C7.5 Lecture 6: Differential geometry 2 Curves, vectors, and tensors

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Curves and tangent vectors

As before, we define a curve

$$\gamma : [0,1] \text{ (or } \mathbb{R} \text{ (or } \mathbb{R}^+) \to \mathcal{M}.$$

What about a tangent vector? Our manifold does not have an affine space structure, nor are there special sets of inertial coordinates we can use.

We *can* still differentiate functions along a curve: given $f : \mathcal{M} \to \mathbb{R}$, we have

$$\frac{\mathrm{d}}{\mathrm{d}\lambda}f\circ\gamma:=V(f),$$

where we use this equation to *define* the "vector" V. It is an operator which acts on scalar fields f via the above formula.

Vectors satisfy the following two important properties: for constants $a, b \in \mathbb{R}$ and functions $f, g : \mathcal{M} \to \mathbb{R}$

1 Linearity:

$$V(af + bg) = aV(f) + bV(g).$$

2 The Leibniz rule:

$$V(fg) = gV(f) + fV(g).$$

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Tangent vectors in local coordinates

In terms of local coordinates x^a , we can set

$$V(f)\big|_{p} = V(f \circ \phi_{U}^{-1} \circ \phi_{U})\big|_{p}$$
$$= V\left(\tilde{f}(x^{a})\right)\big|_{p},$$

where $\tilde{f} = f \circ \phi_U^{-1}$. Note that $\tilde{f} : \mathbb{R}^n \supset U \to \mathbb{R}$, and $x^a = \phi_U(p)$. Using the chain rule we have

$$V(f)\big|_{p} = V(x^{a})\big|_{p} \frac{\partial \tilde{f}}{\partial x^{a}}\Big|_{x^{a}(p)}$$
$$= V^{a}\partial_{a}\tilde{f}.$$

Since this formula holds in *all* local coordinates, we write $V = V^{\mu}\partial_{\mu}$. By a common abuse of notation, people often write f for $\tilde{f} = f \circ \phi_U^{-1}$, although these are two different objects: f is a function on the manifold, while \tilde{f} is a function of the local coordinates x^a (of course, they take the same value at corresponding points!).

Vectors and the tangent space

A vector at a point $p \in \mathcal{M}$ is the tangent vector to some curve¹ through p, at the point p.

The tangent space at p, $T_p(\mathcal{M})$ is the set of all vectors at p.

¹Strictly speaking we need to talk about equivalence classes, because there are multiple curves with the same tangent vector. Two curves γ and γ' , with tangent vectors V and V' at p are said to define the same vector if V(f) = V'(f) for all f.

 $T_p(\mathcal{M})$ is a vector space with the same dimension as the dimension of the manifold.

Given some local coordinates x^a , we can define the vectors $\partial_a = \frac{\partial}{\partial x^a}$ as the vectors tangent to the curves along which x^a changes while x^b , $b \neq a$ remain constant, parametrised by x^a (see figure 1). Such vector fields are sometimes called *coordinate induced vector fields*.



The coordinate induced vector field $\frac{\partial}{\partial x}$ points in the direction where x changes while all the other coordinates (here, the coordinate y) remain the same. Similarly, $\frac{\partial}{\partial y}$ points in the direction where y changes while x remains constant.

Covectors and tensors

The cotangent space $T_p^*(\mathcal{M})$ is the dual space of the vector space $T_p(\mathcal{M})$, i.e. it consists of all linear maps (called *covectors*) from the tangent space to \mathbb{R} .

A tensor of rank (n, m) is an element of $(T_p(\mathcal{M}))^n \times (T_p^*(\mathcal{M}))^m$. Equivalently, it is a multi-linear map from $(T_p(\mathcal{M}))^m \times (T_p^*(\mathcal{M}))^n$ to \mathbb{R} .

Tensor components and their transformation laws

Given local coordinates x^a , the components of the vector X are

$$X^a := X(x^a).$$

By the chain rule, for any scalar function $f:\mathcal{M}\to\mathbb{R}$ we have

$$X(f) = X(f \circ \phi_U^{-1} \circ \phi_U) = X(x^a) \frac{\partial}{\partial x^a} (f \circ \phi_U^{-1}) = X^a \frac{\partial}{\partial x^a} (f \circ \phi_U^{-1}).$$

Note that $\phi_U(p) = (x^0(p), \dots x^{n-1}(p))$, and $f \circ \phi_U^{-1} : \mathbb{R}^n \to \mathbb{R}$.
In particular, the components of the vector ∂_b are

$$(\partial_b)^a = \partial_b(x^a) = \delta_b^a,$$

so

$$(\partial_b)(f) = \frac{\partial}{\partial x^b}(f \circ \phi_U^{-1}).$$

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Now suppose we change coordinates in a neighbourhood of the point p, from the coordinates x^a to coordinates $y^{a'}(x^a)$. Then the new components of the vector X are, using the chain rule,

$$(X')^{a'} = X(y^{a'}) = \frac{\partial y^{a'}}{\partial x^a} X(x^a) = \frac{\partial y^{a'}}{\partial x^a} X^a.$$

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This is the transformation law for vectors.

Let η be a covector. Then the components of η are defined to be

 $\eta_{a} := \eta(\partial_{a}).$

Note that

$$\eta(X) = \eta(X^a \partial_a) = X^a \eta_a$$

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Since this holds in *any* coordinate system, we write $\eta(X) = X^{\mu}\eta_{\mu} = \eta_{\mu}X^{\mu}$.

Under a change of coordinates as before, we have

$$\eta(X) = \eta_a X^a = (\eta')_{a'} (X')^{a'} = (\eta')_{a'} \frac{\partial y^{a'}}{\partial x^a} X^a,$$

so we must have

$$\begin{aligned} \frac{\partial y^{a'}}{\partial x^{a}} (\eta')_{a'} &= \eta_{a} \\ \Rightarrow (\eta')_{a'} &= \frac{\partial x^{a}}{\partial y^{a'}} \eta_{a} \end{aligned}$$

using the inverse function theorem. This is the *covector transformation law*.

Now, under a change of coordinates as above, we have

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More general tensors transform in the "obvious" way – as if they were the product of a bunch of vector and covector fields :

$$(T')^{a'_{1}a'_{2}\dots a'_{n}}_{\begin{array}{c}b'_{1}b'_{2}\dots b'_{m}\end{array}}$$

= $\frac{\partial y^{a'_{1}}}{\partial x^{a_{1}}}\frac{\partial y^{a'_{2}}}{\partial x^{a_{2}}}\dots \frac{\partial y^{a'_{n}}}{\partial x^{a_{n}}}\frac{\partial x^{b_{1}}}{\partial y^{b'_{1}}}\frac{\partial x^{b_{2}}}{\partial y^{b'_{2}}}\dots \frac{\partial x^{b_{m}}}{\partial y^{b'_{m}}}T^{a_{1}a_{2}\dots a_{n}}_{\begin{array}{c}b_{1}b_{2}\dots b_{m}\end{array}$

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