Numerical Solution of Partial Differential Equations

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Lecture 15

First-order hyperbolic equations

Let Ω be a bounded open set in $\mathbb{R}^n$, $n \geq 1$, with boundary $\Gamma = \partial\Omega$, and let $T > 0$. In $Q = \Omega \times (0, T]$, we consider the initial boundary-value problem

$$\frac{\partial u}{\partial t} + \sum_{i=1}^{n} b_i(x) \frac{\partial u}{\partial x_i} + c(x,t)u = f(x,t), \quad x \in \Omega, \ t \in (0, T], \quad (1)$$

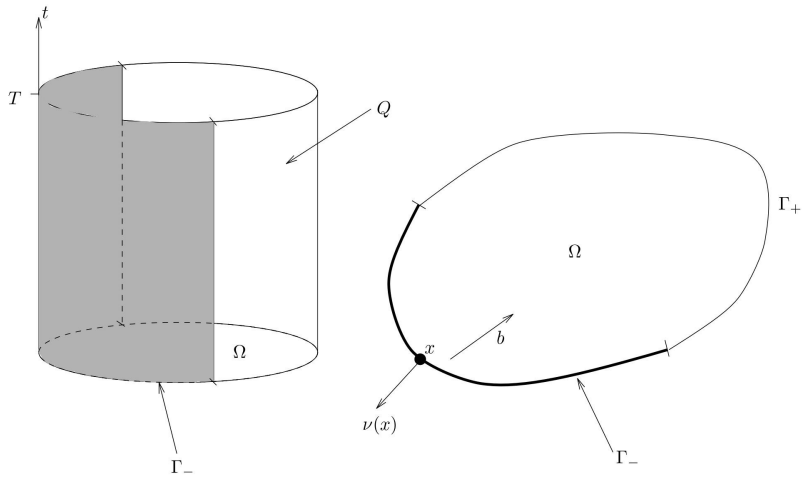$$u(x,t) = 0, \quad x \in \Gamma_-, \ t \in [0, T], \qquad (2)$$

$$u(x,0) = u_0(x), \quad x \in \overline{\Omega}, \qquad (3)$$

where

$$\Gamma_- = \{x \in \Gamma : b(x) \cdot \nu(x) < 0\},$$

$b = (b_1, \ldots, b_n)$ and $\nu(x)$ denotes the unit outward normal to Γ at $x \in \Gamma$. Γ_- will be called the inflow boundary.
Its complement, $\Gamma_+ = \Gamma \setminus \Gamma_-$, is referred to as the outflow boundary.

Continuous dependence of the solution on the data

We shall assume that

$$b_i \in C^1(\overline{\Omega}), \quad i = 1, \ldots, n,$$
$$c \in C(\overline{Q}), \quad f \in L_2(Q),$$
$$u_0 \in L_2(\Omega).$$

To ensure consistency between the initial condition and the boundary condition, we shall suppose that $u_0(x) = 0$, $x \in \Gamma_-$.

We make the additional hypothesis:

$$c(x, t) - \frac{1}{2} \sum_{i=1}^{n} \frac{\partial b_i}{\partial x_i}(x) \geq 0, \quad x \in \overline{\Omega}, \quad t \in [0, T]. \tag{4}$$

By taking the inner product in $L_2(\Omega)$ of the equation (1) with $u(\cdot, t)$, performing partial integration and noting the boundary condition (2):

$$\left(\frac{\partial u}{\partial t}(\cdot, t), u(\cdot, t)\right) + \left(c(\cdot, t) - \frac{1}{2}\sum_{i=1}^{n}\frac{\partial b_i}{\partial x_i}(\cdot), u^2(\cdot, t)\right)$$
$$+ \frac{1}{2}\int_{\Gamma_+}\left[\sum_{i=1}^{n}b_i(x)\nu_i(x)\right]u^2(x, t)\,\mathrm{d}s(x) = (f(\cdot, t), u(\cdot, t)), \tag{5}$$

where $\nu(x) = (\nu_1(x), \ldots, \nu_n(x))$ is the unit outward normal vector to Γ at $x \in \Gamma$.

Thanks to (4) and noting that

$$\left(\frac{\partial u}{\partial t}, u\right) = \int_\Omega \frac{\partial u}{\partial t}(x, t)\, u(x, t)\, \mathrm{d}x$$

$$= \int_\Omega \frac{1}{2}\frac{\partial}{\partial t}u^2(x, t)\, \mathrm{d}x = \frac{1}{2}\frac{\mathrm{d}}{\mathrm{d}t}\int_\Omega u^2(x, t)\, \mathrm{d}x$$

$$= \frac{1}{2}\frac{\mathrm{d}}{\mathrm{d}t}\|u(\cdot, t)\|^2,$$

it follows from (5) that

$$\frac{1}{2}\frac{\mathrm{d}}{\mathrm{d}t}\|u(\cdot, t)\|^2 + \frac{1}{2}\int_{\Gamma_+}\left[\sum_{i=1}^{n} b_i(x)\nu_i(x)\right] u^2(x, t)\, \mathrm{d}s(x) \leq (f, u).$$

By the Cauchy–Schwarz inequality,

$$(f, u) \leq \|f(\cdot, t)\|\, \|u(\cdot, t)\|$$

$$\leq \frac{1}{2}\|f(\cdot, t)\|^2 + \frac{1}{2}\|u(\cdot, t)\|^2,$$

and therefore, for any $t \in [0, T]$,

$$\frac{\mathrm{d}}{\mathrm{d}t}\|u(\cdot, t)\|^2 + \int_{\Gamma_+}\left[\sum_{i=1}^{n} b_i(x)\nu_i(x)\right] u^2(x, t)\, \mathrm{d}s(x) - \|u(\cdot, t)\|^2 \leq \|f(\cdot, t)\|^2.$$

Multiplying both sides by e^{-t}, this inequality can be rewritten as follows:

$$\frac{\mathrm{d}}{\mathrm{d}t}\left(\mathrm{e}^{-t}\|u(\cdot,t)\|^2\right) + \mathrm{e}^{-t}\int_{\Gamma_+}\left[\sum_{i=1}^{n} b_i(x)\nu_i(x)\right]u^2(x,t)\,\mathrm{d}s(x) \le \mathrm{e}^{-t}\|f(\cdot,t)\|^2.$$

Integrating this inequality w.r.t. t and noting the initial condition (3),

$$\mathrm{e}^{-t}\|u(\cdot,t)\|^2 + \int_0^t \mathrm{e}^{-\tau}\int_{\Gamma_+}\left[\sum_{i=1}^{n} b_i(x)\nu_i(x)\right]u^2(x,\tau)\,\mathrm{d}s(x)\,\mathrm{d}\tau$$

$$\le \|u_0\|^2 + \int_0^t \mathrm{e}^{-\tau}\|f(\cdot,\tau)\|^2\,\mathrm{d}\tau, \quad t\in[0,T].$$

It therefore follows that

$$\|u(\cdot,t)\|^2 + \int_0^t \mathrm{e}^{t-\tau}\int_{\Gamma_+}\left[\sum_{i=1}^{n} b_i(x)\nu_i(x)\right]u^2(x,\tau)\,\mathrm{d}s(x)\,\mathrm{d}\tau$$

$$\le \mathrm{e}^t\|u_0\|^2 + \int_0^t \mathrm{e}^{t-\tau}\|f(\cdot,\tau)\|^2\,\mathrm{d}\tau, \quad t\in[0,T]. \tag{6}$$

This 'energy inequality' implies continuous dependence of the solution on the data. It also implies uniqueness of the solution.

Let us consider the important special case when

$$c \equiv 0, \quad f \equiv 0, \quad \text{and } \operatorname{div} b = \sum_{i=1}^{n} \frac{\partial b_i}{\partial x_i} \equiv 0,$$

where $b(x) = (b_1(x), \dots, b_n(x))$. Then, thanks to the identity (5),

$$\frac{1}{2} \frac{\mathrm{d}}{\mathrm{d}t} \|u(\cdot, t)\|^2 + \frac{1}{2} \int_{\Gamma_+} [b(x) \cdot \nu(x)] u^2(x, t) \, \mathrm{d}s(x) = 0,$$

and therefore,

$$\|u(\cdot, t)\|^2 + \int_0^t \int_{\Gamma_+} [b(x) \cdot \nu(x)] u^2(x, \tau) \, \mathrm{d}s(x) \, \mathrm{d}\tau = \|u_0\|^2,$$

which can be viewed as an identity expressing 'conservation of energy' for the initial-boundary-value problem (1)–(3).

Explicit finite difference approximation

We focus on a special case of the problem: the constant-coefficient hyperbolic equation in one space dimension

$$\frac{\partial u}{\partial t} + b\frac{\partial u}{\partial x} = f(x, t), \quad x \in (0, 1), \ t \in (0, T], \tag{7}$$

subject to the boundary and initial conditions

$$u(x, t) = 0, \qquad x \in \Gamma_-, \ t \in [0, T],$$
$$u(x, 0) = u_0(x), \quad x \in [0, 1].$$

If $b > 0$ then $\Gamma_- = \{0\}$, and if $b < 0$ then $\Gamma_- = \{1\}$. Let us assume, for example, that $b > 0$. Then the appropriate boundary condition is

$$u(0, t) = 0, \quad t \in [0, T].$$

Let $\Delta x := 1/J$ be the mesh-size in the x-direction and $\Delta t := T/M$ the mesh-size in the time-direction, t. Let us also define

$$x_j := j\Delta x, \ \ j = 0, \dots, J, \qquad t_m := m\,\Delta t, \ \ m = 0, \dots, M.$$

At the mesh-point (x_j, t_m), (7) is approximated by the explicit scheme

$$\frac{U_j^{m+1} - U_j^m}{\Delta t} + b\, D_x^- U_j^m = f(x_j, t_m), \qquad j = 1, \dots, J,$$
$$m = 0, \dots, M-1,$$

subject to the boundary and initial conditions

$$U_0^m = 0, \qquad m = 0, \dots, M,$$
$$U_j^0 = u_0(x_j), \qquad j = 0, \dots, J.$$

Equivalently, this can be written as follows:

$$U_j^{m+1} = (1 - \mu)U_j^m + \mu U_{j-1}^m + \Delta t\, f(x_j, t_m), \qquad \left\{ \begin{array}{l} j = 1, \ldots, J, \\ m = 0, \ldots, M - 1, \end{array} \right.$$

together with

$$U_0^m = 0, \qquad m = 0, \ldots, M,$$
$$U_j^0 = u_0(x_j), \qquad j = 0, \ldots, J,$$

where

$$\mu = \frac{b \Delta t}{\Delta x};$$

μ is called the Courant–Friedrichs–Lewy (CFL) number.

The explicit finite difference scheme is frequently called the first-order upwind scheme.

Stability in the discrete maximum norm. Suppose that $0 \leq \mu \leq 1$; then

$$
\begin{aligned}
\left| U_j^{m+1} \right| &\leq (1-\mu) \left| U_j^m \right| + \mu \left| U_{j-1}^m \right| + \Delta t \left| f(x_j, t_m) \right| \\
&\leq (1-\mu) \max_{0 \leq j \leq J} \left| U_j^m \right| + \mu \max_{1 \leq j \leq J+1} \left| U_{j-1}^m \right| + \Delta t \max_{0 \leq j \leq J} \left| f(x_j, t_m) \right| \\
&= \max_{0 \leq j \leq J} \left| U_j^m \right| + \Delta t \max_{0 \leq j \leq J} \left| f(x_j, t_m) \right|.
\end{aligned}
$$

Thus we have that

$$
\max_{0 \leq j \leq J} \left| U_j^{m+1} \right| \leq \max_{0 \leq j \leq J} \left| U_j^m \right| + \Delta t \max_{0 \leq j \leq J} \left| f(x_j, t_m) \right|.
$$

Let us define the mesh-dependent norm

$$
\| U \|_\infty := \max_{0 \leq j \leq J} \left| U_j \right|;
$$

then

$$
\| U^{m+1} \|_\infty \leq \| U^m \|_\infty + \Delta t \| f(\cdot, t_m) \|_\infty, \quad m = 0, \ldots, M-1.
$$

Summing through m, we get

$$\max_{1\leq k\leq M}\|U^k\|_\infty \leq \|U^0\|_\infty + \sum_{m=0}^{M-1}\Delta t\|f(\cdot,t_m)\|_\infty,$$

which expresses stability of the scheme, assuming that

$$0 \leq \mu = \frac{b\Delta t}{\Delta x} \leq 1.$$

Thus, the finite difference scheme is conditionally stable in the maximum norm, the condition being that the CFL number $\mu \in [0,1]$.

Stability in the discrete L_2 norm

It is possible to show that the scheme is also conditionally stable in the discrete L_2-norm, $\|\cdot\,]\|$, defined by

$$\| V \,]|^2 = \sum_{i=1}^{J} \Delta x\, V_i^2.$$

The associated inner product is

$$(V, W] := \sum_{i=1}^{J} \Delta x\, V_i W_i.$$

Since

$$U_j^m = \frac{U_j^m + U_{j-1}^m}{2} + \frac{U_j^m - U_{j-1}^m}{2},$$

and $U_0^m = 0$, it follows that

$$
\begin{aligned}
(D_x^- U^m, U^m] &= \sum_{j=1}^{J} \Delta x \frac{U_j^m - U_{j-1}^m}{\Delta x} U_j^m \\
&= \frac{1}{2} \sum_{j=1}^{J} \{(U_j^m)^2 - (U_{j-1}^m)^2\} + \frac{\Delta x}{2} \sum_{j=1}^{J} \Delta x \left(\frac{U_j^m - U_{j-1}^m}{\Delta x} \right)^2 \\
&= \frac{1}{2} (U_J^m)^2 + \frac{\Delta x}{2} \| D_x^- U^m \|^2.
\end{aligned}
\tag{8}
$$

In addition, since

$$U_j^m = \frac{U_j^{m+1} + U_j^m}{2} - \frac{U_j^{m+1} - U_j^m}{2}$$

for $m = 0, \ldots, M-1$, we have for such m that

$$\left(\frac{U^{m+1} - U^m}{\Delta t}, U^m \right] = \frac{1}{2\Delta t} \left(\|U^{m+1}]\|^2 - \|U^m]\|^2 \right) - \frac{\Delta t}{2} \left\| \frac{U^{m+1} - U^m}{\Delta t} \right\|^2. \quad (9)$$

By taking the $(\cdot, \cdot]$-inner product of scheme with U^m and using (8) and (9):

$$\|U^{m+1}]\|^2 + \Delta t\, b\, |U_J^m|^2 + b\, \Delta x\, \Delta t\, \|D_x^- U^m]\|^2 - \|U^m]\|^2$$
$$- \Delta t^2 \left\| \frac{U^{m+1} - U^m}{\Delta t} \right\|^2 = 2\Delta t\, (f^m, U^m], \qquad m = 0, \ldots, M-1. \quad (10)$$

First suppose that $f \equiv 0$; then,

$$\frac{U^{m+1} - U^m}{\Delta t} = -b\, D_x^- U^m,$$

and by substituting this into the last term on the left-hand side of the equality (10) we have that, for $m = 0, \ldots, M-1$,

$$\|U^{m+1}]\|^2 + \Delta t\, b\, |U_J^m|^2 + b\, \Delta x\, \Delta t\, (1 - \mu)\|D_x^- U^m]\|^2 = \|U^m]\|^2.$$

Summing through m,

$$\|U^k]\|^2 + \sum_{m=0}^{k-1} \Delta t\, b\, |U_J^m|^2 + b\, \Delta x\, (1 - \mu) \sum_{m=0}^{k-1} \Delta t\, \|D_x^- U^m]\|^2 = \|U^0]\|^2 \qquad (11)$$

for $k = 1, \ldots, M$, which proves the stability of the scheme when $f \equiv 0$ for

$$0 \le \mu = \frac{b\, \Delta t}{\Delta x} \le 1.$$

In particular, if $\mu = 1$, we have from (11) that

$$\|U^k]|^2 + \sum_{m=0}^{k-1} \Delta t \, b \, |U_J^m|^2 = \|U^0]|^2, \quad k = 1, \ldots, M,$$

which is the discrete version of the identity (6), and expresses 'conservation of energy' in the discrete sense.

More generally, for $0 \le \mu \le 1$, (11) implies

$$\|U^k]|^2 + \sum_{m=0}^{k-1} \Delta t \, b \, |U_J^m|^2 \le \|U^0]|^2, \quad k = 1, \ldots, M.$$

Returning to the <u>general case $f \not\equiv 0$</u>, for all $k = 1, \ldots, M$,

$$\|U^k]\|^2 + \sum_{m=0}^{k-1} \Delta t\, b\, |U_J^m|^2 \le \mathrm{e}^T \left(\|U^0]\|^2 + \left(1 + \frac{\Delta t}{\epsilon}\right) \sum_{m=0}^{M-1} \Delta t \|f^m]\|^2 \right),$$

provided that, for some $\epsilon \in (0, 1)$,

$$0 \le \mu = \frac{b\,\Delta t}{\Delta x} \le 1 - \epsilon.$$

Proof. Since

$$\left\| \frac{U^{m+1} - U^m}{\Delta t} \right]\right|^2 = \|f^m - bD_x^- U^m]\|^2 \leq \{\|f^m]\| + b\|D_x^- U^m]\|\}^2$$
$$\leq \left(1 + \frac{1}{\epsilon'}\right) \|f^m]\|^2 + (1 + \epsilon')b^2\|D_x^- U^m]\|^2, \quad \epsilon' > 0,$$

and

$$(f^m, U^m] \leq \|f^m]\| \, \|U^m]\| \leq \tfrac{1}{2}\|f^m]\|^2 + \tfrac{1}{2}\|U^m]\|^2,$$

it follows from the equality (10) that

$$\|U^{m+1}]\|^2 + \Delta t \, b \, |U_J^m|^2 + b \, \Delta x \, \Delta t \left[1 - (1 + \epsilon')\frac{b\Delta t}{\Delta x}\right] \|D_x^- U^m]\|^2$$
$$\leq \Delta t \left[\left(1 + \frac{1}{\epsilon'}\right)\Delta t + 1\right] \|f^m]\|^2 + (1 + \Delta t)\|U^m]\|^2.$$

Let $\epsilon = 1 - 1/(1 + \epsilon') \in (0, 1)$ and assume the CFL condition

$$0 \le \mu = \frac{b\,\Delta t}{\Delta x} \le 1 - \epsilon,$$

we have, for $m = 0, \ldots, M - 1$, that

$$\|U^{m+1}]\|^2 + \Delta t\, b\, |U_J^m|^2 \le \|U^m]\|^2 + \Delta t \left(1 + \frac{\Delta t}{\epsilon}\right) \|f^m]\|^2 + \Delta t \|U^m]\|^2.$$

Upon summation of this inequality over $m = 0, \ldots, k - 1$, we deduce that

$$\|U^k]\|^2 + \left(\sum_{m=0}^{k-1} \Delta t\, b\, |U_J^m|^2\right) \le \|U^0]\|^2 + \left(1 + \frac{\Delta t}{\epsilon}\right) \sum_{m=0}^{k-1} \Delta t\, \|f^m]\|^2 \tag{12}$$
$$+ \sum_{m=0}^{k-1} \Delta t\, \|U^m]\|^2$$

for $k = 1, \ldots, M$.

We require the following lemma, which is easily proved by induction.

Lemma (Discrete Gronwall lemma)

Let (a_k), (b_k), (c_k) and (d_k) be four sequences of non-negative real numbers such that the sequence (c_k) is non-decreasing and

$$a_k + b_k \leq c_k + \sum_{m=0}^{k-1} d_m a_m, \quad k \geq 1; \quad a_0 + b_0 \leq c_0.$$

Then,

$$a_k + b_k \leq c_k \exp\left(\sum_{m=0}^{k-1} d_m\right), \quad k \geq 1.$$

By applying this lemma to the inequality (12) with

$$a_k := \|U^k]\|^2, \quad k \geq 0,$$

$$b_k := \sum_{m=0}^{k-1} \Delta t \, b \, |U_J^m|^2, \quad k \geq 1; \quad b_0 = 0,$$

$$c_k := \|U^0]\|^2 + \left(1 + \frac{\Delta t}{\epsilon}\right) \sum_{m=0}^{k-1} \Delta t \, \|f^m]\|^2, \quad k \geq 1; \quad c_0 = \|U^0]\|^2,$$

$$d_k := \Delta t, \quad k = 1, 2, \ldots, M,$$

we obtain for $k = 1, \ldots, M$:

$$\|U^k]\|^2 + \sum_{m=0}^{k-1} \Delta t \, b \, |U_J^m|^2 \leq e^{t_k} \left(\|U^0]\|^2 + \left(1 + \frac{\Delta t}{\epsilon}\right) \sum_{m=0}^{k-1} \Delta t \|f^m]\|^2 \right),$$

where $t_k = k\Delta t$. Hence we deduce stability of the scheme, in the sense that

$$\max_{1 \leq k \leq M} \left(\|U^k]\|^2 + \sum_{m=0}^{k-1} \Delta t \, b \, |U_J^m|^2 \right) \leq e^T \left(\|U^0]\|^2 + \left(1 + \frac{\Delta t}{\epsilon}\right) \sum_{m=0}^{M-1} \Delta t \|f^m]\|^2 \right).$$