = \frac{L_0}{g} \quad \left(\text{check } [time]^2 = \frac{[length]}{[length]/[time]^2} \right)$$

$$\text{Dwp } ^\wedge: \quad \frac{d}{d\tau} (\hat{l}^2 \dot{\theta}) = -\hat{l} \sin \theta \quad (1.10)$$

Phase Plane Analysis

Let us assume that l is const $\therefore l=1$.

So (1.10) becomes $\ddot{\Theta} = -\sin \Theta$ (1.11)

Look for steady states.

$$\dot{\Theta} = \ddot{\Theta} = 0 \Rightarrow \sin \Theta_0 = 0 \Rightarrow \Theta_0 = n\pi, \quad n=0,1,2, \dots \text{ etc.}$$

Recast (1.11) as a system of two 1st order ODEs:

$$\begin{aligned} \dot{\Theta} &= \psi \\ \dot{\psi} &= -\sin \Theta \end{aligned} \quad (1.12)$$

st.st. (Θ, ψ) are $(n\pi, 0)$ $n=0,1,2, \dots$

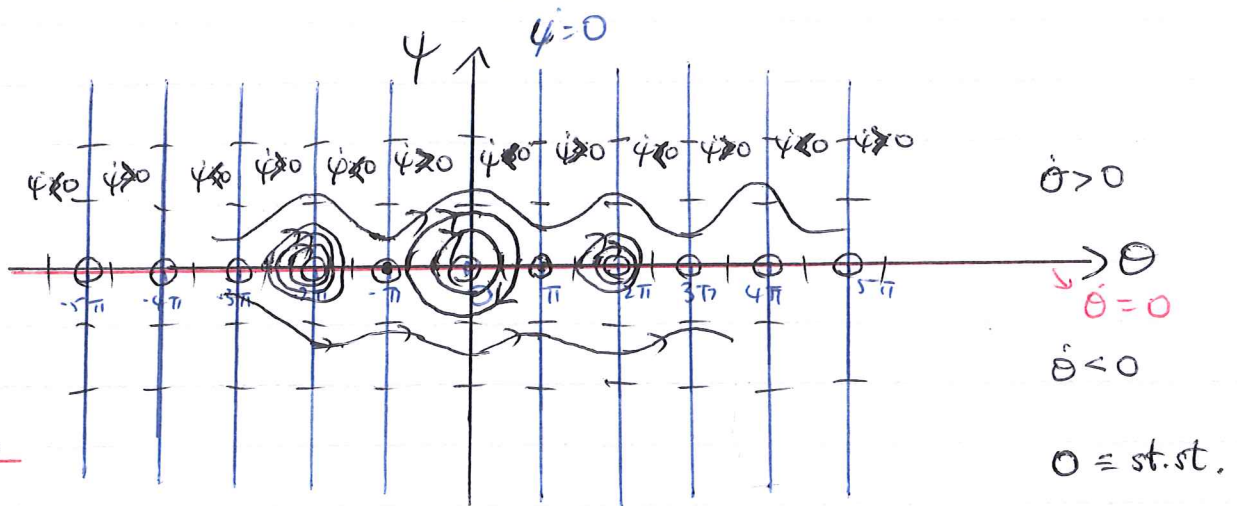


Figure S2

Nullclines: $\dot{\Theta} = 0 \Rightarrow \psi = 0$

Linear Stability Analysis: st.st $(2n\pi, 0)$: $\Theta = \Theta^* + \tilde{\Theta}$
 $\psi = \tilde{\psi}$

$$\begin{aligned} \tilde{\Theta}, \tilde{\psi} \text{ small: } \quad \dot{\tilde{\Theta}} &= \tilde{\psi} \\ \dot{\tilde{\psi}} &= -\tilde{\Theta} \end{aligned}$$

Jacobian is $\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$

$$\Rightarrow \begin{vmatrix} -\lambda & 1 \\ -1 & -\lambda \end{vmatrix} = 0$$

$$\Rightarrow \lambda^2 + 1 = 0 \Rightarrow \lambda^2 = -1 \Rightarrow \text{centre}$$

s.t.s. $((2n+1)\pi, 0) : \begin{aligned} \dot{\tilde{\theta}} &= \tilde{\psi} \\ \dot{\tilde{\psi}} &= -\tilde{\theta} \end{aligned}$

$$\Rightarrow \lambda^2 - 1 = 0 \therefore \lambda = \pm 1 \Rightarrow \text{saddle}$$

What do these results mean?

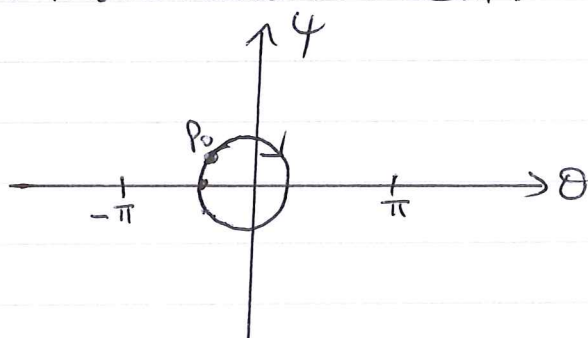


Figure S3

Int. Cond. $\begin{aligned} \theta &= \theta_0 \\ \psi &= \psi_0 \end{aligned}$

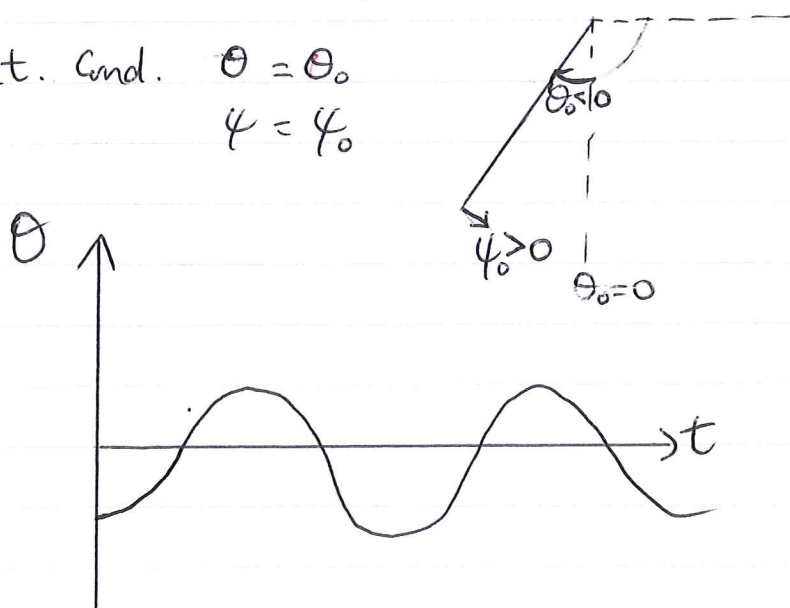


Figure S4

Is this realistic?