

STRING THEORY

II

LECTURE IV

OXFORD UNIVERSITY
MATH PHIS TT 2020

SPECTRUM OF MASSLESS PHYSICAL STATES: 10 d: $SO(8)$

BOSONIC

(NS_+, NS_+)

$$\begin{aligned} \mathcal{H}_v \otimes \mathcal{H}_v &= 1 \oplus 28 \oplus 35_v \\ \xi_i b_{-\frac{1}{2}}^i |0\rangle_{NS} \otimes \tilde{\xi}_i \tilde{b}_{-\frac{1}{2}}^i |\bar{0}\rangle_{NS} & \quad \downarrow \\ & \text{scalar} \quad [2] \quad \text{ANTISYMMETRIC 2-TENSOR } B \\ & \quad \quad \quad (2) \quad \text{SYMMETRIC TRACELESS 2-TENSOR: } G. \end{aligned}$$

(R_+, R_+)

$$\begin{aligned} \mathcal{H}_s \otimes \mathcal{H}_s &= 1 \oplus 28 \oplus 35_s \\ & \quad \downarrow \\ & [0] \text{ scalar} \quad [2] \quad [4]_+ \\ & \quad \quad \quad C_0 \quad C_2 \quad C_4 \quad \uparrow \text{SELF-DUAL.} \end{aligned}$$

(R_+, R_-)

$$\begin{aligned} \mathcal{H}_s \otimes \mathcal{H}_c &= \mathcal{H}_v \oplus 56_v \\ & \quad \downarrow \\ & [1] \quad [3] \\ & \quad \quad C_1 \quad C_3 \end{aligned}$$

RR-p-FORMS

IB IA.

FERMIONIC

(NS_+, R_+)
 (R_+, NS_+)

$$\begin{aligned} \mathcal{H}_v \otimes \mathcal{H}_s &= \mathcal{H}_c \oplus 56_c \\ & \quad \downarrow \\ & \text{spin } 1/2 \quad \text{spin } 3/2 \\ & \lambda \hat{=} \text{DILATINO. } \psi \hat{=} \text{GRAVITINO.} \end{aligned}$$

(NS_+, R_-)
 (R_-, NS_+)

$$\begin{aligned} \mathcal{H}_v \otimes \mathcal{H}_c &= \mathcal{H}_s \oplus 56_s \\ & \quad \downarrow \\ & \tilde{\lambda} \quad \tilde{\psi} \end{aligned}$$

EFFECTIVE ACTIONS: IIA & IIB STRINGS.

BOSONIC STRING:

[REF: VAN PROYEN & FREEDMAN: SUPERGRAVITY (CUP)]

THE EFFECTIVE ACTION OF Φ, B, G (MASSLESS MODES) WERE DETERMINED BY IMPOSING WEYL INVARIANCE:

ANOMALY " " : $T^a_a \sim \beta^G \partial^a X^\mu \partial_a X^\nu + \dots$

\Rightarrow SETTING $\beta^G = \beta^B = \beta^\Phi = 0 \iff$ SPACE TIME EOM FOR Φ, B, G .

WE CAN DO THE SAME HERE (NS, NS) SECTOR WORKS IDENTICAL TO THE BOSONIC CASE:

$$S_{NSNS} = \frac{1}{2\alpha_{10}^2} \int d^{10}x (-G)^{1/2} e^{-2\Phi} \left(R + 4\partial_\mu \Phi \partial^\mu \Phi - \frac{1}{2} |H_3|^2 \right)$$

\uparrow
10d Gravitational coupling constant.

$$|H_3|^2 = \frac{1}{3!} G^{M_1 N_1} G^{M_2 N_2} G^{M_3 N_3} \times (H_3)_{M_1 M_2 M_3} (H_3)_{N_1 N_2 N_3}$$

FOR (RR) & (R, NS) & (NS, R) SECTOR FIELDS WE COULD PROCEED SIMILARLY, BUT WE WILL MAKE OF THE FACT THAT W/ $N=2$ SUSY (MEANING 2×16 dim spinors as supercharges).

THERE ARE TWO SUPERGRAVITY THEORIES: IIA, IIB.

[PREDATED THE STRING THEORIES].

BEFORE WE DISCUSS THESE, LET'S TALK ABOUT P-FORMS.

(DIFFERENTIAL FORMS: [NAKAHARA]).

RR-SECTOR WE HAVE P-FORMS:

$$C_p = \frac{1}{p!} C_{\underbrace{M_1 \dots M_p}_{\text{TOTALLY ANTISYMM.}}} dx^{M_1} \dots dx^{M_p}.$$

$$(A_p \wedge B_q)_{M_1 \dots M_{p+q}} = \frac{(p+q)!}{p! q!} A_{\underbrace{[M_1 \dots M_p}_{\text{TOTAL ANTI-SYMMETRIZATION.}}} B_{M_{p+1} \dots M_{p+q}}]$$

$(p+q)$ -FORM
 "∧" WEDGE
 PRODUCT.

$$A_p \wedge B_q = (-1)^{pq} B_q \wedge A_p.$$

C_p WILL BE LIKE p -dim GAUGE POTENTIALS.

$$dC_p = F_{p+1} \quad (p+1)\text{-FORM:}$$

$$(dC_p)_{\mu_1 \dots \mu_{p+1}} = (p+1) \underbrace{[C_{\mu_1 \mu_2 \dots \mu_{p+1}}]} \quad \text{EXTERIOR DERIVATIVE}$$

HODGE STAR: $*$:

$$* A_{\mu_1 \dots \mu_{d-p}} = \frac{1}{p!} \epsilon_{\mu_1 \dots \mu_{d-p} \nu_1 \dots \nu_p} A_{\nu_1 \dots \nu_p}.$$

$(d-p)$
- FORM

$\xrightarrow{*}$

p -FORM.

NOTE: THIS DEPENDS ON THE METRIC!
THE ACTION FOR A p -FORM FIELD IS:

$$\int d^d x (-G)^{1/2} |F_p|^2 = \int d^d x \frac{(-G)^{1/2}}{p!} G^{\mu_1 \nu_1} \dots G^{\mu_p \nu_p}.$$

$$\cdot F_{\mu_1 \dots \mu_p} F^{\mu_1 \dots \mu_p}.$$

$$F_p = dA_{p-1}$$

EOMS:

$$\text{GAUGE TRANSFORMATION: } \delta A_{p-1} = d\lambda_{p-2}.$$

$$dF_p = 0$$

$$d * F_p = 0 \quad \text{BIANCHI IDENTITY.}$$

IIA & IIB ACTIONS ARE FIXED BY SUSY. BOSONIC PART IS

$$S_{IIA}^{RR} = -\frac{1}{4\alpha_{10}^2} \int d^{10}x (-G)^{1/2} (|F_2|^2 + |\tilde{F}_4|^2)$$

$$\tilde{F}_4 = dC_3 - C_1 \wedge H_3$$

$$H_3 = \text{NSNS 3-FORM} \\ = dB_2$$

C_1 & C_3 RR-FORMS.

$$S_{IIA}^{\text{CHERN-SIMONS}} = -\frac{1}{4\alpha_{10}^2} \int \underbrace{B_2 \wedge F_4 \wedge F_4}_{10 \text{ FORM.}}$$

$$S_{IIB}^{RR} = -\frac{1}{4\alpha_{10}^2} \int d^{10}x (-G)^{1/2} (|F_1|^2 + |\tilde{F}_3|^2 + \frac{1}{2} |\tilde{F}_5|^2)$$

$$\tilde{F}_3 = F_3 - C_0 \wedge H_3$$

$$H_3 = \text{NSNS 3-FORM} \\ = dB_2$$

$$\tilde{F}_5 = F_5 - \frac{1}{2} \underline{C_2} \wedge \underline{H_3} + \frac{1}{2} \underline{B_2} \wedge \underline{F_3}$$

RECALL: C_4 WAS IN $[4]_+$ REP. dC_4

$$\tilde{F}_5 = * \tilde{F}_3$$

SELF-DUALITY^(SD) OF THE 5-FORM FIELD STRENGTH NEEDS TO BE IMPOSED IN ADDITION (~~A~~ ACTION FOR A SD F_5).

$$S_{\text{IIB}}^{\text{CHERN-SIMONS}} = - \frac{1}{4\alpha_{10}^2} \int \underbrace{C_4 \wedge H_3 \wedge F_3}_{10 \text{ FORM.}}$$

NSNS SECTOR: AS IN THE BOSONIC CASE:

\mathbb{G} (EH ACTION)

Φ (SCALAR)

B_2 \rightsquigarrow COUPLES TO THE STRING.

RR SECTOR: C_p & F_{p+1} FORM FIELDS.

\rightsquigarrow WHAT'S THEIR PURPOSE?

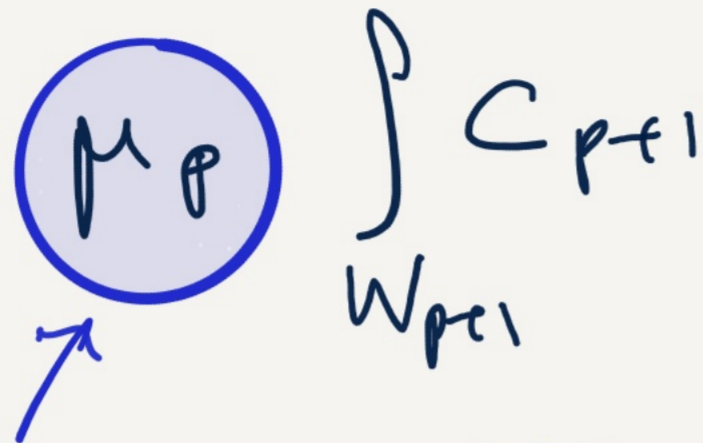
\rightsquigarrow D-BRANES

C_{p+1} : D_p -BRANE
($p+1$ DIM OBJECT)

WE CAN WRITE DOWN
A "WORLD-VOLUME" ACTION

D-BRANES:

D_p-BRANE COUPLES TO THE C_{p+1}-FORM:



$$W_{p+1} = (p+1)\text{-DIM WORLD-VOLUME}$$

CHARGE OF THE D_p-BRANE UNDER THE RR-FORM FIELD.

"D-INSTANTON"
C₀

II B:

$$C_0, C_2, C_4^+$$

→

$$D(2k+1)$$

= P

$$D(-1), D1, D3, D5, D7$$

COUPLE TO THE C_p-FORMS EITHER ELECTRICALLY OR MAGNETICALLY.

NOTE: II B THERE A D1-BRANE (D-STRING) & IT COUPLES TO C₂. VERY SIMILAR TO THE STRING ("FUNDAMENTAL STRING") F1. DUALITY THAT TREATS THESE TWO ON THE SAME FOOTING (S-DUALITY OF II B-STRING).

II A:

$$C_1, C_3$$

$$\rightarrow D(2k)$$

$$D0, D2, D4, D6, D8$$

THE EFFECTIVE ACTION ON A D-BRANE RESEMBLES THAT OF THE NAMBU-GOTO ACTION OF THE STRING:

WORLD-VOLUME W_{p+1} :

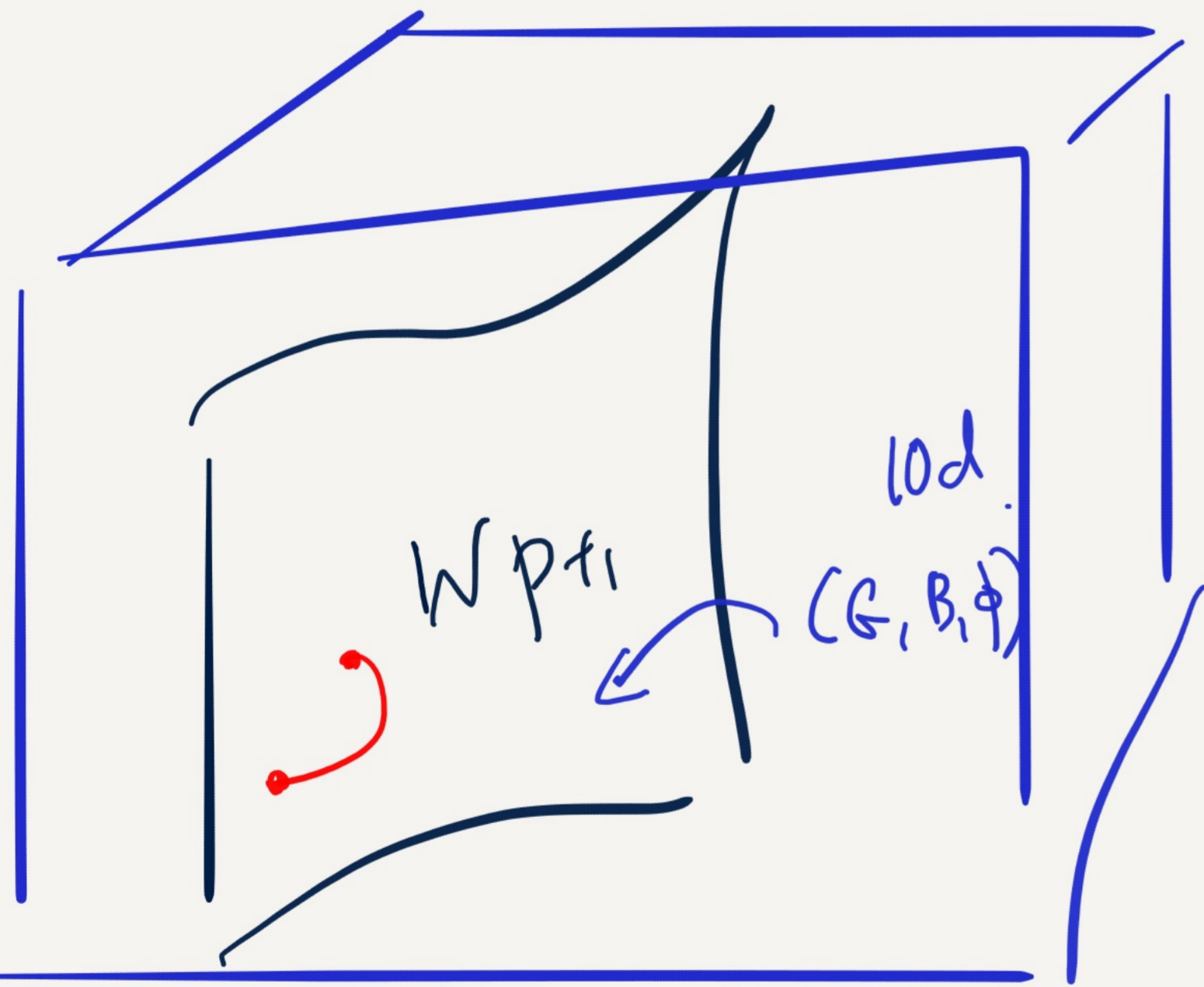
DIRAC-BORN-INFELD ACTION (DBI)

$$S_{D_p} = -T_p \int_{W_{p+1}} d^{p+1} \gamma e^{-\Phi} \sqrt{-\det(G+B+2\alpha' F)}$$

$\int_{W_{p+1}} d^{p+1} \gamma$
 ↑
 COORDINATES ON W_{p+1}

$e^{-\Phi}$
 ↑
 PULLBACK OF Φ (10d METRIC) ONTO W_{p+1}

$\sqrt{-\det(G+B+2\alpha' F)}$
 ↑
 PULLBACK OF G (10d METRIC) ONTO W_{p+1}



$T_p =$ TENSION OF THE D_p -BRANE.

$F = U(1)$ - GAUGE FIELD.
 FROM THE QUANTIZATION OF OPEN STRINGS:

→ $(p+1)$ -DIM SUSY $U(1)$ GAUGE TH.
 T_p & M_p (TENSION & CHARGE)

D_p : 2 PARAMETERS:

IN THE NEXT LECTURE: WE STUDY D_p -BRAES
FROM THE STRING POINT OF VIEW (I.E. 2D CFT).

~o BOUNDARY CFT ASSOCIATED TO THE RNS-STRING.
(IA, IB)

~o COMPUTE T, M, P .
