

$u^T A u = \lambda(u) \cdot u$   
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 $\lambda^T = \frac{1}{h} [2 \cos \frac{h\pi}{m} - 2]$ ,  $h \rightarrow 0$   
 $[(\lambda+1)h = h \cdot a]$   
 so  $\lambda_{max} = \frac{1}{h} [2 \cos \frac{\pi}{m} - 2]$   
 $= -\frac{4}{h^2} + O(h)$   
 $\lambda_{min} = \frac{1}{h} [2 \cos \frac{\pi}{m} - 2]$   
 $= -\frac{4}{h^2} + O(h)$

so  $|\lambda_{min}(h)| = \frac{4}{h^2} + O(h)$  is constant independent of  $h$   
 $\Rightarrow \|e\|_2 \leq \frac{1}{\lambda_{min}} \left( \sum_{j=1}^n |u_j| \right)^2$   
 $\leq \frac{4h^2}{\pi^2} \left( \sum_{j=1}^n |u_j| \right)^2 \approx \frac{4h^2}{\pi^2} \|u\|_2^2$

Also useful in this context (and also in small case) is Gershgorin's Theorem

**Gershgorin's Theorem**  
 $A = (a_{ij})_{i,j=1,\dots,n} \in \mathbb{R}^{n \times n}$   
 then  $\lambda$  lies in at least one of the Gershgorin disks  
 $\{ z \in \mathbb{C} : |z - a_{ii}| \leq \sum_{j \neq i} |a_{ij}| \}$

**Proof**  $\lambda$  an eigenvalue  $\Rightarrow \exists x \neq 0$  with  $Ax = \lambda x$  or  $\sum_{j=1}^n a_{ij} x_j = \lambda x_i$ ,  $i=1,\dots,n$

Suppose  $|x_i| \geq |x_j|$ ,  $i=1,\dots,n$  then  
 $\sum_{j=1}^n a_{ij} x_j = \lambda x_i$   
 $\Leftrightarrow (a_{ii} - \lambda) x_i = -\sum_{j \neq i} a_{ij} x_j$   
 $\Rightarrow |a_{ii} - \lambda| |x_i| \leq \sum_{j \neq i} |a_{ij}| |x_j|$   
 $\Leftrightarrow |a_{ii} - \lambda| \leq \sum_{j \neq i} |a_{ij}| \frac{|x_j|}{|x_i|}$   
 $\leq \sum_{j \neq i} |a_{ij}|$

**Example**  $u'' - u = f$   
 $A = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$

Gershgorin's Th  $\Rightarrow$  for any eigenvalue  $\lambda$   
 $|\lambda + \frac{1}{2} + \lambda| \leq \frac{1}{2}$   
 $\Rightarrow \text{Re } \lambda \leq 0$  i.e.  $\lambda \leq -1$   
 $\Rightarrow \|e\|_1 \leq \|e\|_2$ ,  $\|e\|_2 \leq \|e\|_1$

**Remark**, Gershgorin's Theorem is useful for determining stability of numerical schemes

**Derivative boundary conditions** e.g.  $u'(a) = \alpha$

$\frac{1}{h} \frac{u_1 - u_0}{2} = \alpha$   
 then e.g. for  $u'' = f$  we use the finite difference centered on  $x_0$ :  
 $u_1 - 2u_0 + u_{-1} = h^2 f_0$   
 use  $\textcircled{1}$  to eliminate  $u_{-1}$ :  
 $\Rightarrow 2u_1 - 2u_0 = h^2 f_0 + 2h\alpha$   
 or  $h$  pressure symmetry  
 $u_1 - u_0 = \frac{1}{2} h^2 f_0 + h\alpha$

$\Rightarrow \begin{bmatrix} 1 & -1 \\ 2 & -2 \end{bmatrix} \begin{bmatrix} u_1 \\ u_0 \end{bmatrix} = \begin{bmatrix} \frac{1}{2} h^2 f_0 + h\alpha \\ h^2 f_0 \end{bmatrix}$

Note if  $u'' = f$  with  $u(a) = \alpha$ ,  $u'(b) = \beta$   
 then

$\begin{bmatrix} 1 & -1 & 0 & \dots & 0 \\ 0 & 1 & -2 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & \dots & 0 & \dots & 0 & 1 \end{bmatrix} \begin{bmatrix} u_n \\ \vdots \\ u_1 \\ u_0 \end{bmatrix} = \begin{bmatrix} 0 \\ \vdots \\ h^2 f_0 \\ \alpha \end{bmatrix}$

note this is singular since  $\begin{bmatrix} 1 & -1 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} u_1 \\ u_0 \end{bmatrix} = \begin{bmatrix} 0 \\ \alpha \end{bmatrix}$

in given  $f_0 = 0$  and  $\alpha = 0$  so right side, then it is a homogeneous system  
 $u_0 = u_1 = \dots = u_n = 0$   
 but this corresponds exactly to the ODE BVP:  $u'' = 0$   
 $u(a) = 0 = u(b)$   
 since  $u = 0$  constant satisfies the homogeneous BVP and hence any constant can be added to any particular solution  $u$  of  $u'' = f$ ,  $u(a) = \alpha$ ,  $u'(b) = \beta$ .

'fixed' boundary conditions  
 $C u(a) + u'(a) = \alpha$   
 can also be approximated by  
 $C u_0 + \frac{u_1 - u_0}{h} = \alpha$   
 and  $u_1$  eliminated from  
 $u_1 - 2u_0 + u_{-1} = h^2 f_0$   
 (or  $\tilde{f}_0 = \frac{1}{2} h^2 f_0 + h\alpha$ )  
 $\Rightarrow \begin{bmatrix} 1 & -1 \\ 2 & -2 \end{bmatrix} \begin{bmatrix} u_1 \\ u_0 \end{bmatrix} = \begin{bmatrix} \tilde{f}_0 \\ h^2 f_0 \end{bmatrix}$   
 for  $(\tilde{f}_0) = f$  etc.

**Finite differences for higher order problems**

e.g. beam problem  $u'''' = f$   
 as an alternative to introducing  $w = u''$  can approximate  $u''''$  directly:  
 $u(x_j+2h) - 4u(x_j+h) + 6u(x_j) - 4u(x_j-h) + u(x_j-2h) = \frac{h^4}{12} u''''(x_j) + O(h^6)$   
 $\Rightarrow u(x_{j+2}) - 4u(x_{j+1}) + 6u(x_j) - 4u(x_{j-1}) + u(x_{j-2}) = \frac{h^4}{12} f_j + O(h^6)$   
 $\Rightarrow \frac{1}{h^4} (u_{j+2} - 4u_{j+1} + 6u_j - 4u_{j-1} + u_{j-2}) = \frac{1}{12} f_j + O(h^2)$

**Boundary conditions:**  
 $u(a) = \alpha$ ,  $u'(a) = \beta$ ,  $u(b) = \gamma$ ,  $u'(b) = \delta$   
 $u_0 = \alpha$ ,  $\frac{u_1 - u_0}{h} = \beta$   
 $u_n = \gamma$ ,  $\frac{u_n - u_{n-1}}{h} = \delta$   
 $\Rightarrow u_1 = u_0 + h\beta$   
 $u_n = u_{n-1} + h\delta$   
 $\Rightarrow u_{j+1} - u_j = h\beta$  for  $j=0$   
 $\Rightarrow u_{j+1} - u_j = h\delta$  for  $j=n-1$