

**INTRODUCTION TO QFT 2019: PROBLEM SET 5
SOLUTIONS TO NON-STARRED QUESTIONS**

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1. *Diagrams and Feynman Rules* As the question hints this is best done by Mathematica. On the web-page is a simple Mathematica Notebook that does the job. The basic idea is that for a finite graph we don't really need functional differentiation. If we're just differentiating w.r.t a fixed number, N , of j 's we label them $\{j_n, n = 1, \dots, N\}$ and represent

$$\int d^4x_1 \int d^4x_2 j(x_1) D_F(x_1 - x_2) j(x_2) \quad (1)$$

formally by

$$\sum_{k,n} j_k D_{Fk,n} j_n \quad (2)$$

so the propagator is just a matrix *and* we can do ordinary partial differentiations which is built into Mathematica. This is good enough to generate a representation of the graphs in terms of propagators; at the end we simply replace $D_{Fk,n}$ by the real expression $D_F(x_k - x_n)$ and integrate d^4x_n over internal vertices.

2. *Generating Functionals I*

a) Adopt the short-hand

$$\frac{\delta}{\delta j_k} = -i \frac{\delta}{\delta j(x_k)} \quad (3)$$

then

$$\begin{aligned} & i \left(\prod_{k=1}^4 \left(-i \frac{\delta}{\delta j(x_k)} \right) \right) W[j] \quad (4) \\ &= \left(\prod_{k=1}^3 \left(\frac{\delta}{\delta j_k} \right) \right) \frac{1}{Z[j]} \frac{\delta Z[j]}{\delta j_4} \\ &= \left(\prod_{k=1}^2 \left(\frac{\delta}{\delta j_k} \right) \right) \left(\frac{1}{Z[j]} \frac{\delta^2 Z[j]}{\delta j_3 \delta j_4} - \frac{1}{Z[j]^2} \frac{\delta Z[j]}{\delta j_3} \frac{\delta Z[j]}{\delta j_4} \right) \\ &= \frac{\delta}{\delta j_1} \left(\frac{1}{Z[j]} \frac{\delta^3 Z[j]}{\delta j_2 \delta j_3 \delta j_4} - \frac{1}{Z[j]^2} \left(\frac{\delta^2 Z[j]}{\delta j_2 \delta j_3} \frac{\delta Z[j]}{\delta j_4} + \text{distinct perms} \right) \right. \\ &\quad \left. + \frac{2}{Z[j]^3} \frac{\delta Z[j]}{\delta j_2} \frac{\delta Z[j]}{\delta j_3} \frac{\delta Z[j]}{\delta j_4} \right) \\ &= \frac{1}{Z[j]} \frac{\delta^4 Z[j]}{\delta j_1 \delta j_2 \delta j_3 \delta j_4} - \frac{1}{Z[j]^2} \left(\frac{\delta^2 Z[j]}{\delta j_2 \delta j_3} \frac{\delta^2 Z[j]}{\delta j_4 \delta j_1} + \frac{\delta^3 Z[j]}{\delta j_1 \delta j_2 \delta j_3} \frac{\delta Z[j]}{\delta j_4} + \text{distinct perms} \right) \\ &\quad + \frac{2}{Z[j]^3} \left(\frac{\delta^2 Z[j]}{\delta j_1 \delta j_2} \frac{\delta Z[j]}{\delta j_3} \frac{\delta Z[j]}{\delta j_4} + \text{distinct perms} \right) - \frac{6}{Z[j]^4} \frac{\delta Z[j]}{\delta j_1} \frac{\delta Z[j]}{\delta j_2} \frac{\delta Z[j]}{\delta j_3} \frac{\delta Z[j]}{\delta j_4} \end{aligned}$$

Now set $j = 0$; all odd derivatives of $Z[j]$ vanish because it is an even function of j . We are left with

$$\frac{1}{Z[j]} \frac{\delta^4 Z[j]}{\delta j_1 \delta j_2 \delta j_3 \delta j_4} \Big|_{j=0} - \left(\frac{1}{Z[j]} \frac{\delta^2 Z[j]}{\delta j_2 \delta j_3} \Big|_{j=0} \frac{1}{Z[j]} \frac{\delta^2 Z[j]}{\delta j_4 \delta j_1} \Big|_{j=0} + \text{distinct perms} \right) \quad (5)$$

which is precisely the full 4pt-Green's function with the components that are products of two 2pt-functions subtracted off.

3. Generating Functionals II

a) We have from the lectures

$$Z[j] = \exp \left(-\frac{1}{2} \int d^4 x_1 \int d^4 x_2 j(x_1) D_F(x_1 - x_2) j(x_2) \right) \quad (6)$$

so

$$W(j) = \frac{i}{2} \int d^4 x_1 \int d^4 x_2 j(x_1) D_F(x_1 - x_2) j(x_2) \quad (7)$$

and

$$\phi(x) = \frac{\delta W}{\delta j(x_k)} = i \int d^4 x_1 D_F(x - x_1) j(x_1) \quad (8)$$

Now act on both sides with $Q = (-\partial_\mu \partial^\mu - m^2 + i\epsilon)$ to get (remember D_F is the Green's function of Q)

$$Q_x \phi(x) = i \int d^4 x_1 i \delta^4(x - x_1) j(x_1) = -j(x) \quad (9)$$

So

$$\begin{aligned} \Gamma[\phi] &= \frac{i}{2} \int d^4 x_1 \int d^4 x_2 (Q_{x_1} \phi(x_1)) D_F(x_1 - x_2) (Q_{x_2} \phi(x_2)) \\ &\quad + \int d^4 x \phi(x) Q_x \phi(x) \end{aligned} \quad (10)$$

Integrating the derivatives in Q_{x_1} by parts twice gives

$$\begin{aligned} \Gamma[\phi] &= \frac{i}{2} \int d^4 x_1 \int d^4 x_2 \phi(x_1) (Q_{x_1} D_F(x_1 - x_2)) (Q_{x_2} \phi(x_2)) \\ &\quad + \int d^4 x \phi(x) Q_x \phi(x) \\ &= \frac{1}{2} \int d^4 x \phi(x) (-\partial_\mu \partial^\mu - m^2 + i\epsilon) \phi(x) \end{aligned} \quad (11)$$

which of course is just the action for a free field.