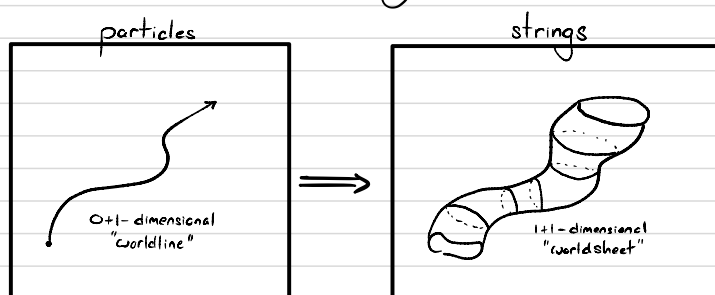


WHAT IS STRING THEORY?

String theory is the theory of fundamental, quantum-mechanical strings.



Perturbative string theory (first discovered, best developed formulation)⁺ is first-quantized*, S-matrix theory.

⁺: many post-1990 developments give non-perturbative insights (duality, matrix models, ...). May see some in STII.

*: string field theory is 2nd quantized (but still perturbative) framework. Still not well understood in general. Don't use it.

Key features: consistently incorporates gravity (so a theory of quantum gravity)

is "unique" theory (this is a subtle issue; relates to dualities)

incorporates many interesting / phenomenologically relevant ingredients from QFT/particle physics.

- non-Abelian gauge symmetries w/ chiral matter.
- spacetime supersymmetry
- "unification"

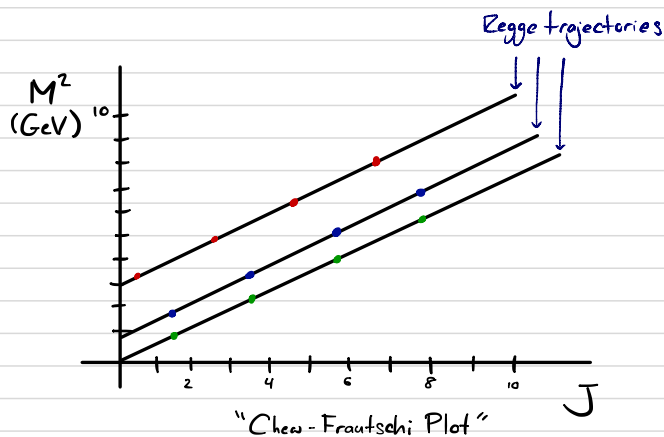
In this course, will develop bosonic string theory. This theory is missing some of the features mentioned above, and suffers from some serious defects / inconsistencies. However, it illustrates the key ideas / techniques in a relatively clean way.

If you continue on to STII, you will learn superstring theory. This is the more exciting theory that has been considered as a candidate to some day describe our world.

Historical Motivation

Late 1950s/early 1960s: Hadronic particle zoo.

Many experimentally observed resonances w/ increasingly large invariant mass/spin.



leading trajectory: $J = \alpha' M^2 + \alpha(0)$

$\alpha' \sim 1 (\text{GeV})^{-2}$

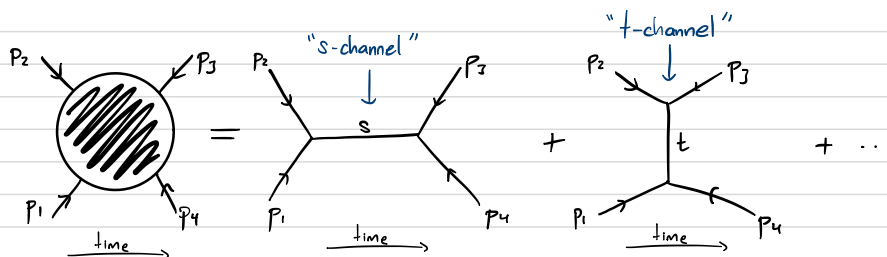
$\alpha(0)$: "Regge intercept"

lightest scalar: $M^2 = -\frac{\alpha(0)}{\alpha'}$

fixed $M^2 \rightarrow$ maximum spin $J(M^2)$

Problems: all resonances fundamental particles?
infinitely many?
can they be described using QFT?

UV difficulties for high spin particles:



$(-, +, +, +)$ signature $\left\{ \begin{array}{l} s = -(p_1 + p_2)^2 \quad (> 0 \text{ for physical elastic scattering}) \\ t = -(p_1 + p_4)^2 \quad (< 0 \text{ for physical elastic scattering}) \\ u = -(p_1 + p_3)^2 \quad (> 0 \text{ for physical elastic scattering}) \end{array} \right\} (1,2) \rightarrow (3,4)$

$$s + t + u = \sum_{i=1}^4 m_i^2$$

Consider t -channel exchange of spin- J particle: $\delta A(s, t) \sim \frac{-g_J^2 (-s)^J}{t - m_J^2}$ for fixed t , large (positive) s (Regge limit)

$$\delta \mathcal{L} \sim g_J (\phi^* \vec{\partial}_{\mu_1} \vec{\partial}_{\mu_2} \dots \vec{\partial}_{\mu_J} \phi) \sigma^{\mu_1 \dots \mu_J}$$

Historical Motivation

This "hard" UV behaviour is not what was observed in, say, pion scattering.

Even worse for loops:

$$\sim \int \frac{d^D p}{(2\pi)^D} \frac{p^{4J}}{(p^2)^4} \begin{cases} \text{safe for scalars} \\ \text{logarithmic divergent for spin } J=1 \\ \text{badly divergent for } J>1 \Rightarrow \text{not renormalizable} \end{cases}$$

(general difficulties w/ theories of elementary higher spin particles.)

For finite sum, term w/ largest spin exchanged in t-channel dominates in Regge limit.

$$\sum_{J=0}^{J_{\max}} \frac{-g_J^2 (-s)^J}{t - M_J^2} \underset{s \rightarrow \infty, t \text{ fixed}}{\sim} s^{J_{\max}}$$

This is not necessarily the case if there are only many exchange diagrams...

$$\sum_{J=0}^{\infty} \frac{-g_J^2 (-s)^J}{t - M_J^2} \sim ?$$

(Dolan-Horn-Schmid) Duality

With infinity of t-channel exchanges, perhaps s-channel poles arise automatically?

$$\sim A(s, t) + t(\lambda_1 \lambda_2 \lambda_3 \lambda_4) + (\text{other flavor orderings})$$

$$A_t(s, t) \sim \sum_J \frac{-g_J^2 (-s)^J}{t - M_J^2}$$

$$A_s(s, t) \sim \sum_J \frac{-g_J^2 (-t)^J}{s - M_J^2}$$

$$\text{QFT: } A(s, t) = A_t(s, t) + A_s(s, t)$$

$$\text{"Dual Model": } A(s, t) = A_t(s, t) \stackrel{?}{=} A_s(s, t)$$

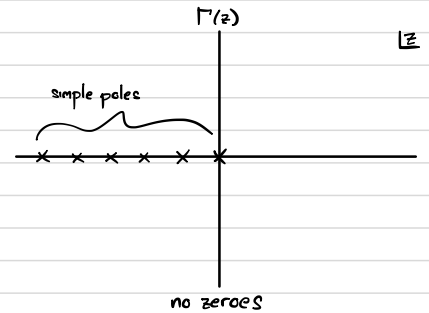
Historical Motivation

Veneziano (1968): $A(s,t) = \frac{\Gamma(-\alpha(s))\Gamma(-\alpha(t))}{\Gamma(-\alpha(s)-\alpha(t))}$

$$\Gamma(z) = \int_0^\infty t^{z-1} e^{-t} dt$$

$$\alpha(s) = \alpha(0) + \alpha' s$$

Regge intercept
Regge slope



$$\begin{cases} s\text{-channel poles: } s = (-\alpha(0) + n) / \alpha', & n=0,1,2,\dots \\ t\text{-channel poles: } t = (-\alpha(0) + n) / \alpha', & n=0,1,2,\dots \\ \text{double poles cancel w/denominator.} \end{cases}$$

Claim: $A(s,t)$ satisfies DHS duality...

$$\begin{aligned} A(s,t) &= - \sum_{n=0}^{\infty} \frac{(\alpha(s)+1)(\alpha(s)+2)\dots(\alpha(s)+n)}{n!} \frac{1}{(\alpha(t)-n)} \\ &= - \sum_{n=0}^{\infty} \frac{(\alpha(t)+1)(\alpha(t)+2)\dots(\alpha(t)+n)}{n!} \frac{1}{(\alpha(s)-n)} \end{aligned}$$

From this representation, see that poles/resonances arrange on linear trajectories w/max spin $J=n \Leftrightarrow M^2 = \frac{-\alpha(0)+n}{\alpha'}$.

Historical Motivation

Virasoro (1969), Shapiro (1970) :

$$A(s, t, u) = \frac{\Gamma(-\alpha_c(s))\Gamma(-\alpha_c(t))\Gamma(-\alpha_c(u))}{\Gamma(-\alpha_c(s)-\alpha_c(t))\Gamma(-\alpha_c(t)-\alpha_c(u))\Gamma(-\alpha_c(u)-\alpha_c(s))}$$

$$\begin{bmatrix} \alpha_c(x) = \alpha(0) + \frac{\alpha' x}{4} \\ \alpha'(s+t+u) = -16\alpha(0) \end{bmatrix}$$

$$\begin{cases} s\text{-channel poles: } s = 4(-\alpha(0) + n), & n = 0, 1, 2, \dots \\ t\text{-channel poles: } & \text{''} \\ u\text{-channel poles: } & \text{''} \end{cases}$$

duality b/w all 3-channels.

max spin @ $m^2 = 4(-\alpha(0) + n)$ is $J = 2n$

Before long, people generalized to external particles other than lightest scalar.

⇒ Veneziano model

⇒ Virasoro-Shapiro model

High Energy Behaviour

Does this solve UV problems?

Veneziano in Regge limit ($s \gg 1$, $t < 0$ fixed)

$$A(s, t) \sim \Gamma(-\alpha(t))(-\alpha(s))^{\alpha(t)} \sim s^{\alpha(t)} \uparrow$$

$-\alpha(t)$

$\alpha(t) = \alpha(0) + \alpha' t$, so suppressed like $s^{-\alpha' |t|}$

(looks like Regge behaviour from exchange of single particle of negative spin $J(t) = -\alpha(0) - \alpha' t$, the "Regge-on")

Veneziano in fixed-angle, hard scattering limit ($s \gg 1$, t/s fixed)

$$A(s, t) \sim [F(\theta_s)]^{-\alpha(s)}$$

$\uparrow$ real function greater than 1

Falls off exponentially fast w/s ∇

These are just formulas, but this suggests the theory behind them has good UV behaviour.

This will turn out to be true @ loop level as well...

A Theory of Strings

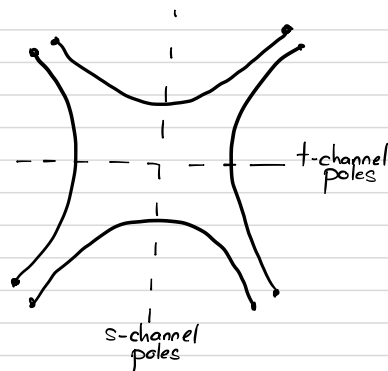
Theory behind Veneziano & Virasoro - Shapiro model turns out to be a theory where elementary particles are replaced by vibrating strings.

This gives heuristic justification for various good properties:

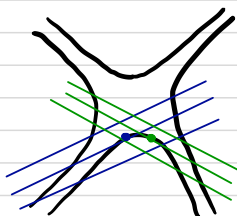
Spectrum:



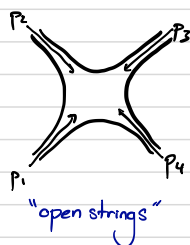
Duality:



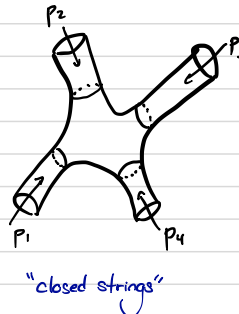
High-energy behaviour:



$A_{\text{Veneziano}}(s, t):$



$A_{\text{Virasoro/Shapiro}}(s, t, u):$



Gravity & the String Scale

Veneziano / VS - amplitudes depend on two parameters: $\alpha(0)$ & α'
 $\alpha(0)$ is a pure number, α' is dimensional, goes like $[mass]^{-2}$

Original idea: $\alpha' \sim (GeV)^{-2}$ (nuclear physics energy scale)

$$-\frac{\alpha(0)}{\alpha'} = \text{mass}^2 \text{ of lightest scalar} = m_\pi^2$$

Problems:

- QCD is correct theory of strong nuclear force (solves UV problems differently).
- Unitarity of amplitude not manifest (require correct signs of residues @ poles when decomposed in Legendre polynomials)
"perturbative partial wave unitarity"

Eventually it was shown that one should have $\alpha(0) = 1$ ($m_\pi^2 < 0$???)

→ Veneziano has pole for massless spin 1 (gauge boson?)

→ VS has pole for massless spin 2 (graviton ???)

Suggests $\alpha' \sim (10^{19} GeV)^{-2}$

J. Scherk & J. Schwarz (1974): "Dual Models For Nonhadrons" (very readable!)

Show: "zero slope limit" (i.e. $\alpha' E^2 \ll 1$) of VS model gives tree-level gravity + scalar theory

$$S_{\text{zero-slope}} = - \int d^4x \sqrt{g} \left\{ \frac{1}{16\pi G_N} R + \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi \right\} \quad \omega / \sqrt{8\pi G_N} = g_{\text{str}} \sqrt{\alpha'}$$