

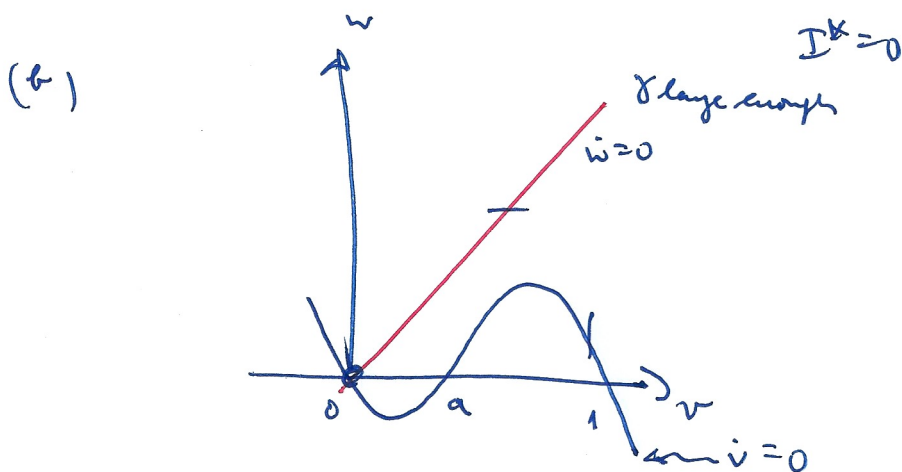
1. $\epsilon \dot{v} = I^* + f(v) - w$
 $\dot{w} = \gamma v - w$

$f = -v(a-v)(1-v) \quad 0 < a < 1$

(a) v intracellular electric potential
 w gate variable
 I^* applied current

$-f(v) + v$ ionic current at

Start with Hodgkin-Huxley, 3 gate variables ^{m, h, n} & telegraph equation for potential (via I_H & I_L): $\tau_m \ll 1 \Rightarrow$ in equilibrium, $n+h = \text{constant}$.
 \Rightarrow FN reduction ...



Endstate are equilibria $(0,0)$ for δ layer

Linearize $\begin{pmatrix} \dot{v} \\ \dot{w} \end{pmatrix} \approx \underbrace{\begin{pmatrix} \frac{f'}{\epsilon} & -\frac{1}{\epsilon} \\ \gamma & -1 \end{pmatrix}}_M \begin{pmatrix} v \\ w \end{pmatrix}$

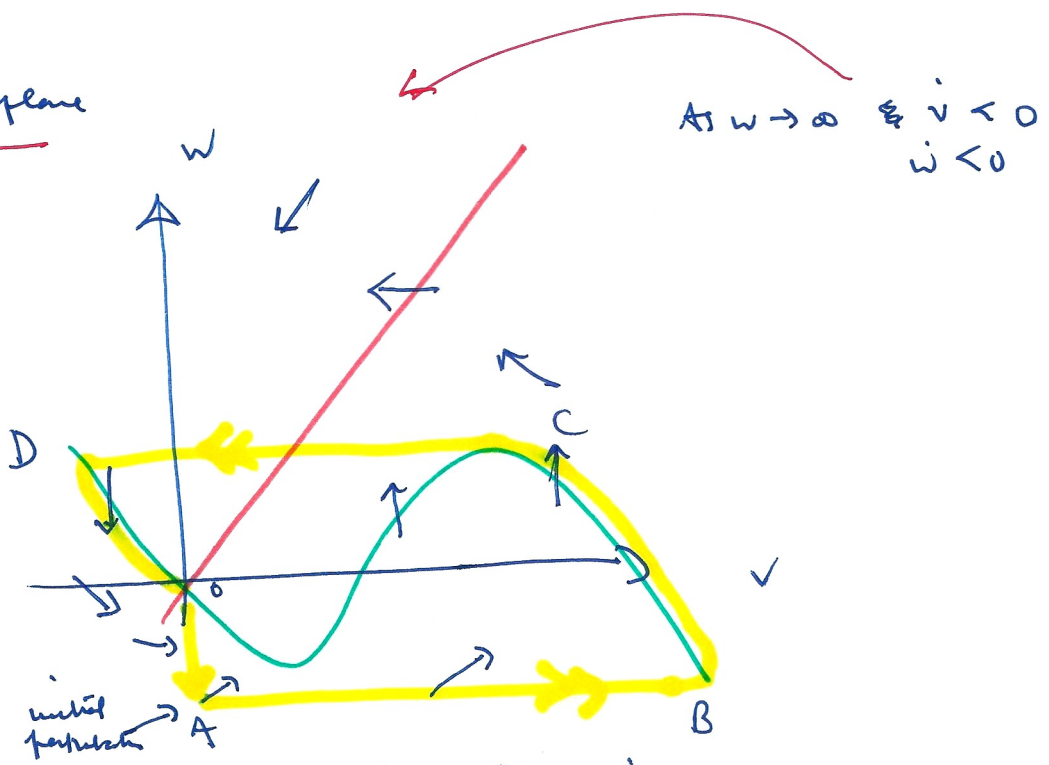
$\text{tr } M = \frac{f'}{\epsilon} - 1 < 0$ at $(0,0)$, $\det M = -\frac{f'}{\epsilon} + \gamma > 0 \Rightarrow$ stable

$t \sim 1$: $w \approx f(v)$

$\dot{w} = f(v) - w$

$t \sim \varepsilon$, $t = \varepsilon T$ $v' = f(v) - w$
 $w' = 0$

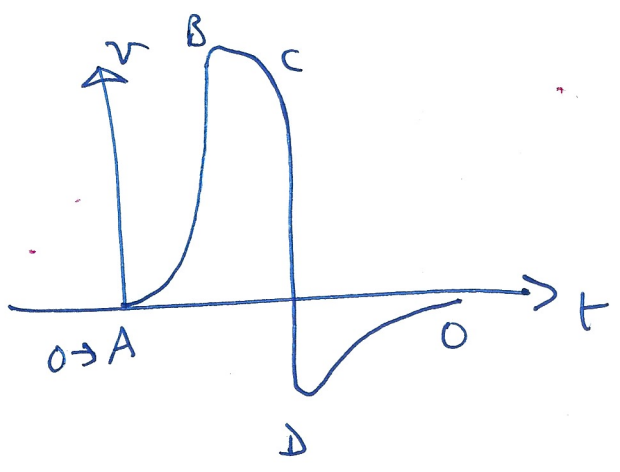
phase plane



trajectories cycle round stable origin

threshold : as w approaches I^* $\dot{v} > 0$ for short time \leftrightarrow lower

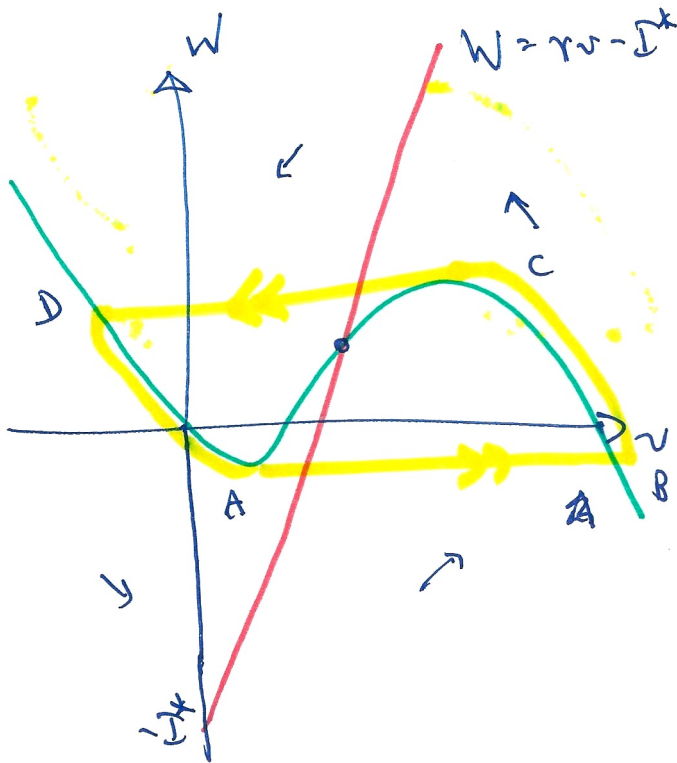
$w \Rightarrow$ trajectory as shown : fast phases with \rightarrow



$$(e) \quad \dot{I}^* = 0, \quad W = w - \dot{I}^*$$

$$\dot{v} = f(v) - W$$

$$\dot{W} = r v - \dot{I}^* - W$$



$$P \approx \int_B^C dt + \int_D^A dt$$

$$\approx \int_B^C \frac{dw}{\dot{w}} + \int_D^A \frac{dw}{\dot{w}} \quad \text{along } W = f(v)$$

$$= \int_B^C \frac{dw}{rv - \dot{I}^* - W} + \int_D^A \frac{dw}{rv - \dot{I}^* - W}$$

$$= \int_B^C \frac{f'(v) dv}{rv - \dot{I}^* - f(v)} + \int_D^A \frac{f'(v) dv}{rv - \dot{I}^* - W}$$

Answer on question leaves: as W, how odd

$$P \approx \int_{w_B}^{w_C} \frac{dw}{\gamma r - \mathcal{I}^* - w} + \int_{w_D}^{w_A} \frac{dw}{\gamma r - \mathcal{I}^* - w}$$

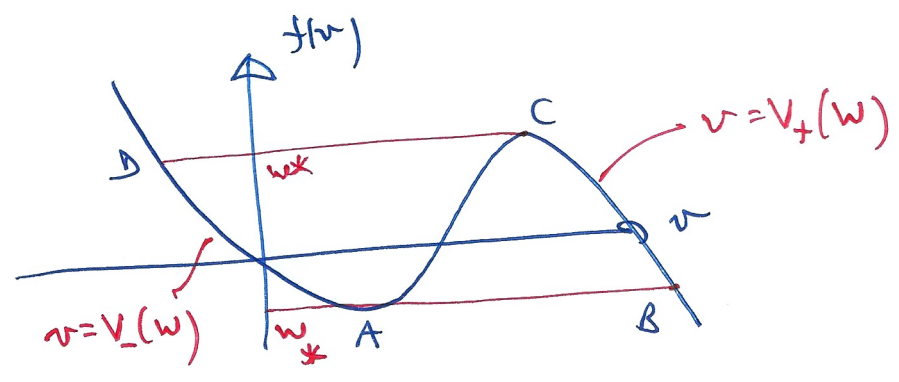
On BC $w < \gamma r - \mathcal{I}^* - w$, denote $v = V_+(w)$ on BC

on DA $w > \gamma r - \mathcal{I}^* - w$ denote $v = V_-(w)$ on AD

(see figure)

Also let $w = f(v) = w_*$ at A & B

$w = w^*$ at C & D



$$\text{So } P \approx \int_{w_*}^{w^*} \frac{dw}{\gamma V_+(w) - \mathcal{I}^* - w} + \int_{w^*}^{w_*} \frac{dw}{\gamma V_-(w) - \mathcal{I}^* - w}$$

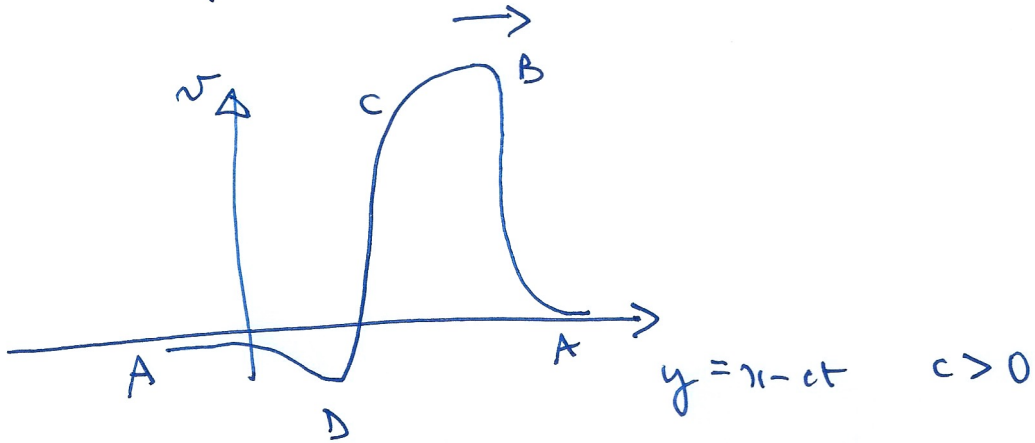
$$= \int_{w_*}^{w^*} \left[\frac{1}{\gamma V_+(w) - \mathcal{I}^* - w} - \frac{1}{\gamma V_-(w) - \mathcal{I}^* - w} \right] dw$$

note this is > 0

(d)

$$\varepsilon v_F = f(v) - w + \varepsilon^2 v_{xxx}$$

$$w_F = \gamma v - w$$



(i)
$$-\varepsilon c v' = f(v) - w + \varepsilon^2 v''$$

$$-c w' = \gamma v - w$$

$v(\eta), w(\eta)$
 bc's $v, w \rightarrow 0$ at $\pm \infty$.

(ii) AB part where $y = \varepsilon \eta$

$$\Rightarrow -c v' = f(v) - w + v''$$

$$-c w' = \varepsilon(\gamma v - w) \Rightarrow w \approx 0$$

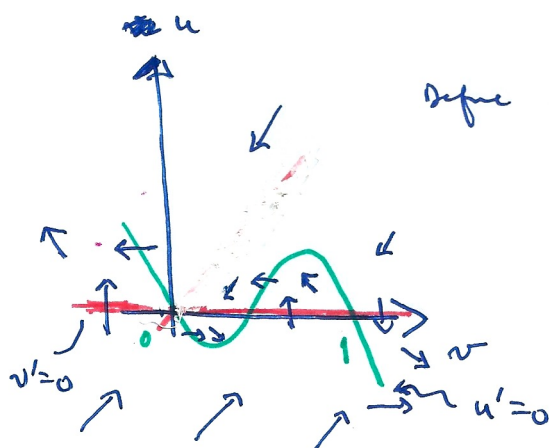
$$\Rightarrow v'' + c v' + f(v) = 0 \quad w \approx AB$$

$$\eta \rightarrow \infty \quad v \rightarrow 0$$

$$\eta \rightarrow -\infty \quad v \rightarrow v_B \quad (\text{where } f(v_B) = 0 \text{ and } v_B = 1)$$

ii $v \rightarrow 1$ (as a).

(iii)



define $u' = -u$
 then $u' = f(v) - cu$

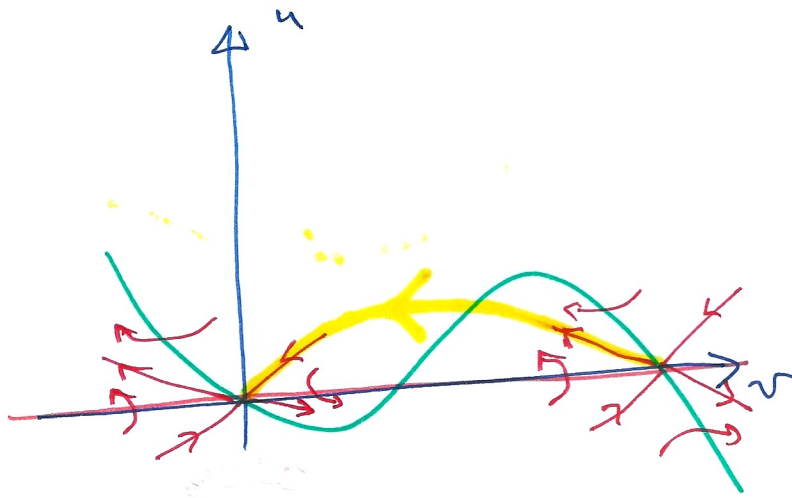
At large u $v' < 0$ $u' < 0$

\Rightarrow trajectories as shown

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Clearly $(0,0)$ & $(1,0)$ are saddles, $(\frac{1}{2},0)$ a node or spiral

To connect $(1,0)$ ($\gamma \rightarrow -\infty$) to $(0,0)$ ($\gamma \rightarrow \infty$) we need to connect the unstable separatrix from $(1,0)$ in $v < 1$ to the stable separatrix v_2 to $(0,0)$ in $v > 0$:



This requires a choice of c to make them join.

$$v'' + cv' + f(v) = 0$$

$$\text{So } \int_{-\infty}^{\infty} v'^2 dv + c \int_{-\infty}^{\infty} v'^2 dv + \int_1^0 f(v) dv = 0$$

$$\Rightarrow c = \frac{\int_0^1 f(v) dv}{\int_{-\infty}^{\infty} v'^2 dv}$$

which requires $\int_0^1 f(v) dv > 0$
 $v_2 \approx \frac{1}{2}$