

# B7.2 ELECTROMAGNETISM

Chapter 5 : Electromagnetic waves (part 2)

Lecture 16

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# 5) Electromagnetic waves

5.1 Plane electromagnetic waves ✓

5.2 Polarization ✓

This lecture:

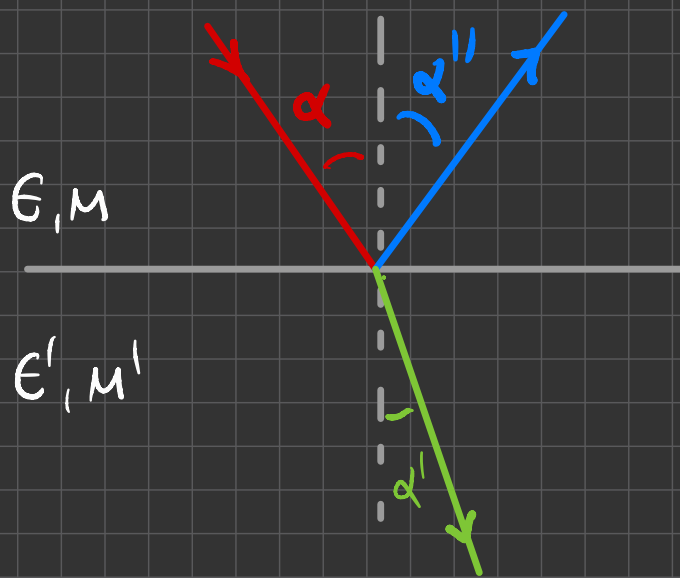
5.3 Reflection & refraction of electromagnetic waves

5.4 Comments on the electromagnetic spectrum

# 15.3] Reflection and refraction of electromagnetic waves

illustrate the wave-like behaviour of the solutions of the wave equation

Consider a **light ray** which crosses a boundary between two media. This ray can



► **reflect** on the surface boundary  
We will see that the angle of reflection equals the angle of incidence

$$\alpha = \alpha'' \quad (i)$$

► **refract** through the surface boundary  
We will see that the ray bends as it enters the new medium

$$n \sin \alpha = n' \sin \alpha' \quad (ii)$$

**Snell's law**

where  $n = \frac{c}{v}$  refraction index of the material

[ air  $n \sim 1$  as  $v \sim c$ ; water  $n \sim 1.3$  ]

Laws (i) & (ii) can be deduced from Maxwell's equations

$$\left( v = \frac{1}{\sqrt{\epsilon \mu}} \right)$$

Let the plane  $z=0$  be the surface boundary between the two media.

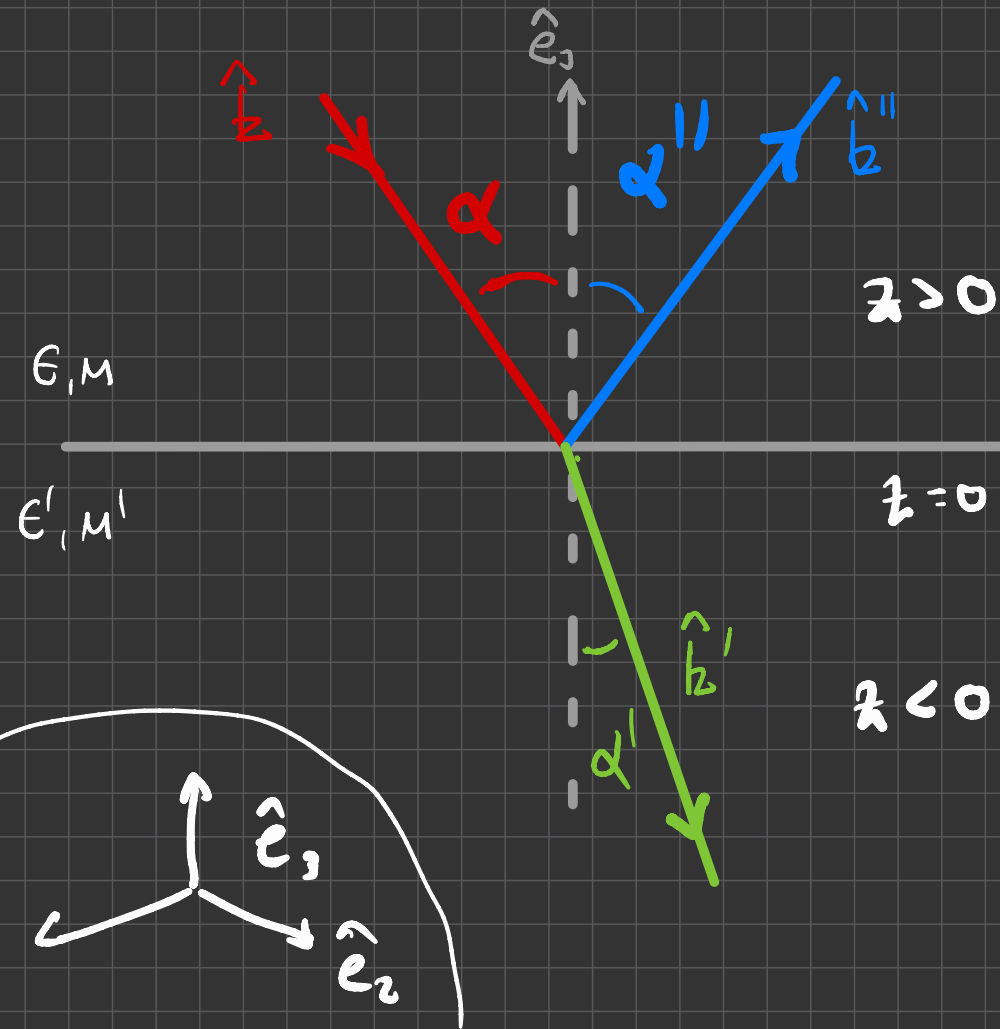
Let  $\hat{k}$  be the direction of propagation of the incident wave be it makes an angle  $\alpha$  with  $\hat{e}_3$

For the reflected wave we have:

$\hat{k}''$  = direction of propagation making an angle  $\alpha''$  with  $\hat{e}_3$

For the refracted wave we have

$\hat{k}'$  = direction of propagation making an angle  $\alpha'$  with  $\hat{e}_3$



Experimentally: all waves have the same frequency  $\omega$  of the incident wave.

So, in a medium with  $\epsilon, \mu$

$$k = \omega v, \quad v = \frac{1}{\sqrt{\epsilon \mu}} = \frac{1}{n} c$$

velocity of the electromagnetic waves in medium with  $\epsilon, \mu$  ( $n > 0$ )

Hence  $k'' = k = \frac{\omega}{v}$

and  $k' = \frac{\omega}{v'}, \quad v' = \frac{1}{\sqrt{\epsilon' \mu'}} = \frac{1}{n'} c$

velocity of the electromagnetic waves in medium with  $\epsilon', \mu'$  ( $n' < 0$ )

# Electromagnetic fields

$\lambda > 0$

by the superposition principle

$$\vec{E} = \underbrace{\vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}}_{\vec{E}_{\text{incident}} = \vec{E}_I} + \underbrace{\vec{E}_0'' e^{i(\vec{k}'' \cdot \vec{r} - \omega t)}}_{\vec{E}_{\text{reflected}} = \vec{E}''}$$

$$\vec{B} = \frac{1}{\omega} (\vec{k} \wedge \vec{E}_I + \vec{k}'' \wedge \vec{E}'')$$

$\lambda < 0$

$$\vec{E}' = \vec{E}_0' e^{i(\vec{k}' \cdot \vec{r} - \omega t)}$$

refracted wave

$$\vec{B}' = \frac{1}{\omega} \vec{k}' \wedge \vec{E}'$$

Want to find  $\vec{k}', \vec{k}'', \vec{E}_0', \vec{E}_0''$  in terms of the given (arbitrary) quantities  $\vec{k}$  &  $\vec{E}_0$ .

How? require the correct boundary conditions at  $z=0$

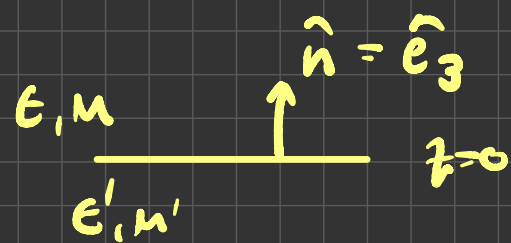
These are:

(A) normal components of  $\epsilon \vec{E}$  and  $\vec{B}$  are continuous across the boundary

(B) tangential components of  $\vec{E}$  and  $\frac{1}{\mu} \vec{B}$  are continuous across the boundary



Then at  $\boxed{z=0}$



$$\textcircled{A} \Rightarrow \begin{cases} (\epsilon (\vec{E}_I + \vec{E}''') - \epsilon' \vec{E}') \cdot \hat{n} = 0 & (1) \\ (\vec{B}_I + \vec{B}''' - \vec{B}') \cdot \hat{n} = 0 & (2) \end{cases}$$

$$\textcircled{B} \Rightarrow \begin{cases} (\vec{E}_I + \vec{E}''' - \vec{E}') \wedge \hat{n} = 0 & (3) \\ \left( \frac{1}{\mu} (\vec{B}_I + \vec{B}''') - \frac{1}{\mu'} \vec{B}' \right) \wedge \hat{n} = 0 & (4) \end{cases}$$

(1) - (4) must be true at  $z=0$   $\boxed{\text{for all}}$   $(x, y)$

At  $t=0$

$$(1) \quad [E(\vec{E}_0 e^{i\vec{k}\cdot\vec{r}} + \vec{E}_0'' e^{i\vec{k}''\cdot\vec{r}}) - e' \vec{E}_0' e^{i\vec{k}'\cdot\vec{r}}] \cdot \hat{n} = 0$$

$$(2) \quad [\vec{k} \wedge \vec{E}_0 e^{i\vec{k}\cdot\vec{r}} + \vec{k}'' \wedge \vec{E}_0'' e^{i\vec{k}''\cdot\vec{r}} - \vec{k}' \wedge \vec{E}_0' e^{i\vec{k}'\cdot\vec{r}}] \cdot \hat{n} = 0$$

$$(3) \quad [\vec{E}_0 e^{i\vec{k}\cdot\vec{r}} + \vec{E}_0'' e^{i\vec{k}''\cdot\vec{r}} - \vec{E}_0' e^{i\vec{k}'\cdot\vec{r}}] \wedge \hat{n} = 0$$

$$(4) \quad \left[ \frac{1}{\mu} (\vec{k} \wedge \vec{E}_0 e^{i\vec{k}\cdot\vec{r}} + \vec{k}'' \wedge \vec{E}_0'' e^{i\vec{k}''\cdot\vec{r}}) - \frac{1}{\mu'} \vec{k}' \wedge \vec{E}_0' e^{i\vec{k}'\cdot\vec{r}} \right] \wedge \hat{n} = 0$$

This is true at  $t=0$   $\nabla(x, y)$  only if

$$\vec{k} \cdot \vec{r} \Big|_{t=0} = \vec{k}' \cdot \vec{r} \Big|_{t=0} = \vec{k}'' \cdot \vec{r} \Big|_{t=0} = 0$$

# WLOG

assume  $\vec{b}$  lies on the  $(x, z)$  plane

$$\text{ic } b_y = 0$$

$$\text{Then } \vec{r}|_{z=0} = x \hat{e}_1 + y \hat{e}_2$$

$$\vec{b} = b_x \hat{e}_1 + b_z \hat{e}_3$$

$$\vec{b} \cdot \vec{r}|_{z=0} = x b_x$$

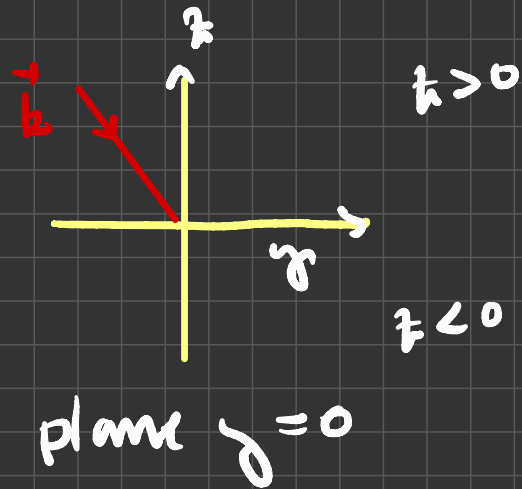
$$\vec{b}'' \cdot \vec{r}|_{z=0} = x b_x'' + y b_y''$$

$$\vec{b}' \cdot \vec{r}|_{z=0} = x b_x' + y b_y'$$



$$x b_x = x b_x'' + y b_y'' = x b_x' + y b_y'$$

$$\nabla^2(x, y)$$



$\Rightarrow b_y' = b_y'' = 0$  i.e.  $\vec{b}, \vec{b}'$  &  $\vec{b}''$  lie on the same plane  
(the  $(x, z)$  plane)

and  $b_x = b_x'' = b_x'$

$\Rightarrow$

$$\left. \begin{aligned} \vec{b} &= b_x \hat{e}_1 + b_z \hat{e}_3 \\ \vec{b}'' &= b_x \hat{e}_1 + b_z'' \hat{e}_3 \\ \vec{b}' &= b_x \hat{e}_1 + b_z' \hat{e}_3 \end{aligned} \right\} \text{with } b_z'' = b_z$$

$$k_x = k_x'' = k_x'$$

$$\Leftrightarrow k \sin \alpha = k'' \sin \alpha'' = k' \sin \alpha'$$

$$k \sin \alpha''$$

$$\Leftrightarrow \sin \alpha = \sin \alpha''$$

$$k \sin \alpha = k' \sin \alpha'$$

$$\left( k = \frac{\omega}{v}, \quad k' = \frac{\omega}{v'} \right)$$

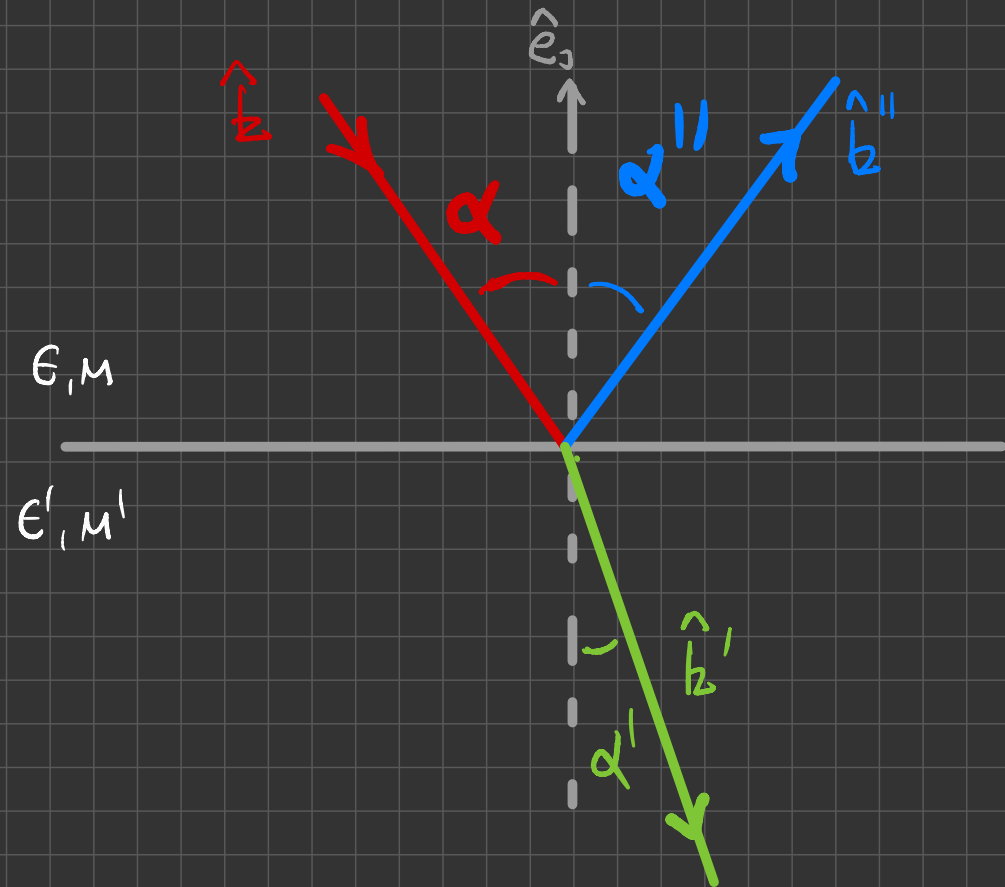
$$\Leftrightarrow$$

$$\alpha = \alpha'$$

$$n \sin \alpha = n' \sin \alpha'$$

$$\left( n = \frac{c}{v}, \quad n' = \frac{c}{v'} \right)$$

Snell's law



Remark: The refractive index depends on the frequency of the light

You have described the dispersion of light by a prism

↑  
separation of light into different colors  
(frequencies!)

waves with higher frequencies (blue) "bend" more  
than those with lower frequencies (red)

You have also observed rainbows:  
(can you explain the double rainbow?)



(Gilded Post Meadow)

Returning to the boundary conditions at  $z=0$

$$(1) \quad [\epsilon (\vec{E}_0 + \vec{E}_0'') - \epsilon' \vec{E}_0'] \cdot \hat{n} = 0$$

$$(2) \quad [\vec{b} \wedge \vec{E}_0 + \vec{b}'' \wedge \vec{E}_0'' - \vec{b}' \wedge \vec{E}_0'] \cdot \hat{n} = 0$$

$$(3) \quad [\vec{E}_0 + \vec{E}_0'' - \vec{E}_0'] \wedge \hat{n} = 0$$

$$(4) \quad \left[ \frac{1}{\mu} (\vec{b} \wedge \vec{E}_0 + \vec{b}'' \wedge \vec{E}_0'') - \frac{1}{\mu'} \vec{b}' \wedge \vec{E}_0' \right] \wedge \hat{n} = 0$$

with  $\alpha = \alpha''$  &  $n \sin \alpha = n' \sin \alpha'$   $\hat{n} = \hat{e}_3$

$$k = k'' = \frac{\omega}{c} n \quad k' = \frac{\omega}{c} n'$$

$\vec{b}$ ,  $\vec{b}'$  &  $\vec{b}''$  all lie on the same plane ( $(x, z)$ -plane)



Example: Suppose the electric field of the incident ray is linearly polarized with

$$\vec{E}_0 = E_0 \hat{e}_2 \quad (\text{perpendicular to } (x, z) \text{ plane})$$

Then  $\vec{E}_0'$  &  $\vec{E}_0''$  are also in the  $y$ -direction

$$(3) : [\vec{E}_0 + \vec{E}_0'' - \vec{E}_0'] \wedge \hat{n} = 0 \quad \hat{n} = \hat{e}_3$$

$$\Rightarrow \underline{E_0 + E_0'' - E_0' = 0}$$

$$(4) : \left[ \frac{1}{\mu} (\vec{b} \wedge \vec{E}_0 + \vec{b}'' \wedge \vec{E}_0'') - \frac{1}{\mu_1} \vec{b}' \wedge \vec{E}_0' \right] \wedge \hat{n} = 0$$

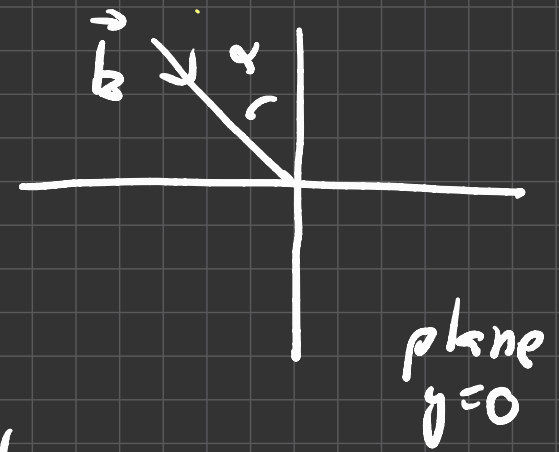
$$(\vec{b} \wedge \hat{e}_2) \wedge \hat{e}_3 = -b_x \hat{e}_2$$

$$\therefore \frac{1}{\mu} (b_x E_0 + b_x'' E_0'') - \frac{1}{\mu_1} b_x' E_0' = 0$$

$$k_z = -k \cos \alpha$$

$$k_z'' = k \cos \alpha$$

$$k_z' = -k' \cos \alpha'$$



$$\frac{1}{\mu} k \cos \alpha (-E_0 + E_0'') + \frac{1}{\mu'} k' \cos \alpha' E_0' = 0$$

$$\frac{k}{\mu} = \frac{\omega}{v_m} = \omega \sqrt{\frac{\epsilon}{\mu}}$$

$$\Rightarrow \sqrt{\frac{\epsilon}{\mu}} \cos \alpha (E_0'' - E_0) = \sqrt{\frac{\epsilon'}{\mu'}} \cos \alpha' E_0'$$

Assuming  $\mu = \mu'$ :  $E_0' = \frac{2 \cos \alpha \sin \alpha'}{\sin(\alpha + \alpha')} E_0$

$$E_0'' = -\frac{\sin(\alpha - \alpha')}{\sin(\alpha + \alpha')} E_0$$

## 15.4 Comments on the electromagnetic spectrum

Electromagnetic spectrum  $\rightarrow$  "full" range of electromagnetic frequencies

Recall: everything radiates electromagnetic energy  
(our own bodies, the stars -)

the frequency of the electromagnetic waves depends on the temperature of the body

When you look at the night sky: what is the light being detected is telling you about the celestial bodies you are observing?

you can ask for example:

why is the sun "yellow" or some other star "red"?

red  
↓  
softer



blue  
↓  
hotter

higher frequencies →

sun

surface 5778 kelvin ←  
core  $10^8$  kelvin

↖ hot, high pressure (due to gravitational pull)

⇒ nuclear fusion

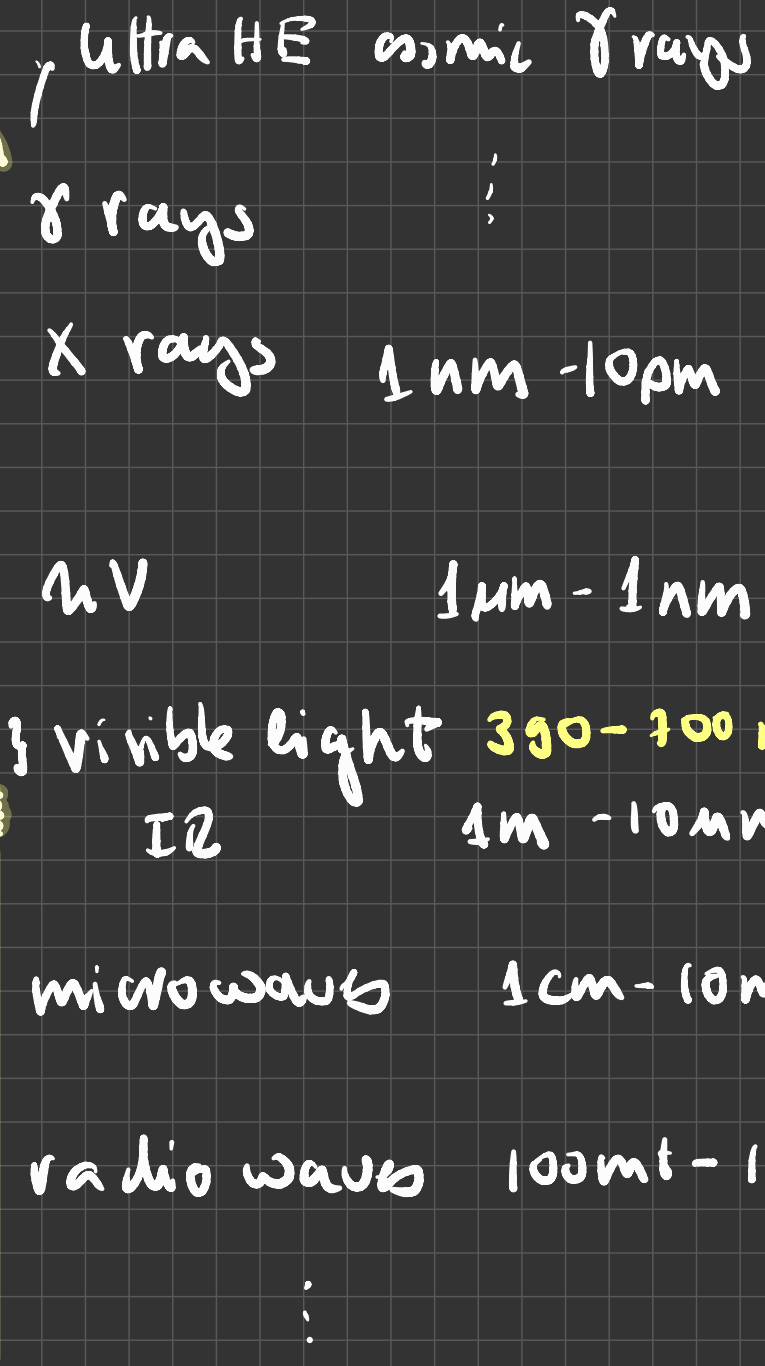
$H \rightarrow He + \text{photons of different frequencies}$

(also heat & energetic charged particles)

↖ solar wind

increasing frequency/energy ↑

γ-ray & X-ray astronomy  
(study radiation of celestial bodies)  
radiation expected from bodies which contain very hot gases ( $10^6 - 10^8$  Kelvin)



the funny names are historical

thermal body radiation →

think of the microwave oven: can transfer energy to heat food →

↓ increasing wavelength

Next (and the last lecture)

electromagnetism & special relativity