Tuesday, 19 May 2020 16:25 VECOUS flow, Wed ZOMay 10am car everyone see this? 2017 Q b, end  $x | b, e n x$ <br> $w = x (\frac{x^2}{x^2} + \frac{y^2}{b^2}) + \beta$ We need to find x and  $\frac{1}{2}$ , so that<br> $w = -U$  on  $\frac{z^2}{a^2} + \frac{y^2}{b^2} = 1$ , and  $\nabla^{2}w = -\frac{1}{\mu}G.$ 

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2016 \text{ Q1}
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2018 \text{ Q2}
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Tuesday, 19 May 2020 16:31<br>2017, Q | (iV)  $w = \alpha \left(\frac{x^2}{\alpha^2} + \frac{y^2}{\alpha^2}\right) + \beta$ Flux  $Q = \iint w \ dxdy$  $\frac{x^2}{a^2} + \frac{y^2}{b^2} \le 1$ , put  $\begin{cases} x = x \cos \theta \\ y = b \cos \theta \end{cases}$  $Q = \int_{a}^{2\pi} d\theta \int_{\infty}^{l} dr$  rab W  $=\int_{0}^{2\pi}d\theta\int_{c}^{1}dr$  rab  $(\propto(r^{2}cos^{2}\theta)$ =  $\int_{0}^{2\pi} d\theta \int_{c}^{1} dr$  rab  $(ar^{2}+ \beta)^{+\beta}$  $=2\pi ab \left[\frac{1}{4}\alpha r^{4}+\frac{1}{2}\beta r^{2}\right]_{0}^{1}$  $=2\pi a b \left(\frac{\alpha}{4}+\frac{\beta}{2}\right)$ The flux  $Q=0$  when  $\alpha=-215$ hetting in  $x = \frac{Ge^{2}b^{2}}{2\mu (a^{2}+b^{2})}$  $\beta = -U - \frac{1}{2\mu(a^2+b^2)}$ from (iii) gives<br>  $S = -0 - \frac{1}{2\mu}$ <br>  $= -\frac{1}{2\mu}$ 

 $2016$  Q1 (iii) Solve  $\frac{3u}{dt} = \cos t + \frac{2^{2}u}{3}$ with  $u = 0$  on  $y = 0$ ,  $u \rightarrow$ surt as<br>Put  $u = \text{sat} + w(y,t)$   $y \neq \infty$ Put  $u = sat + W(y,t)$ so  $\frac{\partial w}{\partial t} = \frac{\partial^2 w}{\partial y^2}$  with  $w = -\frac{2}{\partial x}$ on y=0  $Try$   $W(y,t)$ <br>=  $f(y) e^{it}$  as  $y \neq \infty$ <br>=  $f(y) e^{it}$  and  $x = y$ and take real ports<br> $\frac{\partial w}{\partial t} = i f(s) e^{it}$ inplicitly.  $\frac{\partial E}{\partial y^2} = f''(y) e^{it}$  $\frac{2}{3}$  and  $\frac{1}{3}$  we hat  $w = ie^{it} = -sint + \frac{1}{2}cos t$ on  $y = 0$  $\frac{\partial w}{\partial t} = \frac{\partial^2 w}{\partial y^2}$ and  $f(0) = 0$  $\Rightarrow$  if  $=$  f"  $f \rightarrow 0$  as  $y \rightarrow \infty$  $(\tau_{\mu} - \frac{f(y)}{f(z)}) = e^{\lambda y} \frac{f(z)}{f(z)} = e^{\frac{\lambda y}{\lambda}} = \pm e^{\frac{\lambda y}{\lambda}}$ For decay,  $670$  as  $y \ge 10$ ,  $w = 10$ ,  $w = 10$  $f(y) = f(0) e^{-\frac{1+i}{\sqrt{2}}y} = -\frac{1+i}{\sqrt{2}}$ <br> $f(y) = f(0) e^{-\frac{1+i}{\sqrt{2}}y}$  $U = \frac{1}{2}e^{-\frac{1}{2}te^{-t}}$  $w(y,t) = f(y) e^{i t}$ <br>=  $\overline{c} e^{i t} - \frac{1+i}{\sqrt{2}} y$ <br>=  $e^{-y} \sqrt{2} \overline{c} e^{i (t-\frac{y}{\sqrt{2}})}$ The real part  $\zeta$ <br>Wreal =  $e^{-\frac{3}{2}\sqrt{z}}$  sur  $(\frac{y}{\sqrt{z}}-t)$ Firally, nally,<br> $u = \sin t + e^{-\frac{y}{2}}$  sur  $\left(\frac{y}{\sqrt{2}} - t\right)$ 



Wednesday, 20 May 2020 10:40<br>20 8 Q 2 (c)  $F'' + 4F + \frac{F^2}{2} + C = 0$ with  $F=0$  on  $\theta=\pm\alpha$  $F(0) = F(0) \phi(1)$   $\theta = \alpha y$  $F'' = F(c) \phi''(\psi) \left(\frac{d\psi}{d\theta}\right)^2 = \frac{F(c)}{d^2} \phi''$  $\notimes \frac{F(0)}{\alpha^2}$   $\not=$   $\frac{1}{4}$   $\left(1 + 4F(0)\psi + \frac{F(0)^2}{v}\psi + C = 0\right)$ wer BC  $\phi = 0$  on  $y = \pm 1$ Structurally, this is  $\phi'' = g(\phi)$ Multiply  $\circledast$  by  $\phi'$  and  $\int d\psi$  $rac{F(0)}{x^{2}}$   $\phi'' \phi'$   $\in$  4 $F(0)$   $\phi \phi' + \frac{F(0)^{2}}{v}$   $\phi \phi'$  $+c=0$  $\frac{F(0)}{2\alpha^{2}}\phi'^{2}+\frac{4}{2}F(0)\phi^{2}+\frac{F(0)^{2}}{5\upsilon}\phi^{3}+C\phi^{4}$  $A = A$ The symmetre solution The symmetric soutient  $\phi'(0) = 0$ and we hnow  $\hat{\varphi}(0) = 1$  $2F(0) + \frac{F(0)^2}{30} + C = A$  $rac{F^{(0)}}{2\alpha^{2}}\phi'^{2}+2F^{(0)}\phi^{2}+\frac{F^{(0)}^{2}}{3\nu}\phi^{3}+C\phi^{2}$  $ZF(c) + \frac{F(c)^{2}}{3c} + C$  $rac{F(0)}{2\alpha^{2}}\phi'^{2} = 2F(0)(1-\phi^{2}) + \frac{F(0)^{2}}{3D}(1-\phi^{3})$  $+C(1-\phi)$  $2\frac{1}{x^{2}}\psi'^{2} = 2(1-\psi^{2}) + \frac{F(0)}{3v}(1-\psi^{2})$  $+\frac{c}{F(c)}(1-e)$  $= (1-f)(\frac{c}{F(0)} + 2(1+f) + \frac{F(0)}{32}(1+f))$  $=$   $G(\phi)$  $\frac{1}{2}d^{2}(\frac{df}{dt})^{2}=(0)$  $\frac{1}{\sqrt{2\alpha}}\frac{d\alpha}{d\gamma} = \pm \sqrt{G(\ell)}$  $\int \mathbb{Z} \alpha \int d\psi = \pm \int \frac{\partial \psi}{\partial \psi}$  $JZ \propto \gamma = \mp \int_{\phi}^{1} \frac{ds}{\sqrt{G(s)}}$  $\overline{f}$  as  $\overline{\psi}$  is on the lower limit<br> $\overline{\psi}(\pm I) = 0$  and  $\overline{\psi}(0) = 1$ , so  $\frac{df}{d\psi} < 0$ <br> $y = f \frac{ds}{d\psi}$  in  $1 > 0$  $=520y = 6 \int_{\phi}^{1} \frac{ds}{\sqrt{G(s)}}$ =  $\int_{\phi}^{1} \frac{ds}{\sqrt{(1-s)(\frac{C}{F(c)}+2(1+s)+\frac{F(c)}{g(0)}(1+s+s))}}$ The guen around the pat of  $K$ <br>The guen around her  $R(1-s)s^2$ + lover powers of s. Adming for  $mmg$  for<br>  $\alpha y = \int_{\beta}^{1} \frac{ds}{\sqrt{\frac{z}{3}R(1-s)(K+(1+\frac{6}{R})s+s^{2})}}$  $\Rightarrow R = \frac{F(0)}{2}$  and  $K = |+6\frac{v}{F(0)} + \frac{zvc}{F(0)}$ <br>  $\Rightarrow R = \frac{F(0)}{2}$  and  $K = |+6\frac{v}{F(0)} + \frac{zvc}{F(0)}$ when  $y=1$ You can also see Iar Hewitt's<br>video version of this calculation. video version of this calculation.<br>It's easier in hondsight to schiptify<br>the expression for  $\left(\frac{d\ell}{dy}\right)^2$  to look  $\left(\frac{l\phi}{l\eta}\right)^{2}=\frac{2}{3}\alpha^{2}K(l\phi)(K+(l+\frac{C}{R}\phi)+\phi)$ <br>kefae separating varables.