

3/

$$\nabla \cdot \underline{u} = 0$$

$$\underline{u}_t + (\underline{u} \cdot \nabla) \underline{u} + \nabla \times \underline{u} = -\frac{1}{\rho_0} \nabla p + \frac{\rho}{\rho_0} \underline{g} + \frac{\mu}{\rho_0} \nabla^2 \underline{u}$$

$$T_t + \underline{u} \cdot \nabla T = \kappa \nabla^2 T$$

$$\underline{g} = g \underline{k} \quad \alpha = -\beta k$$

$$\rho = \rho_0 [1 - \alpha (T - T_0)]$$

(i) $x \sim d, t \sim \frac{d^2}{\kappa} \quad \underline{u} \sim \frac{\kappa}{d}, \rho + \rho_0 \delta \rho \sim \frac{\mu \underline{u}}{d^2}, T - T_0 \sim \Delta T$

$\Rightarrow$ non-d

$$\nabla \cdot \underline{u} = 0$$

$$\frac{\kappa}{d} \frac{\kappa}{d^2} \left[\frac{d\underline{u}}{dt} \right] + \frac{\kappa}{d} \beta \underline{k} \times \underline{u} = -\frac{\mu \kappa}{\rho_0 d^3} \nabla p + \alpha \Delta T g T \underline{k} + \frac{\mu}{\rho_0} \frac{\kappa}{d^3} \nabla^2 \underline{u}$$

$\div$ by $\frac{\mu \kappa}{\rho_0 d^3}$

$$\Rightarrow \frac{\rho_0 \kappa}{\mu} \frac{d\underline{u}}{dt} + \frac{\rho_0 d^2 \beta}{\mu} \underline{k} \times \underline{u} = -\nabla p + \frac{\alpha \Delta T \rho_0 d^3}{\mu \kappa} T \underline{k} + \nabla^2 \underline{u}$$

or $\frac{1}{Pr} \frac{d\underline{u}}{dt} + Ta \underline{k} \times \underline{u} = -\nabla p + Ra T \underline{k} + \nabla^2 \underline{u}$

$Pr = \frac{\mu}{\rho_0 \kappa}$ Prandtl $Ta = \frac{\rho_0 d^2 \beta}{\mu}$ Taylor $Ra = \frac{\alpha \rho_0 \Delta T d^3}{\mu \kappa}$ Rayleigh

Energy equation is

(2)

$$\tau_t + \underline{u} \cdot \nabla \tau = \nabla^2 \tau$$

(ii)

$$\underline{\tau} = 1 - z$$

$$\underline{u} = 0$$

$$p_z = Ra(1-z)$$

$$\Rightarrow \underline{p} = -\frac{1}{2} Ra(1-z)^2$$

(iii) Small $\underline{u}$ linear $\underline{p} = -\frac{1}{2} Ra(1-z)^2 + P$, $\tau = 1 - z + \theta$

$$\Rightarrow \frac{1}{Pr} \frac{\partial \underline{u}}{\partial t} + Ta \underline{k} \times \underline{u} = -\nabla P + Ra \underline{\theta} \underline{k} + \nabla^2 \underline{u}$$

$$\theta_t - w = \nabla^2 \theta$$

$$[\underline{u} = (u, v, w)] \quad \nabla \cdot \underline{u} = 0$$

[To derive the dispersion relation...

Note that if $\nabla \cdot \underline{u} = 0$ we can find $\underline{\Psi}$ s.t. $\underline{u} = \text{curl } \underline{\Psi}$

because solve $\nabla^2 \underline{\Psi} = -\underline{\omega} = -\text{curl } \underline{u}$ with appropriate bc's.

$$\Rightarrow \text{grad div } \underline{\Psi} - \text{curl curl } \underline{\Psi} = -\text{curl } \underline{u}$$

$$\text{Suppose } \text{div } \underline{\Psi} = A, \text{ then } \nabla A = \text{curl}(\text{curl } \underline{\Psi} = \underline{u})$$

$$\text{So } \nabla^2 A = 0,$$

$$\text{so solve } \nabla^2 \underline{\Psi} = -\underline{\omega} \text{ subject to } \nabla \cdot \underline{\Psi} = 0 \text{ on boundary } (eg at \infty)$$

$$= A = 0$$

$$\Rightarrow \underline{u} = \text{curl } \underline{\Psi}$$

etc.

]

(3)

Assume the characteristic equation for solutions of $e^{\sigma t} e^{ikx} \sin n\tilde{z}$

$$(\sigma + k_n^2) \left(\frac{\sigma}{Pr} + k_n^2 \right)^2 k_n^2 - Ra \left(\frac{\sigma}{Pr} + k_n^2 \right) k^2 + Ta (\sigma + k_n^2) n^2 \tilde{n}^2 = 0$$

$$k_n^2 = k^2 + n^2 \tilde{n}^2$$

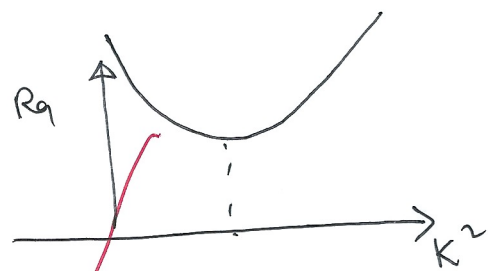
$$Ta = \frac{\rho_0 d^3 f}{\mu} \quad (f = 2\Omega \text{ notation})$$

If σ is real the critical Rayleigh number occurs when $\sigma = 0$

$$\Rightarrow Ra = \frac{k_n^6 + Ta n^2 \tilde{n}^2}{k^2}$$

$$= \frac{(k^2 + n^2 \tilde{n}^2)^3 + n^2 \tilde{n}^2 Ta}{k^2}$$

or $Ta \rightarrow \infty$? if $Ta \gg 1$ $Ra_c \approx \frac{n^2 \tilde{n}^2 Ta}{k^2}$ well no



where $n=1$ gives minimum

there is always a minimum

$$\frac{dRa}{dk^2} = \frac{3(k^2 + n^2 \tilde{n}^2)^2}{k^2} - \frac{\{(k^2 + n^2 \tilde{n}^2)^3 + n^2 \tilde{n}^2 Ta\}}{k^4}$$

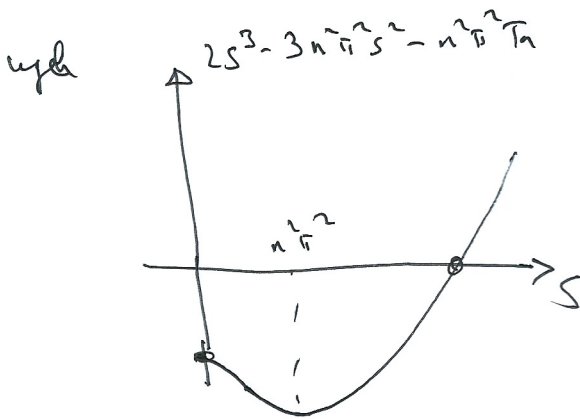
$$= 0 \text{ if } 3k^2(k^2 + n^2 \tilde{n}^2)^2 = (k^2 + n^2 \tilde{n}^2)^3 + n^2 \tilde{n}^2 Ta$$

$$\text{or } (k^2 + n^2 \tilde{n}^2) k^2$$

note $K_n^2 = k^2 + n^2 \bar{\eta}^2 = S$

so $3(S - n^2 \bar{\eta}^2) S^2 = S^3 + n^2 \bar{\eta}^2 T_a$

i $3S^3 - 3n^2 \bar{\eta}^2 S^2 - n^2 \bar{\eta}^2 T_a = 0$



turning point at

$$6S^2 = 6n^2 \bar{\eta}^2 S$$

$$S = 0 \text{ or } n^2 \bar{\eta}^2$$

When T_a is large root is $S \approx \left(\frac{n^2 \bar{\eta}^2 T_a}{2} \right)^{1/3}$

(other roots complex $S \approx \pm i \left(\frac{T_a}{3} \right)^{1/2} + \dots$)

so min R_a is at $S = \left(\frac{T_a n^2 \bar{\eta}^2}{2} \right)^{1/3}$

so $K \approx \sqrt{S}$ (take $n=1$)

$$\Rightarrow K \approx \bar{\eta}^{1/3} \left(\frac{T_a}{2} \right)^{1/6}$$

$$R_{ac} \approx \frac{K^6 + \bar{\eta}^2 T_a}{K^2} \approx \frac{\frac{\bar{\eta}^2 T_a}{2} + \bar{\eta}^2 T_a}{\left(\frac{T_a \bar{\eta}^2}{2} \right)^{1/3}}$$

$$\approx 3 \left(\frac{\bar{\eta}^2 T_a}{2} \right)^{2/3}$$

rotation makes it more stable
(as $R_{ac} \uparrow T_a$).