

As $x \rightarrow \infty$

$$\int_0^{\pi/2} e^{ix \cos t} dt - \text{method of stationary phase for first term} \\ (\text{then method of steepest descents for more...})$$

$$\int_0^1 \ln t e^{ixt} dt - \text{method of steepest descents} \\ (\text{note that using IBPs and the method of} \\ \text{stationary phase does not work...})$$

$$\int_0^x t^{-\frac{1}{2}} e^{-t} dt - \text{write as } \int_0^{\infty} t^{-\frac{1}{2}} e^{-t} dt - \int_x^{\infty} t^{-\frac{1}{2}} e^{-t} dt \text{ and then} \\ \text{use IBPs for the second integral}$$

$$\int_0^{\pi/2} e^{-x \sin^2 t} dt - \text{Laplace's method}$$

$$\int_0^1 e^{ix e^{-1/t}} dt - \text{method of steepest descents with } s = e^{-1/t}$$

As $x \rightarrow 0^+$

$$\int_0^{10} \frac{e^{-xt}}{1+t} dt - \text{Taylor expand the integrand and integrate term-by-term}$$

$$\int_0^{\pi/2} \frac{1}{\sqrt{\cos^2 t + x \sin^2 t}} dt - \text{write as } \int_0^{\pi/2 - \delta} + \int_{\pi/2 - \delta}^{\pi/2} \text{ where } x \ll \delta \ll 1$$

$$\int_0^1 \frac{\sin(xt)}{t} dt - \text{Taylor expand and integrate term-by-term.}$$

$$\int_x^{\infty} t^{a-1} e^{-t} dt - \text{write as } \int_0^{\infty} - \int_0^x \text{ and then Taylor expand} \\ \text{and integrate term-by-term for the second} \\ \text{integral when } \operatorname{Re}(a) > 0. (\text{NB v. tricky o/w!})$$

$$\int_0^1 \frac{\ln t}{x+t} dt - \text{write as } \int_0^{\delta} + \int_{\delta}^1 \text{ where } x \ll \delta \ll 1.$$

PS2 Q2 (as $x \rightarrow \infty$)

(2)

$$I_1(x) = \int_{-1}^1 e^{-x \cos ht} dt = \underbrace{\int_{-1}^{-\varepsilon} e^{-x \cos ht} dt}_{I_{11}} + \underbrace{\int_{-\varepsilon}^{\varepsilon} e^{-x \cos ht} dt}_{I_{12}} + \underbrace{\int_{\varepsilon}^1 e^{-x \cos ht} dt}_{I_{13} = I_1 \text{ (by symmetry)}}$$

$\varepsilon \ll 1$

- Split integral this way because $\cos ht$ has a maximum at $t=0$.

$$I_{11}(x) = \int_{\varepsilon}^1 e^{-x \cos ht} dt = O(e^{-x \cos h \varepsilon}) = O(e^{-x} e^{-x \varepsilon^2/2})$$

$$\begin{aligned} I_{12}(x) &= \int_{-\varepsilon}^{+\varepsilon} e^{-x \cos ht} dt = \int_{-\varepsilon}^{+\varepsilon} e^{-x(1 + \frac{1}{2}t^2 + O(t^4))} dt \\ &= e^{-x} \int_{-\varepsilon}^{+\varepsilon} e^{-\frac{1}{2}xt^2} e^{O(xt^4)} dt \quad \left. \begin{array}{l} \text{OK since } \varepsilon \ll 1 \\ \text{OK provided } xt^4 \ll 1 \end{array} \right\} \\ &= e^{-x} \int_{-\varepsilon}^{+\varepsilon} e^{-\frac{1}{2}xt^2} [1 + O(xt^4)] dt \\ &= e^{-x} \int_{-\varepsilon \sqrt{\frac{x}{2}}}^{+\varepsilon \sqrt{\frac{x}{2}}} e^{-s^2} [1 + O(s^4/x)] \cdot \sqrt{\frac{2}{x}} ds \quad \left. \begin{array}{l} s = \sqrt{\frac{x}{2}} t \\ \text{OK provided } \varepsilon \sqrt{\frac{x}{2}} \gg 1 \end{array} \right\} \\ &= e^{-x} \sqrt{\frac{2}{x}} \left\{ \int_{-\varepsilon \sqrt{\frac{x}{2}}}^{+\varepsilon \sqrt{\frac{x}{2}}} e^{-s^2} ds + O\left(\frac{1}{x}\right) \right\} \\ &= e^{-x} \sqrt{\frac{2}{x}} \left\{ \int_{-\infty}^{\infty} e^{-s^2} ds + O\left(\frac{1}{x}\right) \right\} \\ &= e^{-x} \sqrt{\frac{2}{x}} \left\{ \sqrt{\pi} + O\left(\frac{1}{x}\right) \right\} \text{ as } x \rightarrow \infty \end{aligned}$$

Then for $\varepsilon \sqrt{x} \gg 1$ we have $I_{11}(x) \ll I_{12}(x)$ as $x \rightarrow \infty$.

$$\therefore I_1(x) \sim \sqrt{\frac{2\pi}{x}} e^{-x} \text{ as } x \rightarrow \infty$$

NB we need to select ε s.t. $x^{-\frac{1}{2}} \ll \varepsilon \ll x^{-\frac{1}{4}}$

$$I_2(x) = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} e^{-x(t^2 - \sin^2 t)} dt$$

③

Let $\phi(t) = \sin^2 t - t^2$. Then $\phi(0) = 0$ and $\phi(t) < 0$ for $t \neq 0$. Hence $\phi(t)$ has an interior maximum at $t = 0$. Note that

$\phi'(t) = 2 \sin t \cos t - 2t$ and so $\phi'(0) = 0 \Rightarrow$ degenerate case.

$$\phi(t) = \left(t - \frac{1}{3}t^3 + o(t^5)\right)^2 - t^2 = -\frac{1}{3}t^4 + o(t^6) \text{ as } t \rightarrow 0.$$

Split the integral as

$$I_2(x) = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \sim dt = \underbrace{\int_{-\frac{\pi}{2}}^{-\varepsilon} \sim dt}_{I_{21}} + \underbrace{\int_{-\varepsilon}^{+\varepsilon} \sim dt}_{I_{22}} + \underbrace{\int_{+\varepsilon}^{\frac{\pi}{2}} \sim dt}_{I_{23} = I_{21} \text{ by symmetry.}} \text{ (where } \varepsilon \ll 1)$$

$$I_{21}(x) = \int_{+\varepsilon}^{\frac{\pi}{2}} e^{x\phi(t)} dt = o(e^{x\phi(\varepsilon)}) = o(e^{-x\varepsilon^4})$$

$$\begin{aligned} I_{22}(x) &= \int_{-\varepsilon}^{+\varepsilon} e^{-x\left(\frac{1}{3}t^4 + o(t^6)\right)} dt \quad \leftarrow \text{ok since } \varepsilon \ll 1 \\ &= \int_{-\varepsilon}^{+\varepsilon} e^{-\frac{1}{3}xt^4} (1 + o(xt^6)) dt \quad \left. \begin{array}{l} \text{ok provided } xt^6 \ll 1 \\ \\ \text{provided } \left(\frac{x}{3}\right)^{\frac{1}{4}} \varepsilon \gg 1 \end{array} \right\} \\ &= \int_{-\left(\frac{x}{3}\right)^{\frac{1}{4}} \varepsilon}^{+\left(\frac{x}{3}\right)^{\frac{1}{4}} \varepsilon} e^{-s^4} (1 + o(x^{-\frac{1}{2}}s^6)) \cdot \left(\frac{3}{x}\right)^{\frac{1}{4}} ds \\ &= \left(\frac{3}{x}\right)^{\frac{1}{4}} \left\{ \int_{-\infty}^{+\infty} e^{-s^4} ds + o(x^{-\frac{1}{2}}) \right\} \\ &= \left(\frac{3}{x}\right)^{\frac{1}{4}} \left\{ \frac{2\Gamma(1/4)}{4} + o(x^{-\frac{1}{2}}) \right\} \quad \left. \begin{array}{l} \\ \\ \text{by hint on problem sheet} \end{array} \right\} \end{aligned}$$

For $x^{\frac{1}{4}}\varepsilon \gg 1$ then we have $I_{21}(x) \ll I_{22}(x)$ as $x \rightarrow \infty$

$$\therefore I_2(x) \sim \frac{\Gamma(1/4)}{2} \cdot \left(\frac{3}{x}\right)^{\frac{1}{4}} \text{ as } x \rightarrow \infty$$

NB we need to select ε s.t. $x^{-\frac{1}{4}} \ll \varepsilon \ll x^{-\frac{1}{6}}$

$$I_3(x) = \int_0^{\infty} e^{-2t - x/t^2} dt$$

$$\frac{d}{dt} \left(2t + \frac{x}{t^2} \right) = 2 - \frac{2x}{t^3} \Rightarrow \text{the maximum moves as } x \text{ is varied.}$$

\Rightarrow Make a change of variables: let $y = x^{\frac{1}{3}}$ and $t = x^{\frac{1}{3}}u$ so that

$$I_3(x) = \int_0^{\infty} e^{-y \underbrace{(2u + \frac{1}{u^2})}_{\varphi(u)}} y du$$

$$\text{Then } \varphi(u) = -2u + \frac{1}{u^2} \Rightarrow \varphi'(u) = -2 - \frac{2}{u^3} \text{ and } \varphi'(u) = 0 \text{ for } u = 1.$$

$$\begin{aligned} \text{Expand } \varphi(u) \text{ about } u=1: \varphi(u) &= \varphi(1) + (u-1)\varphi'(1) + \frac{1}{2}(u-1)^2\varphi''(1) + o((u-1)^3) \\ &= -3 - 3(u-1)^2 + o((u-1)^3) \end{aligned}$$

Split the range of integration:

$$I_3(y) = \underbrace{\int_0^{1-\varepsilon} \sim du}_{I_{31}} + \underbrace{\int_{1-\varepsilon}^{1+\varepsilon} \sim du}_{I_{32}} + \underbrace{\int_{1+\varepsilon}^{\infty} \sim du}_{I_{33}} \quad (\varepsilon \ll 1)$$

$$I_{31}(y) = \int_0^{1-\varepsilon} y e^{-y(2u + 1/u^2)} du = o(y e^{-y\varphi(1-\varepsilon)}) = o(y e^{-3y} e^{-3y\varepsilon^2})$$

$$I_{33}(y) = \int_{1+\varepsilon}^{\infty} y e^{-y(2u + 1/u^2)} du = o(y e^{-y\varphi(1+\varepsilon)}) = o(y e^{-3y} e^{-3y\varepsilon^2})$$

$$\begin{aligned} I_{32}(y) &= \int_{1-\varepsilon}^{1+\varepsilon} y e^{-y(3 + 3(u-1)^2 + o((u-1)^3))} du \quad (\text{OK since } \varepsilon \ll 1) \\ &= y e^{-3y} \int_{1-\varepsilon}^{1+\varepsilon} e^{-3y(u-1)^2} (1 + o(y(u-1)^3)) du \quad \downarrow \text{provided } y\varepsilon^3 \ll 1 \\ &= y e^{-3y} \int_{-\varepsilon\sqrt{3y}}^{+\varepsilon\sqrt{3y}} e^{-s^2} (1 + o(y^{-\frac{1}{2}}s^3)) \frac{1}{\sqrt{3y}} ds \quad \downarrow \sqrt{3y}(u-1) = s \\ &= \sqrt{\frac{y}{3}} e^{-3y} \left\{ \int_{-\infty}^{\infty} e^{-s^2} ds + o(y^{-\frac{1}{2}}) \right\} \quad \downarrow \text{provided } \varepsilon\sqrt{3y} \gg 1 \end{aligned}$$

Then, for $\varepsilon \sqrt{y} \gg 1$ we have $I_{31}(y), I_{33}(y) \ll I_{32}(y)$ as $y \rightarrow \infty$. ⑤

$$\text{Hence } I_3(y) \sim \sqrt{\frac{y}{3}} e^{-3y} \cdot \sqrt{\pi}$$

$$= \sqrt{\frac{\pi}{3}} x^{1/6} e^{-3x^{1/3}} \quad \text{as } x \rightarrow \infty.$$

NB we need to select ε s.t. $x^{-1/6} \ll \varepsilon \ll x^{-1/9}$.

PS2 Q3

$$J_1(x) = \int_0^1 e^{ixt^2} \cosh(t^2) dt \quad \text{as } x \rightarrow \infty.$$

Then $\psi(t) = t^2$ and $\psi'(t) = 2t \Rightarrow \psi'(t) = 0$ for $t = 0$.

Hence split the region of integration as

$$J_1(x) = \underbrace{\int_0^\varepsilon \dots dt}_{J_{11}} + \underbrace{\int_\varepsilon^1 \dots dt}_{J_{12}} \quad \text{for } \varepsilon \ll 1.$$

$$J_{11}(x) = \int_0^\varepsilon (1 + o(t^4)) e^{ixt^2} dt \quad (\text{OK since } \varepsilon \ll 1)$$

$$= \int_0^{\varepsilon\sqrt{x}} (1 + o(s^4/x^2)) e^{is^2} \cdot \frac{1}{\sqrt{x}} ds \quad \downarrow \quad s = \sqrt{x}t$$

$$= \frac{1}{\sqrt{x}} \int_0^\infty e^{is^2} ds + o\left(\frac{1}{\sqrt{x} \cdot \varepsilon\sqrt{x}}\right) \quad \downarrow \quad \text{provided } \varepsilon\sqrt{x} \gg 1$$

*

NB we can do $\int_0^{\varepsilon\sqrt{x}} \rightarrow \int_0^\infty$ because

$$\begin{aligned} \frac{1}{\sqrt{x}} \int_{\varepsilon\sqrt{x}}^\infty e^{is^2} ds &= \frac{1}{\sqrt{x}} \int_{\varepsilon\sqrt{x}}^\infty \underbrace{\frac{1}{2is}}_u \cdot \underbrace{2is e^{is^2}}_{dv/ds} ds \\ &= \frac{1}{\sqrt{x}} \left\{ \left[\frac{1}{2is} e^{is^2} \right]_{\varepsilon\sqrt{x}}^\infty - \int_{\varepsilon\sqrt{x}}^\infty -\frac{1}{2is^2} e^{is^2} ds \right\} \\ &= \underbrace{\frac{1}{2i\varepsilon x} e^{i\varepsilon^2 x}}_{o\left(\frac{1}{\varepsilon x}\right)} - \underbrace{\frac{1}{2i\sqrt{x}} \int_{\varepsilon\sqrt{x}}^\infty \frac{1}{s^2} e^{is^2} ds}_{o\left(\frac{1}{\varepsilon x}\right)} \end{aligned}$$

Further, * comes from

$$\frac{1}{x^{5/2}} \int_0^{\varepsilon\sqrt{x}} s^4 e^{is^2} ds = o\left(\frac{(\varepsilon\sqrt{x})^3}{x^{5/2}}\right) = o\left(\frac{\varepsilon^3}{x}\right) \ll 1 \quad \text{provided } \varepsilon^3 \ll x.$$

→ from IBPs and looking at the first term...

$$J_{1/2}(x) = \int_{\varepsilon}^1 \underbrace{\frac{\cosh(t^2)}{2ixt}}_u \cdot \underbrace{2ixt e^{ixt^2}}_{dv/dt} dt$$

$$= \left[\frac{\cosh(t^2)}{2ixt} e^{ixt^2} \right]_{\varepsilon}^1 - \int_{\varepsilon}^1 \frac{\partial}{\partial t} \left(\frac{\cosh(t^2)}{2ixt} \right) e^{ixt^2} dt$$

$$= O\left(\frac{1}{\varepsilon x}\right) \qquad = O\left(\frac{1}{x}\right) \text{ as } x \rightarrow \infty \text{ by RLL}$$

Hence $J_{1/2}(x) \sim \frac{e^{i\pi/4}}{2} \sqrt{\frac{\pi}{x}}$ as $x \rightarrow \infty$

NB we need to select $x^{-\varepsilon} \ll \varepsilon \ll x^{-\frac{1}{3}}$ ($\Rightarrow x^{-\frac{1}{2}} \gg \frac{1}{\varepsilon x}$).

$$J_2(x) = \text{Re} [J_4(x)] = \text{Re} \left[\int_0^1 \tan t e^{ixt^4} dt \right] \text{ as } x \rightarrow \infty$$

Then $\psi(t) = t^4$ has $\psi'(0) = 0$, so we split the region of integration as:

$$J_4(x) = \underbrace{\int_0^\epsilon \sim dt}_{J_{41}} + \underbrace{\int_\epsilon^1 \sim dt}_{J_{42}} \quad (\epsilon \ll 1)$$

$$\begin{aligned} J_{41}(x) &= \int_0^\epsilon (t + o(t^3)) e^{ixt^4} dt \quad (\text{OK since } \epsilon \ll 1) \\ &= \int_0^{x^{1/4}\epsilon} \left(\frac{s}{x^{1/4}} + o\left(\frac{s^3}{x^{3/4}}\right) \right) e^{is^4} \cdot x^{-1/4} ds \quad \downarrow s = x^{1/4}t \\ &= x^{-1/2} \int_0^\infty s e^{is^4} ds + o\left(\frac{1}{\epsilon^2 x}\right) \quad \downarrow \text{provided } x^4 \epsilon \gg 1 \end{aligned}$$

NB the above holds since

$$\begin{aligned} x^{-1/2} \int_{x^{1/4}\epsilon}^\infty s e^{is^4} ds &= x^{-1/2} \int_{x^{1/4}\epsilon}^\infty \underbrace{\frac{1}{4is^2}}_u \cdot \underbrace{4is^3 e^{is^4}}_{dv/ds} ds \\ &= x^{-1/2} \left\{ \underbrace{\left[\frac{1}{4is^2} e^{is^4} \right]_{x^{1/4}\epsilon}^\infty}_{= o\left(\frac{1}{x^{1/2}\epsilon^2}\right)} - \underbrace{\int_{x^{1/4}\epsilon}^\infty \frac{-2}{4is^3} e^{is^4} ds}_{= o\left(\frac{1}{x^{1/2}\epsilon^2}\right)} \right\} \\ &= o\left(\frac{1}{x\epsilon^2}\right) \end{aligned}$$

and also

$$\frac{1}{x} \int_0^{x^{1/4}\epsilon} s^3 e^{is^4} ds = \frac{1}{4ix} [e^{is^4}]_0^{x^{1/4}\epsilon} = o\left(\frac{1}{x}\right)$$

$$J_{42}(x) = \int_\epsilon^1 \underbrace{\frac{\tan t}{4ixt^3}}_u \cdot \underbrace{4ixt^3 e^{ixt^4}}_{dv/dt} dt$$

$$J_{4,2}(x) = \left[\frac{\tan t}{4ixt^3} e^{ixt^4} \right]_{\varepsilon}^1 - \int_{\varepsilon}^1 \underbrace{\frac{\partial}{\partial t} \left(\frac{\tan t}{4ixt^3} \right) e^{ixt^4}}_{= O\left(\frac{1}{x}\right) \text{ as } x \rightarrow \infty \text{ by RL}} dt$$

$$= O\left(\frac{1}{x\varepsilon^2}\right)$$

Hence

$$J_4 \sim x^{-\frac{1}{2}} \int_0^{\infty} s e^{is^4} ds = \frac{x^{-\frac{1}{2}} e^{i\pi/4} \Gamma\left(\frac{1}{2}\right)}{4}$$

$$\therefore J_2(x) \sim \frac{\cos\left(\frac{\pi}{4}\right) \Gamma\left(\frac{1}{2}\right)}{4} x^{-\frac{1}{2}} \quad \text{as } x \rightarrow \infty.$$

$$J_3(x) = \int_0^1 e^{ix(t - \sin t)} dt$$

Let $\psi(t) = t - \sin t \Rightarrow \psi'(t) = 1 - \cos t$ and hence $t=0$ is the only stationary point. We have $\psi(t) = t - [t - \frac{1}{3}t^3 + O(t^5)]$

$$= \frac{1}{3!}t^3 + O(t^5)$$

Hence we split the region of integration:

$$J_3(x) = \underbrace{\int_0^\varepsilon \sim dt}_{J_{31}} + \underbrace{\int_\varepsilon^1 \sim dt}_{J_{32}}$$

$$J_{31}(x) = \int_0^\varepsilon e^{ix[\frac{1}{6}t^3 + O(t^5)]} dt \quad (\text{since } \varepsilon \ll 1)$$

$$= \int_0^{\varepsilon(\frac{x}{6})^{1/3}} e^{is^3} [1 + O(x(\frac{s}{x^{1/3}})^5)] (\frac{6}{x})^{1/3} ds \quad \downarrow \quad s = (\frac{x}{6})^{1/3} t$$

(provided $x\varepsilon^5 \ll 1$)

$$= (\frac{6}{x})^{1/3} \int_0^\infty e^{is^3} ds + O(\frac{1}{\varepsilon^2 x}) \quad (\text{provided } x^{1/3} \varepsilon \gg 1)$$

NB the above holds since

$$\begin{aligned} x^{-1/3} \int_{\varepsilon(\frac{x}{6})^{1/3}}^\infty e^{is^3} ds &= x^{-1/3} \int_{\varepsilon(\frac{x}{6})^{1/3}}^\infty \underbrace{\frac{1}{3is^2}}_u \cdot \underbrace{3is^2 e^{is^3}}_{dv/ds} ds \\ &= x^{-1/3} \left\{ \left[\frac{1}{3is^2} e^{is^3} \right]_{\varepsilon(\frac{x}{6})^{1/3}}^\infty - \int_{\varepsilon(\frac{x}{6})^{1/3}}^\infty \frac{\partial}{\partial s} \left(\frac{1}{3is^2} \right) e^{is^3} ds \right\} \\ &= O\left(\frac{1}{\varepsilon^2 x}\right) \end{aligned}$$

and also

$$\begin{aligned} \frac{1}{x} \int_0^{\varepsilon(\frac{x}{6})^{1/3}} s^5 e^{is^3} ds &= \frac{1}{x} \int_0^{\varepsilon(\frac{x}{6})^{1/3}} \underbrace{\frac{s^3}{3i}}_u \cdot \underbrace{3is^2 e^{is^3}}_{dv/ds} ds \\ &= \frac{1}{x} \left\{ \left[\frac{s^3}{3i} e^{is^3} \right]_0^{\varepsilon(\frac{x}{6})^{1/3}} - \int_0^{\varepsilon(\frac{x}{6})^{1/3}} \frac{s^2}{i} e^{is^3} ds \right\} \\ &= O(\varepsilon^3) + O\left(\frac{1}{x}\right) \end{aligned}$$

$$\begin{aligned}\therefore J_3(x) &\sim \left(\frac{6}{x}\right)^{1/3} \int_0^\infty e^{is^3} ds \\ &= \left(\frac{6}{x}\right)^{1/3} \frac{e^{i\pi/6} \Gamma\left(\frac{1}{3}\right)}{3} \\ &= \left(\frac{2}{9}\right)^{1/3} \Gamma\left(\frac{1}{3}\right) e^{i\pi/6} x^{-1/3} \text{ as } x \rightarrow \infty\end{aligned}$$

NB we need to select $x^{-1/3} \ll \epsilon \ll x^{-1/3}$.

$$I(x) = \int_{-1}^1 (1-t^2)^N e^{ixt} dt$$

N -integer, and contour of integration is a line segment from $-1 \rightarrow 1$.

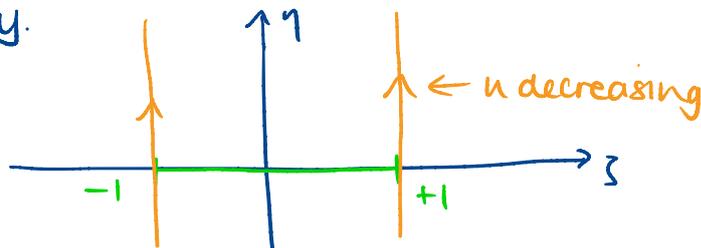
(a) $\varphi(t) = it = i(\zeta + i\eta) \Rightarrow u(\zeta, \eta) = -\eta, v(\zeta, \eta) = \zeta$

On a steepest descent contour: $v(\zeta, \eta) = \zeta = \text{constant}$

$\nabla u = \begin{pmatrix} 0 \\ -1 \end{pmatrix}$ is tangent to the steepest descent contour, and

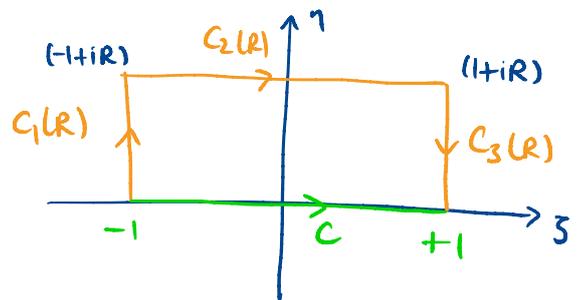
$-\nabla u = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ is a tangent in the direction in which u decreases

most rapidly.



(b) Since the integrand is holomorphic in $t \in \mathbb{C}$, we can deform the contour to give

$$I(x) = \left\{ \int_{C_1} + \int_{C_2} + \int_{C_3} \right\} f(t) e^{x\varphi(t)} dt$$



On $C_2(R)$: $|f(t)e^{x\varphi(t)}| = |f(t)|e^{-xR} \rightarrow 0$ as $R \rightarrow \infty$.

$$\therefore I(x) = \int_{-1}^{-1+i\infty} (1-t^2)^N e^{ixt} dt + \int_{1+i\infty}^1 (1-t^2)^N e^{ixt} dt$$

$$= I_-(x) + I_+(x)$$

(c) $I_{\pm}(x) = \int_{\pm 1}^{\pm 1+i\infty} f(\pm 1+is) e^{ix(\pm 1+is)} i ds \quad s > 0, t = \pm 1+is$

$$= ie^{\pm ix} \int_{\pm 1}^{\pm 1+i\infty} f(\pm 1+is) e^{-xs} ds$$

$$\text{where } f(\pm 1+is) = (1-(\pm 1+is)^2)^N = \begin{cases} (-is)^N (2+is)^N & \oplus \\ (2+is)^N (is)^N & \ominus \end{cases}$$

Using Laplace's method: we write

$$I_{\pm}(x) = ie^{\pm ix} \left\{ \int_0^{\varepsilon} + \int_{\varepsilon}^{\infty} \right\} f(\pm 1 + is) e^{-xs} ds \quad (\varepsilon \ll 1)$$

Consider I_+ first:

$$\begin{aligned} \int_0^{\varepsilon} f(1+is) e^{-xs} ds &= \int_0^{\varepsilon} (-is)^N (2+is)^N e^{-xs} ds \\ &= (-2i)^N \int_0^{\varepsilon} (s^N + o(s^{N+1})) e^{-xs} ds \quad \left. \begin{array}{l} \text{since} \\ \varepsilon \ll 1. \end{array} \right\} \\ &= \frac{(-2i)^N}{x^{N+1}} \left\{ \int_0^{\varepsilon x} u^N e^{-u} du + o\left(\frac{1}{x}\right) \right\} \quad \left. \begin{array}{l} \text{letting} \\ s = \frac{u}{x} \end{array} \right\} \\ &= \frac{(-2i)^N}{x^{N+1}} \left\{ \underbrace{\int_0^{\infty} u^N e^{-u} du}_{N!} + o\left(\frac{1}{x}\right) \right\} \quad \left. \begin{array}{l} \text{provided} \\ \varepsilon x \gg 1 \end{array} \right\} \\ &= \frac{(-2i)^N N!}{x^{N+1}} + o\left(\frac{1}{x^{N+2}}\right) \end{aligned}$$

$$\int_{\varepsilon}^{\infty} f(1+is) e^{-xs} ds = o(e^{-\varepsilon x})$$

$$\therefore I_+(x) \sim \frac{ie^{ix} (-2i)^N N!}{x^{N+1}} \text{ as } x \rightarrow \infty$$

NB Need to select $x^{-1} \ll \varepsilon \ll 1$.

$$\text{Similarly, } I_-(x) \sim \frac{ie^{-ix} (2i)^N N!}{x^{N+1}} \text{ as } x \rightarrow \infty$$

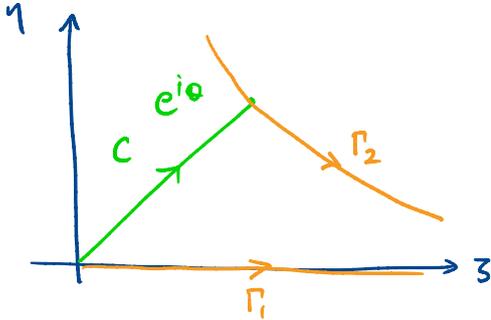
Hence

$$I(x) \sim \frac{2^N i^{N+1} N!}{x^{N+1}} (e^{-ix} - (-1)^N e^{ix}) \text{ as } x \rightarrow \infty$$

(NB - this is real - as required by symmetry!).

$$\operatorname{erf}(z) = \frac{z}{\sqrt{\pi}} \int_0^z e^{-z^2} dz = \frac{2r}{\sqrt{\pi}} \int_0^{e^{i\theta}} e^{-r^2 t^2} dt \quad \text{with } z = re^{i\theta} \text{ and } s = rt$$

$$\phi(t) = -t^2 = -(z+i\eta)^2 = \underbrace{\eta^2 - z^2}_{u(z,\eta)} - \underbrace{2z\eta i}_{v(z,\eta)}$$



Contour of steepest descent through $(0,0)$ is $\eta = 0$.

Contour of steepest descent through $t = e^{i\theta}$ is $2z\eta = 2\cos\theta \sin\theta = \sin 2\theta$
 $(\theta \in (0, \pi/2) \Rightarrow z, \eta > 0)$

Then, by the deformation theorem, $\operatorname{erf}(z) = \left(\int_{\Gamma_1} - \int_{\Gamma_2} \right) \frac{2r}{\sqrt{\pi}} e^{r^2 \phi(t)} dt$
 $\Gamma_1 = \Gamma_1(r) \quad \Gamma_2 = \Gamma_2(r, \theta)$

$$I_1(r) = \frac{2r}{\sqrt{\pi}} \int_0^\infty e^{-r^2 z^2} dz = \frac{2}{\sqrt{\pi}} \int_0^\infty e^{-u^2} du = 1$$

with $u = rz$

For I_2 , we have

$$\Gamma_2 = \left\{ z + i \frac{\sin 2\theta}{2z}, z > \cos\theta \right\}$$

$2z\eta = \sin 2\theta$
 $\Rightarrow \eta = \frac{\sin 2\theta}{2z}$

and $\underbrace{u}_{\eta^2 - z^2} + \underbrace{iv}_{-2z\eta i}$
 $\phi(t) = \eta^2 - z^2 - 2z\eta i$
 $= \eta^2 - z^2 - i \sin 2\theta$

which gives

$$I_2(r, \theta) = \frac{2r}{\sqrt{\pi}} \int_{\cos\theta}^\infty e^{r^2(\eta^2 - z^2 - i \sin 2\theta)} \cdot (1 + \eta'(z)i) dz$$

$$\eta'(z) = -\frac{\sin 2\theta}{2z^2}$$

$$= \frac{2r}{\sqrt{\pi}} e^{-r^2 i \sin 2\theta} \int_{\cos\theta}^\infty F(z) e^{r^2 \Phi(z)} dz$$

$$F(z) = 1 - \frac{i \sin 2\theta}{2z^2}$$

$$\Phi(z) = \frac{\sin^2 2\theta}{4z^2}$$

Since Γ_2 is a contour of steepest descent then

$\Phi(z)$ is a decreasing function of z on Γ_2

\Rightarrow Apply Laplace's method to give

$$I_2(r, \theta) = \frac{-2r}{\sqrt{\pi}} e^{-r^2 i \sin 2\theta} \frac{F(\cos \theta) e^{r^2 \Phi(\cos \theta)}}{r^2 \Phi'(\cos \theta)} \quad \text{as } r \rightarrow \infty.$$

$$F(\cos \theta) = \frac{e^{-i\theta}}{\cos \theta}, \quad \Phi(\cos \theta) = -\cos 2\theta, \quad \Phi'(\cos \theta) = \frac{-2}{\cos \theta}$$

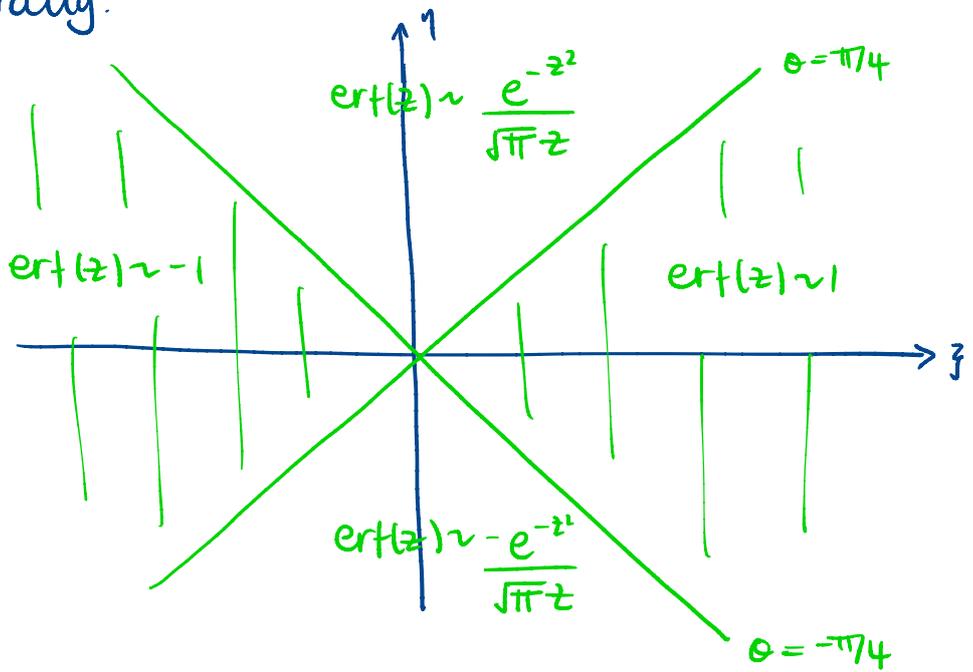
$$\Rightarrow I_2(r, \theta) \sim \frac{1}{\sqrt{\pi} r e^{i\theta}} e^{-r^2 e^{2i\theta}} \quad \text{as } r \rightarrow \infty.$$

Hence $I_1(r) \sim 1$ and $I_2(r, \theta) \sim \frac{1}{\sqrt{\pi} z} e^{-z^2}$ as $r = |z| \rightarrow \infty$
 for $0 < \theta = \arg(z) < \frac{\pi}{2}$

$$|I_2(r)| \sim \frac{1}{r} e^{-r^2 \cos 2\theta} = \begin{cases} \ll 1 & \text{for } 0 < \theta \leq \frac{\pi}{4} \\ \gg 1 & \text{for } \frac{\pi}{4} < \theta < \frac{\pi}{2} \end{cases} \quad \text{as } z \rightarrow \infty$$

$$\therefore \text{erf}(z) = \begin{cases} 1 & \text{for } 0 < \theta \leq \frac{\pi}{4} \\ \frac{1}{\sqrt{\pi} z} e^{-z^2} & \text{for } \frac{\pi}{4} < \theta < \frac{\pi}{2} \end{cases} \quad \text{as } |z| \rightarrow \infty.$$

More generally:



- Different asymptotic expansions in different regions → Stokes Phenomena
- While e^{-z^2} is active, it has an essential singularity at ∞ ↑
- $\theta = \pm \frac{\pi}{2}$ - Stokes' lines (across which topology of SD contour changes).
- $|\theta| = \frac{\pi}{4}, \frac{3\pi}{4}$ - anti Stokes' lines (across which dominance of end point and saddle point changes).

$$I(\varepsilon) = \int_0^1 \frac{f(x)}{x+\varepsilon} dx \quad \text{as } \varepsilon \rightarrow 0^+ \text{ with } f \text{ smooth.}$$

$$= \underbrace{\int_0^\delta \frac{f(x)}{x+\varepsilon} dx}_{I_1} + \underbrace{\int_\delta^1 \frac{f(x)}{x+\varepsilon} dx}_{I_2} \quad \text{where } 0 < \varepsilon \ll \delta \ll 1$$

$$I_1(\varepsilon) = \int_0^{\delta/\varepsilon} \frac{f(\varepsilon y)}{y+1} dy \quad (\text{letting } x = \varepsilon y)$$

$$= \int_0^{\delta/\varepsilon} \frac{1}{y+1} [f(0) + \varepsilon y f'(0) + o(\varepsilon^2)] dy$$

↓ since $\varepsilon y \ll \delta \ll 1$.

$$= [f(0) \ln(y+1)]_0^{\delta/\varepsilon} + o(\delta) \quad \downarrow \quad o(\varepsilon \cdot \frac{\delta}{\varepsilon})$$

$$= f(0) \ln\left(1 + \frac{\delta}{\varepsilon}\right) + o(\delta)$$

$$= f(0) \ln\left(\frac{\delta}{\varepsilon}\right) + f(0) \ln\left(1 + \frac{\varepsilon}{\delta}\right) + o(\delta)$$

$$= -f(0) \ln \varepsilon + f(0) \ln \delta + o\left(\delta, \frac{\varepsilon}{\delta}\right)$$

$$I_2(\varepsilon) = \int_\delta^1 \frac{f(x)}{x+\varepsilon} dx$$

$$= \int_\delta^1 \frac{f(x)}{x(1+\varepsilon/x)} dx$$

$$= \int_\delta^1 \frac{f(x)}{x} \left(1 - \frac{\varepsilon}{x} + o(\varepsilon^2)\right) dx$$

↓ since $\frac{\varepsilon}{x} < \frac{\varepsilon}{\delta} \ll 1$

$$= \int_\delta^1 \frac{f(x) - f(0)}{x} dx + \int_\delta^1 \frac{f(0)}{x} dx + \dots$$

$$= \int_\delta^1 \frac{f(x) - f(0)}{x} dx - f(0) \ln \delta + \dots$$

$$\therefore I(\varepsilon) \sim -f(0) \ln \varepsilon + f(0) \ln \delta + \int_\delta^1 \frac{f(x) - f(0)}{x} dx - f(0) \ln \delta + \dots$$

$$\sim -f(0) \ln \varepsilon + \int_0^1 \frac{f(x) - f(0)}{x} dx + \dots \quad \text{as } \varepsilon \rightarrow 0^+.$$