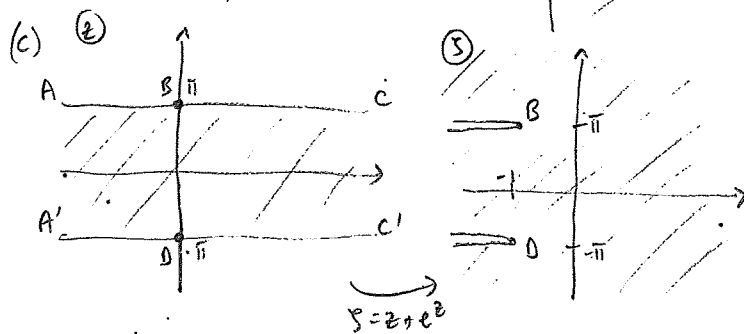
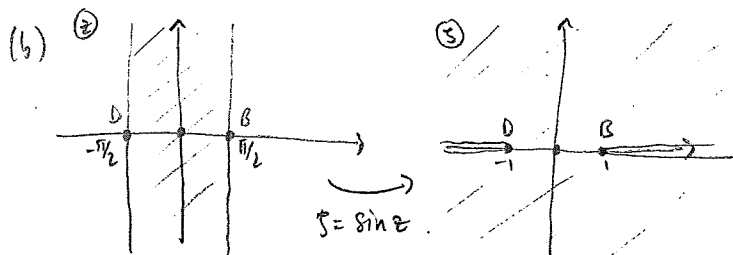
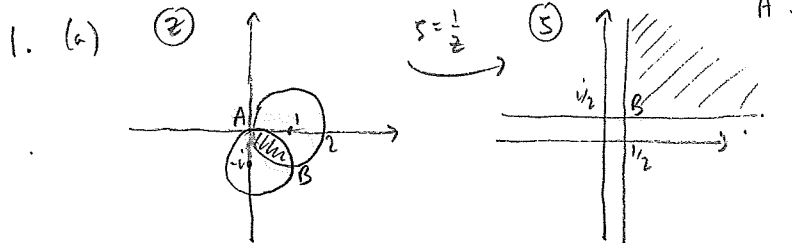




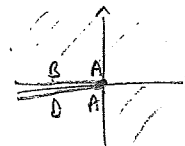
$$S = \sin z = \sin x \cosh y + i \cos x \sinh y$$

Note $f(z) = \sin z$ has $f'(z) = \cos z = 0$ at $z = \pm \pi/2$

where $f'(z) \neq 0$, so angles at B and D are doubled

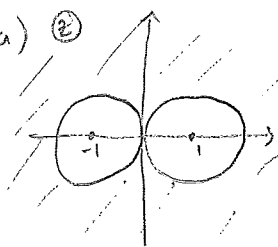
Note e^z maps to:

then have to add z , which shifts the 'cut' up and down by π .

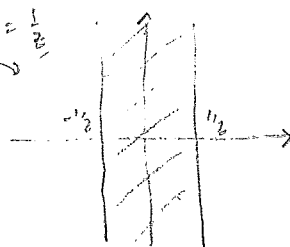


$f'(z) = 1 + e^z = 0$ at $z = \pm i\pi$, where $f'(z) \neq 0$ so angles at B and D are doubled

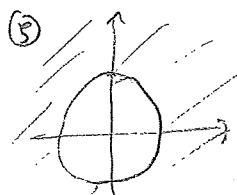
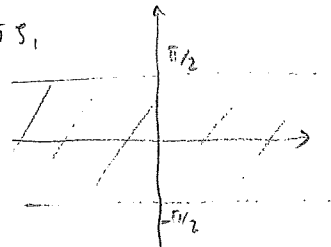
2. (a) ②



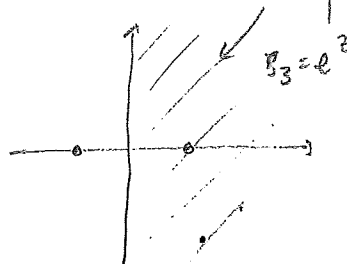
$$S_1 = \frac{1}{z}$$



$$S_2 = i\pi S_1$$



$$S = \frac{S_2 + 1}{S_2 - 1}$$



$$S_3 = e^z$$

so
$$\zeta = \frac{e^{i\pi/2} + 1}{e^{i\pi/2} - 1}$$

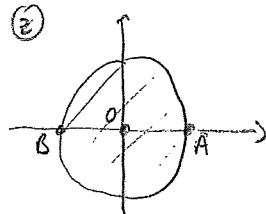
$$= \frac{e^{i\pi/2} + e^{-i\pi/2}}{e^{i\pi/2} - e^{-i\pi/2}}$$

$$= -i \cot \frac{\pi}{2z}$$

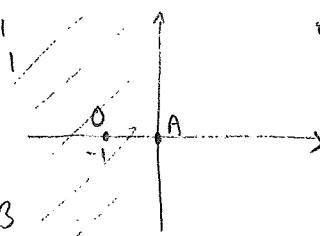
(the map is not unique of course)

(this form is fine)

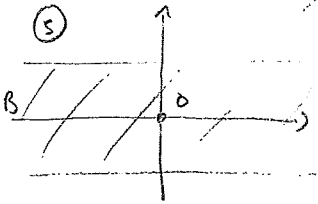
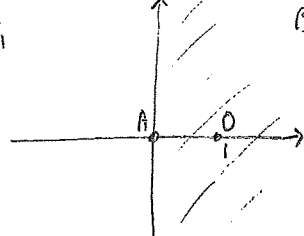
(b) ②



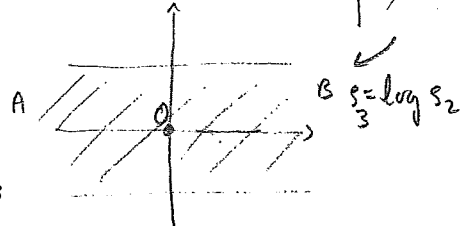
$$S_1 = \frac{z-1}{z+1}$$



$$S_2 = -S_1$$



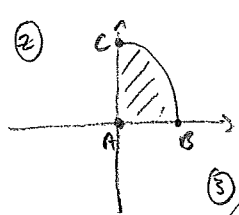
$$S = -S_3$$



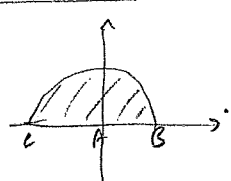
so
$$\zeta = -\log \left(\frac{1-z}{1+z} \right) = \log \left(\frac{1+z}{1-z} \right)$$

(some steps could be easily combined)

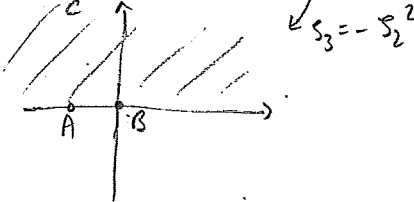
(c) ②



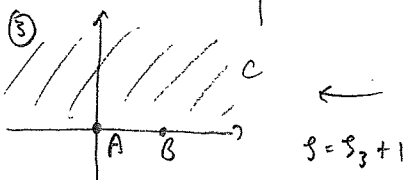
$$S_1 = z^2$$



$$S_2 = \frac{S_1 - 1}{S_1 + 1}$$



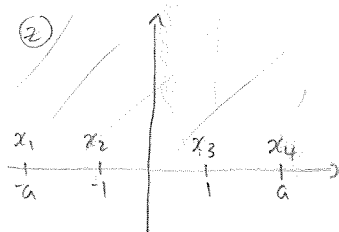
$$S_3 = -S_2^2$$



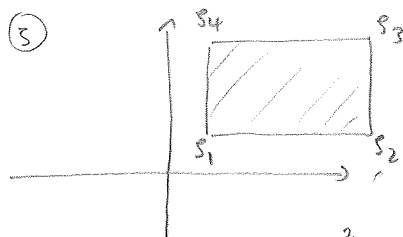
$$S = S_3 + 1$$

so
$$\zeta = 1 - \left(\frac{z^2 - 1}{z^2 + 1} \right)^2 = \frac{4z^2}{(z^2 + 1)^2}$$

3.



$$S = f(z)$$


 $\beta = \frac{1}{2}$ at each vertex

Schwarz-Christoffel formula $\Rightarrow f(z) = A + C \int^z (t-1)^{-1/2} (t+1)^{-1/2} (t-a)^{-1/2} (t+a)^{-1/2} dt$

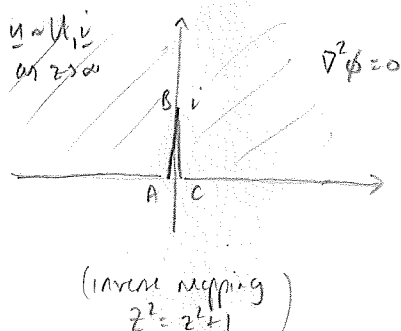
$$= A + C \int^z \frac{dt}{(t^2-1)^{1/2} (t^2-a^2)^{1/2}}$$

A and C only control the location, orientation and scaling of the whole rectangle, so the aspect ratio must be controlled by a (i.e. there is nothing else that could control it).

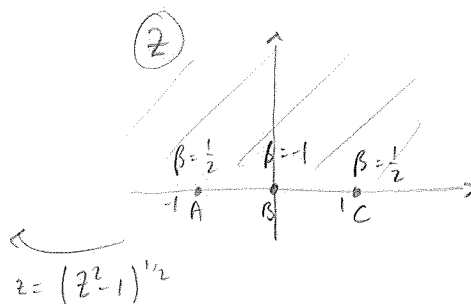
Note the length of the sides are $l_1 = |f(1) - f(-1)|$ and $l_2 = |f(a) - f(1)|$, which are

$$l_1 = |C| \left| \int_{-1}^1 \frac{dt}{(t^2-1)^{1/2} (t^2-a^2)^{1/2}} \right| \quad l_2 = |C| \left| \int_1^a \frac{dt}{(t^2-1)^{1/2} (t^2-a^2)^{1/2}} \right| \quad \text{so} \quad \frac{l_1}{l_2} = \left| \frac{\int_{-1}^1 \frac{dt}{(t^2-1)^{1/2} (t^2-a^2)^{1/2}}}{\int_1^a \frac{dt}{(t^2-1)^{1/2} (t^2-a^2)^{1/2}}} \right|$$

4.



$$\nabla^2 \phi = 0$$



Schwarz-Christoffel.

$$z = f(z) = A + C \int^z \frac{t dt}{(t^2-1)^{1/2}}$$

$$= A + C (z^2 - 1)^{1/2}$$

$1 \mapsto 0 \Rightarrow A = 0; 0 \mapsto i \Rightarrow C = 1$

so $z = (z^2 - 1)^{1/2}$

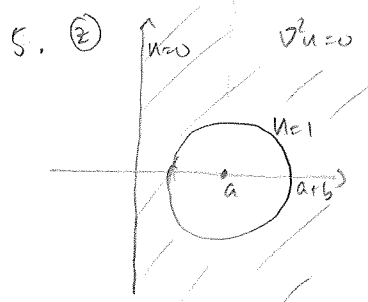
We'd like to find complex potential $w(z)$ satisfying $w \sim U, z$ as $z \rightarrow \infty$, $\text{Im } w = 0$ on boundary (with suitable choice of branch)

Converting to z plane, noting $z \sim z$ for large z , we want $w(z)$ with $w \sim U, z$ as $z \rightarrow \infty$, &

$\text{Im } w = 0$ on boundary, which is now just the real axis. So $w = U, z$ works everywhere.

Hence $w(z) = W(z) = U_1 (z^2 + 1)^{1/2}$

(Note we could work out from this the velocity potential and stream function, since $w = \phi + i\psi$)



Want to map 0 to an annulus. (There's no way we could map both boundaries to straight lines, since in that case they would 'knech' at ∞ , and they clearly don't knuch. If at least one of them has to remain circular, then it would obviously be easier if they are both circular, and concentric, so we could use polar coordinates in mapped domain)

- Consider mapping $s = \frac{z-\alpha}{z+\alpha}$, which, since Möbius, maps the boundaries to circles.
- For $z=iy$, $s = \frac{-(\alpha-iy)}{\alpha+iy}$ which has unit modulus, so the imaginary axis maps to the unit circle.
- For $z=a+be^{i\theta}$, $s = \frac{be^{i\theta}-(\alpha-a)}{be^{i\theta}+(\alpha+a)}$. If this is to map to a concentric circle $|s|=R$, we need

$$s\bar{s} = R^2 = \frac{(be^{i\theta}-(\alpha-a))(be^{-i\theta}-(\alpha-a))}{(be^{i\theta}+(\alpha+a))(be^{-i\theta}+(\alpha+a))} = \frac{b^2 - 2(\alpha-a)b\cos\theta + (\alpha-a)^2}{b^2 + 2(\alpha+a)b\cos\theta + (\alpha+a)^2} = -\frac{(\alpha-a)}{(\alpha+a)} + \frac{b^2 + \frac{\alpha-a}{\alpha+a}b^2 + (\alpha-a)^2 + \frac{\alpha-a}{\alpha+a}(\alpha+a)^2}{b^2 + 2(\alpha+a)b\cos\theta + (\alpha+a)^2}$$

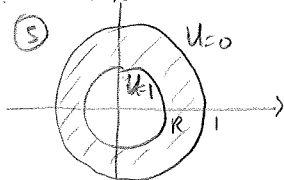
A to ensure this quantity is independent of θ , we must have the last numerator equal to zero.

ie. $2\alpha b^2 + 2\alpha(\alpha^2 - a^2) = 0 \Rightarrow$ we need $\boxed{\alpha^2 = a^2 - b^2}$

So with this choice of α , the circle $|z-a|=b$ maps to the circle $\boxed{|s|=R = \left(\frac{a-\alpha}{a+\alpha}\right)^{1/2} < 1}$

(Note $R = \frac{a-\alpha}{b} = \frac{a-\sqrt{a^2-b^2}}{b}$)

- So $s = \frac{z-\alpha}{z+\alpha}$ maps to



Solution of Laplace in polar coordinates
($\nabla^2 u = \frac{1}{r} \frac{d}{dr} \left(r \frac{du}{dr} \right) = 0$) are $u = A \log r + B$.

The solution here has $u = \frac{\log r}{\log R}$ (satisfies $u=0$ on $r=1$, $u=1$ on $r=R$)
 $= \operatorname{Re} \left(\frac{\log s}{\log R} \right)$

Hence $u = \operatorname{Re} \left(\frac{\log \left(\frac{z-\alpha}{z+\alpha} \right)}{\log \left(\frac{a-\alpha}{a+\alpha} \right)^{1/2}} \right) = \frac{\log \left| \frac{z-\sqrt{a^2-b^2}}{z+\sqrt{a^2-b^2}} \right|}{\log \left| \frac{a-\sqrt{a^2-b^2}}{a+\sqrt{a^2-b^2}} \right|}$

6. (a) $z = \cosh^{-1}(Z) \Leftrightarrow Z = \cosh z = \frac{1}{2}(e^z + e^{-z}) \Leftrightarrow e^{2z} - 2Ze^z + 1 = 0 \Leftrightarrow (e^z - Z)^2 = Z^2 - 1$

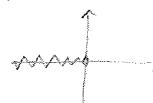
$\Leftrightarrow e^z = Z + (Z^2 - 1)^{1/2} \Leftrightarrow z = \log(Z + (Z^2 - 1)^{1/2})$

Define $(Z^2 - 1)^{1/2} = |Z^2 - 1|^{1/2} e^{i(\theta_1 + \theta_2)/2}$ where $\theta_1 = \arg(Z-1) \in (-\pi, \pi)$

$\theta_2 = \arg(Z+1) \in [-\pi, \pi]$



and take $\log(s) = \log|s| + i\theta$ where $\theta = \arg(s) \in (-\pi, \pi)$

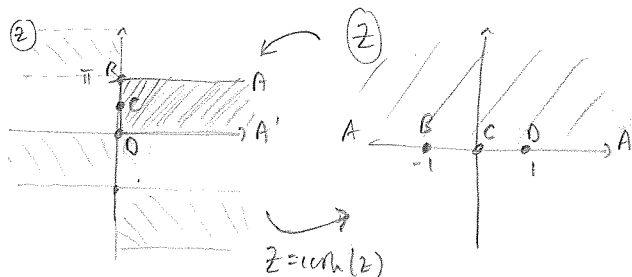


Then $\boxed{\cosh^{-1}(Z) = \log(Z + (Z^2 - 1)^{1/2})}$ is holomorphic in the upper half plane.

With this definition, $Z=0$ has $\theta_1 = \pi, \theta_2 = 0$ so $(Z^2 - 1)^{1/2} = e^{i\pi/2} = i$, so $Z + (Z^2 - 1)^{1/2} = i$, which has $\theta = \pi/2$.

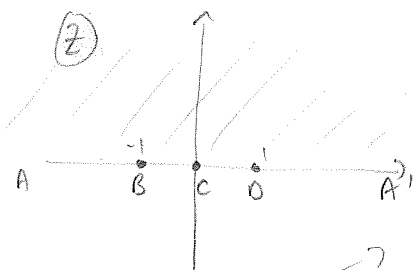
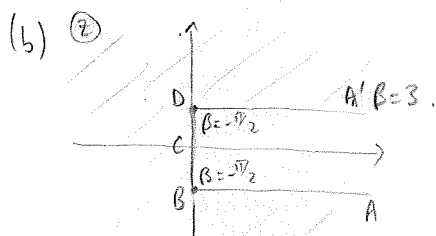
Hence $\cosh^{-1}(0) = i\pi/2$.

$\frac{d}{dZ} \cosh^{-1}(Z) = \frac{1 + \frac{Z}{(Z^2 - 1)^{1/2}}}{Z + (Z^2 - 1)^{1/2}} = \frac{1}{(Z^2 - 1)^{1/2}}$



$= \cosh x \cos y - i \sinh x \sin y$

- Note, the branch can also be 'defined' graphically. There are many regions of the Z plane that map to the upper half z plane under $Z = \cosh z$. The choice of branch is equivalent to deciding which of these we want to map back to.



$z = A + C \int (t^2 - 1)^{1/2} dt$

$= A + C \left(\frac{1}{2} (Z(Z^2 - 1)^{1/2} - \cosh^{-1} Z) \right)$

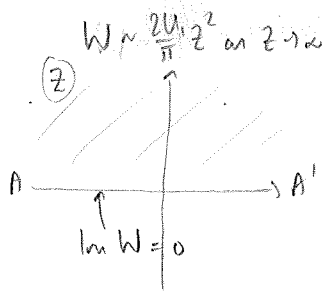
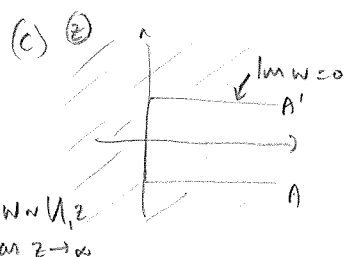
using the result below, and (a)

To map C to C, take $A + C \left(\frac{1}{2} \frac{i\pi}{4} \right) = 0$

D to D, take $A = i \Rightarrow C = \frac{4}{\pi}$

so $\boxed{z = i + \frac{2}{\pi} \left(Z(Z^2 - 1)^{1/2} - \cosh^{-1} Z \right)}$

NOTE $\int (t^2 - 1)^{1/2} dt = \int \frac{t^2}{(t^2 - 1)^{1/2}} dt - \int \frac{dt}{(t^2 - 1)^{1/2}}$
 $= t(t^2 - 1)^{1/2} - \int (t^2 - 1)^{1/2} dt - \int \frac{dt}{(t^2 - 1)^{1/2}}$



For large Z , $\cosh^{-1} Z$ is only logarithmic, so

$z \sim \frac{2}{\pi} Z^2$

Hence $w \sim U_1 Z \Rightarrow W \sim U_1 \frac{2}{\pi} Z^2$ for far field.

By inspection $\boxed{W = \frac{2U_1}{\pi} Z^2}$ satisfies the required conditions

Hence $\boxed{W(Z) = W(Z(z))}$ this inverse mapping does not look like it has a nice form.

We could write implicitly as $z = i + \frac{2}{\pi} \left[\left(\frac{\pi W}{2U_1} \right)^{1/2} \left(\frac{\pi W}{2U_1} - 1 \right)^{1/2} - \cosh^{-1} \left(\frac{\pi W}{2U_1} \right) \right]$